



## Research Paper

## Thermal design of rectangular microscale inorganic light-emitting diodes

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## HIGHLIGHTS

- A 3D analytical model is developed to study the thermal properties of rectangular  $\mu$ -ILEDs.
- A scaling law for normalized temperature increase of a single  $\mu$ -ILED is established.
- Thermal properties of rectangular  $\mu$ -ILED array are investigated.
- Design guidelines for thermal management of  $\mu$ -ILEDs are provided.

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## ABSTRACT

The recently developed microscale inorganic light-emitting diodes ( $\mu$ -ILEDs) have attracted much attention due to their potential use in biointegrated applications such as optogenetics. It is critical to understand the thermal properties of  $\mu$ -ILEDs since excessive heating may reduce the performance significantly. A three-dimensional analytical model based on the Fourier Cosine transform is developed to study the thermal properties of rectangular  $\mu$ -ILEDs in this paper. The analytical prediction agrees well with finite element analysis. A scaling law for the normalized temperature increase of a single  $\mu$ -ILED is established in terms of four non-dimensional parameters: the normalized shape factor  $\gamma = b/a$ , the normalized  $\mu$ -ILED thickness  $H_{LED} = H_{LED}/a$ , and  $\chi_1 = k_m H_m / (k_s a)$  and  $\chi_2 = k_s H_B / (k_B a)$ , where  $a$  and  $b$  are the half-lengths of the  $\mu$ -ILED,  $k$  and  $H$  are the thermal conductivity and the thickness with the subscripts LED,  $m$ ,  $B$  and  $s$  for the  $\mu$ -ILED, metal layer, encapsulation layer and substrate, respectively. The influences of these non-dimensional parameters on the normalized temperature increase are systematically investigated. The temperature increase for rectangular  $\mu$ -ILED array is then obtained analytically by the method of superposition. These results provide design guidelines to minimize the adverse thermal response of rectangular  $\mu$ -ILEDs.

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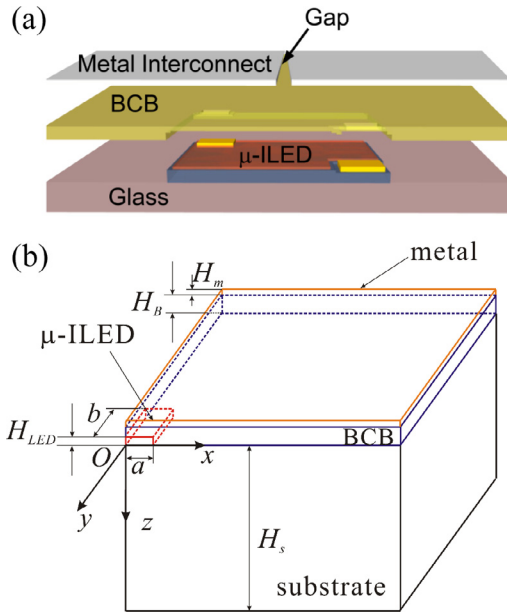
## 1. Introduction

Microscale inorganic light-emitting diodes ( $\mu$ -ILEDs) with characteristics in size, efficiency and lifetime enable many novel applications, which are not possible using conventional LEDs, such as stretchable optoelectronics [1–5], optogenetics [6–10] and solid-state lighting.

The recent strategy by Kim et al. [11] to fabricate  $\mu$ -ILEDs has received much attention. It involves the growth of high quality epitaxial material on sapphire or Si substrates followed by the definition of arrays of individual  $\mu$ -ILED. Transfer printing is then used to release the completed  $\mu$ -ILED array from the grown substrate onto other unusual substrates. A photodefinable layer of benzocyclobutene (BCB) encapsulates the  $\mu$ -ILED array with the p- and n-type ohmic contacts of each  $\mu$ -ILED exposed for electrical interconnection. A metal layer (e.g., aluminum) is then deposited on the top to serve as interconnects and heat spreader. This strategy is applicable to interconnect multiple  $\mu$ -ILEDs over a large area to create  $\mu$ -ILED arrays on unusual substrates. Fig. 1(a) shows the

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**Fig. 1.** (a) Schematic diagram of the layout for  $\mu$ -ILED system; (b) Schematic illustration of the 3D analytical model showing a quarter geometry of the  $\mu$ -ILED system.

schematic diagram of a single rectangular  $\mu$ -ILED encapsulated by BCB and metal layers on a glass substrate.

It is critical to understand the thermal properties of  $\mu$ -ILEDs to avoid excessive heating, which could induce significant reduction in the  $\mu$ -ILED efficiency. Many researchers have performed thermal analysis of  $\mu$ -ILEDs. Lü et al. [12] studied the heat conduction of square  $\mu$ -ILEDs with layouts in Fig. 1a under a constant power analytically. For simplicity, an axisymmetric model, instead of three-dimensional model, is established to capture the square  $\mu$ -ILED temperature quantitatively. The simple analytic expression for the  $\mu$ -ILED temperature is very useful in the design optimization of  $\mu$ -ILEDs to minimize the adverse thermal effects. The operation of  $\mu$ -ILEDs in a pulsed mode could further reduce the temperature of  $\mu$ -ILEDs. Kim et al. [13] investigated the thermal properties of square  $\mu$ -ILEDs on various substrates under a pulsed power experimentally. Li et al. [14] then established an analytic axisymmetric model to explore the dependence of the  $\mu$ -ILED temperature on the geometric dimensions, material properties and loading parameters. All above studies are for square  $\mu$ -ILEDs and the analytic models, reviewed in [15], adopted the axisymmetric treatment. Although these axisymmetric models can predict the temperature distributions accurately for square  $\mu$ -ILEDs, they cannot be applied for rectangular  $\mu$ -ILEDs, which are more general in practical applications.

