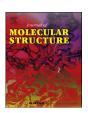
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Geometrical and electronic structure of the molybdenum and tungsten halides MX_3 and MX_4 (M = Mo, W; X = F, Cl): Jahn-Teller effect and spin-orbit coupling[★]

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ABSTRACT

The ground and lower-lying excited electronic states of MX_3 and MX_4 (M = Mo, W; X = F, Cl) molecules were systematically studied by the complete active space self-consistent field (CASSCF) and multiconfigurational quasi-degenerate second-order perturbation (XMCQDPT2) methods. Scalar-relativistic effects and spin-orbit coupling were taken into account employing the third-order Douglas-Kroll-Hess (DKH) Hamiltonian and full Breit-Pauli operator. The ground orbital state of MoF3 and MoCl3 complexes is quadruplet ⁴A₂' state in the equilibrium configuration with D_{3h} symmetry. If the spin-orbit interaction is taken into account the calculations result in two ${}^4E_{1/2}$ and ${}^4E_{3/2}$ spin-orbit states with 11–14 cm⁻¹ energy differences instead of one ⁴A₂' state. Spin-orbit interaction quenches the Jahn-Teller effect in ²E" ground orbital state of WF₃ complex in D_{3h} configuration, and C_{2v} configuration with $\binom{4}{E_{1/2}} + \binom{4}{E_{3/2}}$ spinmixed state corresponds to the minimum on the PES. The WCl₃ complex reveals strong interaction of the spin-orbit states of different multiplicity. For WCl3 complex the configuration with D3h symmetry corresponds to the minimum on the PES for four lowest spin-mixed electronic states. The MoF₄, MoCl₄ and WCl₄ complexes possess the tetrahedral equilibrium configuration with ³A₂ orbital electronic state. If the spin-orbit interaction is taken into account the ground state changes to the threefold degenerated ³T₂ spin-orbit state with weak Jahn-Teller effect. In the case of WF4 complex the strong Jahn-Teller effect has been discovered in the first excited ^{1}E singlet electronic state of the tetrahedral structure, and D_{2d} configuration with ¹A₁ electronic state possesses the lowest energy. The comparison with available experimental data has been performed. The tendencies in molecular parameters of complexes studied were analyzed.

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1. Introduction

The molybdenum and tungsten halides MX_3 and MX_4 (M = Mo, W; X = F, Cl) have attracted considerable interest because of their important role as the precursors and catalysts in many chemical reactions. Some of these compounds are currently used as transients in the halide method of refinement. The molybdenum and tungsten are widely used as the construction materials for different engines, which operate at high temperature and in the aggressive halogen containing medium. Consequently there has been interest in the molecular structure of the molybdenum and tungsten

forecasting a constructional material behavior in corresponding conditions. On the other hand, special attention has been paid to a molecular and electronic structure of the molybdenum and tungsten halides. The close energies of the electronic states with different multiplicities, the combination of strong spin-orbit coupling and Jahn-Teller effect are to be expected for these complexes. As a result, the high-symmetry structure can be distorted in an unpredictable way.

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According to mass spectrometric study [1-3] the molybdenum and tungsten halides exist in the gas phase as monomeric molecules MX_n (n = 1–6) and their polymers. The molecular structure of free molecules of several molybdenum and tungsten halides (WCl₃[4], MoF₄[5], MoCl₄[6], MoBr₄[6,7], WCl₄[8]) has been studied by gas-phase electron diffraction (GED). The authors of [4] conclude that WCl_3 possesses T-like structure of $C_{2\nu}$ symmetry. The

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Dedicated to Prof. Georgiy V. Girichev on the Occasion of his 70th birthday.

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tetrafluoride and tetrachloride of molybdenum were found to be tetrahedral [5,6], whereas equilibrium structure of WCl₄ corresponds to C_{2v} symmetry [8]. Two different GED studies of molybdenum tetrabromide have been reported with rather conflicting conclusions about the structure of MoBr₄: either with tetrahedral T_d [6] or C_{2v} [8] symmetry. Vibrational IR spectra were recorded for the products of MoF₆ photolysis in argon matrix [9], and for the species obtained by transmission of F_2 gas-phase over the molybdenum and tungsten heated at 100-1100 °C and isolated in argon matrix [10]. According to the spectral data in case of molybdenum the gas phase consists of MoF_n (n = 3,4,5,6) [9,10] molecules, trimers (MoF₅)₃[10] and MoOF₄[10]. In case of tungsten only WF₆ and WOF₄ are present in the gas phase [10]. The bands observed in IR spectra were assigned in terms of tetrahedral T_d configuration for MoF₄[9,10] and symmetric D_{3h} configuration for MoF₃[10].

The first quantum chemical (QC) studies (ROHF, MP2, CASSCF + SOCI) of MoF3, WF3, MoF4 and WF4 were performed in 2000 [11,12]. One year later the results of DFT/B3LYP study of WF3[13], WCl3[14], WF4[13], WCl4[14] were published. According to QC data [11,13,14] high-spin electronic state $^4\text{A}'_2$ with D3h geometric configuration is the ground state of MoF3, WF3 and WCl3 complexes. The first excited electronic state for these molecules was predicted to be Jahn-Teller $^2\text{E}''$ term [11,13,14]. The ground and first excited electronic states of WF4 were found to be $^1\text{A}_1$ (D2d) and $^3\text{A}_2$ (Td), respectively [12,13]. On the contrary, in case of MoF4 and WCl4 triplet $^3\text{A}_2$ with tetrahedral Td configuration corresponds to the ground state, whereas singlet $^1\text{A}_1$ (D2d) is the first excited state [12,14].

Thus, the available up to date information about geometrical and electronic structure of the molybdenum and tungsten halides is poor and imperfect. The effects of spin-orbit coupling (SOC) have not been studied yet. The aim of the present work is a systematic study of the geometrical and electronic structure of MX_3 and MX_4 molecules (M=Mo,W;X=F,Cl) in ground and low-lying excited electronic states by quantum-chemical methods. The influence of spin-orbit coupling (SOC) on the molecular parameters will be analyzed.

2. Preliminary consideration

According to the authors of [11,12] the molybdenum and tungsten tri- and tetrahalides are the complexes of strong ligand field, and the low-spin states are stabilized. Furthermore, in this case the low-spin and high-spin states can possess close energies. The ground and lower-lying excited electronic states of MX₃ and MX₄ molecules, which can be considered as ionic systems $[M^{n+}][X^-]_n$, arise from the states of free cations with open valence d-shells, d^3 and d^2 for M^{3+} and M^{4+} , respectively. Table 1 contains electronic states of M^{3+} and M^{4+} cations along with the corresponding states of MX_3 and MX_4 complexes. The correlations of the orbital and spin-orbit electronic states are shown in Table 2. Possible electronic

states for cations and complexes (Tables 1 and 2) have been determined using group theory and the character tables [15].

3. Calculation details

All calculations were performed using the complete active space (CASSCF) method by GAMESS US program set [16]. The active space included all d-electrons of M³⁺ or M⁴⁺ cation (see Table 1) on five molecular orbitals consisting mainly of valence d-metal atom orbitals. A dynamic correlation was taken into account at the level of XMCQDPT2 multiconfiguration quasi-degenerate second-order perturbation theory [17]. The doubly occupied orbitals corresponding to 1s- orbitals of F, Cl, Mo, W; 2s- and 2p- orbitals of Cl, Mo, W; and 3s-, 3p- and 3d-orbitals of Mo, W were frozen in XMCQDPT2 calculations. Scalar-relativistic effects were taken into account employing the third-order Douglas-Kroll-Hess (DKH) Hamiltonian [18,19]. The largest relativistic effects can be expected for the electronic shells of d-metal atoms. Therefore, these effects were considered only for Mo and W. Spin-orbit coupling (SOC) was taken into account using wave function from XMCQDPT2 calculations (SO-XMCQDPT2). The contribution of SOC to the Hamiltonian was described by full Breit-Pauli operator [20,21]. The problem of small-energy denominators that may appear in a perturbative expansion was solved by the intruder state avoidance (ISA) method [22]. The optimal value of b parameter defining the energydenominator shifts was chosen to be 0.02 and 0.1 [22] for XMCQDPT2 and SO-XMCQDPT2 calculations, respectively. All CASSCF and XMCQDPT2 calculations have been carried out using state-average wave functions. For each multiplicity the averaging was performed over all electronic states (see Table 1) with equal weight coefficients. SPKrATZP basis sets [23-28] were used for molecular properties calculations. The geometry optimization, the calculations of force constants and harmonic vibrational frequencies were performed by the method of numerical differentiation of energies in the symmetry coordinates using ANOCOR program [29]. The vibrational symmetry coordinates of MX₃ and MX₄ molecules with different configurations (Fig. 1) are listed in Table S1.