Recently, three-dimensional models were established to study the thermal performance of  $\mu$ -ILEDs. Zhang et al. [16] studied the thermal properties of wirelessly powered  $\mu$ -ILEDs, demonstrated for potential applications in implantable devices [17], by coupling the one-dimensional heat conduction in the straight inductive coil, and three-dimensional heat conduction in the substrate. Cui et al. [18] presented a three-dimensional analytic model to illustrate the thermal properties of single rectangular  $\mu$ -ILED under a pulsed power. These studies expand the application scope which is limited by axisymmetric model. However, the thermal behaviors of rectangular  $\mu$ -ILEDs and rectangular  $\mu$ -ILED arrays under a constant power are not studied yet, which are important in the application of  $\mu$ -ILEDs in solid state lighting.

The objective of this study is to develop a rigorous analytical model, validated by finite element analysis, to perform three-dimensional thermal analysis of rectangular  $\mu$ -ILEDs and provide

explicit expressions for  $\mu$ -ILED temperature increase in terms of geometric, material and loading parameters.

## 2. Three-dimensional thermal analysis of a single rectangular $\mu$ -ILED

We take rectangular  $\mu$ -ILEDs in experiments [11] with the same layouts in Fig. 1a as an example to illustrate our approach. Due to the symmetry of the system, only a quarter of the system is considered. Fig. 1(b) shows the schematic diagram of the analytical model with the  $\mu$ -ILED on the glass substrate encapsulated by BCB and metal layers. The thicknesses of metal, BCB and substrate are denoted by  $H_m$ ,  $H_B$  and  $H_s$ , respectively. The in-plane dimension of  $\mu$ -ILED is  $2a \times 2b$ , and the thickness is  $H_{LED}$ . The  $\mu$ -ILED is modeled as a planar heat source with the input power  $P$  at the BCB/substrate interface [12,14]. This treatment enables analytical investigation of thermal properties of  $\mu$ -ILEDs, which is to be validated by three-dimensional finite element analysis (FEA).

The coordinate system is established with the origin at the center of  $\mu$ -ILED,  $x$  axis along the side  $a$ ,  $y$  axis along the side  $b$ , and  $z$  axis pointing from the top metal layer to the bottom substrate. The temperature increase  $\Delta T = T(x, y, z) - T_0$  from the ambient temperature satisfies the three-dimensional steady-state heat conduction equation

$$\frac{\partial^2 \Delta T}{\partial x^2} + \frac{\partial^2 \Delta T}{\partial y^2} + \frac{\partial^2 \Delta T}{\partial z^2} = 0. \quad (1)$$

The boundary and continuity conditions are given in the following. At the top surface of metal layer [ $z = -(H_m + H_B)$ ], the natural convection condition gives

$$-k_m \frac{\partial \Delta T}{\partial z} \Big|_{z=-(H_m+H_B)} = -h \cdot \Delta T \Big|_{z=-(H_m+H_B)}, \quad (2)$$

where  $k_m$  is the thermal conductivity of metal and  $h$  is the coefficient of heat convection. The continuity of temperature and heat flux across the metal/BCB interface ( $z = -H_B$ ) gives

$$\Delta T \Big|_{z=-H_B^+} = \Delta T \Big|_{z=-H_B^-} \text{ and } -k_m \frac{\partial \Delta T}{\partial z} \Big|_{z=-H_B^+} = -k_B \frac{\partial \Delta T}{\partial z} \Big|_{z=-H_B^-}, \quad (3)$$

where  $k_B$  is the thermal conductivity of BCB. At the BCB/substrate interface ( $z = 0$ ), the temperature is continuous while the heat flux is only continuous outside the  $\mu$ -ILED and satisfies the heat source condition inside the  $\mu$ -ILED, which give

$$\Delta T \Big|_{z=0^+} = \Delta T \Big|_{z=0^-} \text{ and } k_B \frac{\partial \Delta T}{\partial z} \Big|_{z=0^-} - k_s \frac{\partial \Delta T}{\partial z} \Big|_{z=0^+} = \begin{cases} 0, & (x, y) \notin D \\ \frac{P}{4ab}, & (x, y) \in D \end{cases} \quad (4)$$

where  $k_s$  is the thermal conductivity of substrate,  $D$  is the inside region of the  $\mu$ -ILED at the BCB/substrate interface, i.e.,  $D = \{(x, y), 0 \leq x \leq a, 0 \leq y \leq b\}$ . The ambient temperature at the bottom surface ( $z = H_s$ ) of substrate gives

$$\Delta T \Big|_{z=H_s} = 0. \quad (5)$$

Through the Fourier Cosine transform along  $x$  and  $y$  directions,

$$\hat{\Delta T}(x, \beta, z) = \int_0^\infty \int_0^\infty \Delta T(x, y, z) \cos(\alpha x) \cos(\beta y) dx dy, \quad (6)$$

Eqs. (1)–(5) become

$$\frac{\partial^2 \hat{\Delta T}}{\partial z^2} - (\alpha^2 + \beta^2) \hat{\Delta T} = 0, \quad (7)$$

$$-k_m \frac{\partial \hat{\Delta T}}{\partial z} \Big|_{z=-(H_m+H_B)} = -h \cdot \hat{\Delta T} \Big|_{z=-(H_m+H_B)}, \quad (8)$$

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