The energy levels of free cations M³⁺ were calculated and compared with available experimental data [30] (see Table 3). The calculations with SOC taken into account were performed both for every multiplicity separately (spin-orbit state) and with consideration of interaction of spin-orbit states with different multiplicity (spin-mixed states). For Mo³⁺ cation the states with different multiplicity take part only in wave functions of four electronic states (see Table 3). In these cases the differences between calculated energies of spin-orbit and spin-mixed states are not greater than 740 cm⁻¹. For W³⁺ cation practically all states possesses spin-mixed wave functions. Furthermore, the differences in calculated energies of W³⁺ for spin-orbit and spin-mixed states reach several thousand cm⁻¹ (Table 3). The average deviations of the calculated

 Table 1

 Electronic states of free M^{3+} and M^{4+} cations, degree of degeneracy (n) and orbital electronic states of MX_3 and MX_4 molecules.

$M^{3+}(d^3)$		MX_3	$M^{4+}(d^2)$		MX_4	
Term	n	D _{3h} structure	Term	n	T _d structure	D _{4h} structure
⁴ F	7	${}^{4}A_{2}' + {}^{4}E' + {}^{4}A_{1}'' + {}^{4}A_{2}'' + {}^{4}E''$	³F	7	$^{3}A_{2} + ^{3}T_{1} + ^{3}T_{2}$	$^{3}A_{2g} + ^{3}B_{1g} + ^{3}B_{2g} + 2^{3}E_{g}$
⁴ P	3	${}^{4}\text{A}_{2}' + {}^{4}\text{E}''$				-0 -0 0
² P	3	${}^{2}A_{2}' + {}^{2}E''$	³ P	3	³ T ₁	${}^{3}A_{2g} + {}^{3}E_{g}$
² D	5	${}^{2}A_{1}' + {}^{2}E' + {}^{2}E''$	¹ S	1	¹ A ₁	$^{1}A_{1g}$
² D	5	${}^{2}A_{1}' + {}^{2}E' + {}^{2}E''$	¹ D	5	$^{1}E + {}^{1}T_{2}$	${}^{1}A_{1g}^{1g} + {}^{1}B_{1g} + {}^{1}B_{2g} + {}^{1}E_{g}$
² F	7	${}^{2}A_{2}' + {}^{2}E' + {}^{2}A_{1}'' + {}^{2}A_{2}'' + {}^{2}E''$. 2	15 15 25 5
^{2}G	9	${}^{2}A_{1}' + 2{}^{2}E' + {}^{2}A_{1}'' + {}^{2}A_{2}'' + {}^{2}E''$	^{1}G	9	${}^{1}A_{1} + {}^{1}E + {}^{1}T_{1} + {}^{1}T_{2}$	$2^{1}A_{1g}+{}^{1}A_{2g}+{}^{1}B_{1g}+{}^{1}B_{2g}+2^{1}E_{g}\\$
² H	11	${}^{2}A_{2}' + 2{}^{2}E' + {}^{2}A_{1}'' + {}^{2}A_{2}'' + 2{}^{2}E''$.5 25 .5 25 .

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