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# Time-stepping error bounds for fractional diffusion problems with non-smooth initial data



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#### ABSTRACT

We apply the piecewise constant, discontinuous Galerkin method to discretize a fractional diffusion equation with respect to time. Using Laplace transform techniques, we show that the method is first order accurate at the *n*th time level  $t_n$ , but the error bound includes a factor  $t_n^{-1}$  if we assume no smoothness of the initial data. We also show that for smoother initial data the growth in the error bound as  $t_n$  decreases is milder, and in some cases absent altogether. Our error bounds generalize known results for the classical heat equation and are illustrated for a model problem.

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(1)

#### 1. Introduction

Consider an initial-value problem for an abstract, time-fractional diffusion equation [7, p. 84]

$$\partial_t u + \partial_t^{1-\nu} A u = 0$$
 for  $t > 0$ , with  $u(0) = u_0$  and  $0 < \nu < 1$ .

Here, we think of the solution u as a function from  $[0, \infty)$  to a Hilbert space  $\mathcal{H}$ , with  $\partial_t u = u'(t)$  the usual derivative with respect to t, and with

$$\partial_t^{1-\nu} u(t) = \frac{\partial}{\partial t} \int_0^t \frac{(t-s)^{\nu-1}}{\Gamma(\nu)} u(s) \, ds$$

the Riemann–Liouville fractional derivative of order  $1 - \nu$ . The linear operator A is assumed to be self-adjoint, positivesemidefinite and densely defined in  $\mathcal{H}$ , with a complete orthonormal eigensystem  $\phi_1, \phi_2, \phi_3, \ldots$ . We further assume that the eigenvalues of A tend to infinity. Thus,

 $A\phi_m = \lambda_m \phi_m, \quad \langle \phi_m, \phi_n \rangle = \delta_{mn}, \ 0 \le \lambda_1 \le \lambda_2 \le \lambda_3 \le \cdots,$ 

where  $\langle u, v \rangle$  is the inner product in  $\mathcal{H}$ ; the corresponding norm in  $\mathcal{H}$  is denoted by  $||u|| = \sqrt{\langle u, u \rangle}$ . In particular, we may take  $Au = -\nabla^2 u$  and  $\mathcal{H} = L_2(\Omega)$  for a bounded spatial domain  $\Omega$ , with u subject to homogeneous Dirichlet or Neumann boundary conditions on  $\partial \Omega$ . Our problem (1) then reduces to the classical heat equation when  $v \to 1$ .

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Many authors have studied techniques for the time discretization of (1), but obtaining sharp error bounds has proved challenging. In studies of explicit and implicit finite difference schemes [1,3,8,14,17,20] the error analyses typically assume that the solution u(t) is sufficiently smooth, including at t = 0, which amounts to imposing compatibility conditions on the initial data and source term. In our earlier work on discontinuous Galerkin (DG) time-stepping [11,15,16], we permitted more realistic behaviour, allowing the derivatives of u(t) to be unbounded as  $t \rightarrow 0$ , but were seeking error bounds that are uniform in t using variable time steps. In the present work, we again consider a piecewise-constant DG scheme but with a completely different method of analysis that leads to sharp error bounds even for non-smooth initial data, at the cost of requiring a constant time step  $\Delta t$ . Our previous analysis [11, Theorem 5] of the scheme (5), in conjunction with relevant estimates [10] of the derivatives of u, shows, in the special case of uniform time steps, only a sub-optimal error bound at the *n*th time level  $t_n = n\Delta t$ ,

$$\left\| U^n - u(t_n) \right\| \le C \Delta t^{r\nu} \left\| A^r u_0 \right\| \quad \text{for } 0 \le r < 1/\nu.$$
<sup>(2)</sup>

In our main result, we substantially improve on (2) by showing that

$$\|U^n - u(t_n)\| \le C t_n^{r\nu - 1} \Delta t \|A^r u_0\| \quad \text{for } 0 \le r \le \min(2, 1/\nu).$$
(3)

Thus, for a general  $u_0 \in \mathcal{H}$  the error is of order  $t_n^{-1}\Delta t$  at  $t = t_n$ , so the method is first-order accurate but the error bound includes a factor  $t_n^{-1}$  that grows if  $t_n$  approaches zero, until at  $t = t_1$  the bound is of order  $t_1^{-1}\Delta t = 1$ . However, if  $1/2 \leq \nu < 1$  and  $u_0$  is smooth enough to belong to  $D(A^{1/\nu})$ , the domain of  $A^{1/\nu}$ , then the error is of order  $\Delta t$ , uniformly in  $t_n$ . For  $0 < \nu < 1/2$ , no matter how smooth  $u_0$  a factor  $t_n^{2\nu-1}$  is present. To the best of our knowledge, only Cuesta et al. [2] and McLean and Thomée [12, Theorem 3.1] have hitherto investigated the time discretization of (1) for the interesting case when the initial data might not be regular, the former using a finite difference–convolution quadrature scheme and the latter a method based on numerical inversion of the Laplace transform.

In the present work, we do not discuss the spatial discretization of (1). By contrast, Jin, Lazarov and Zhou [6] studied the piecewise linear finite element solution  $u_h(t) \approx u(t)$  using a quasi-uniform partition of the spatial domain  $\Omega$  into elements with maximum diameter h, but with no time discretization. They worked with an alternative formulation of the fractional diffusion problem,

$$\partial_{t,C}^{\nu} u - \nabla^2 u = 0 \quad \text{for } x \in \Omega \text{ and } 0 < t \le T,$$
(4)

where  $\partial_{t,C}^{\nu}$  denotes the Caputo fractional derivative of order  $\nu$ , and proved [6, Theorems 3.5 and 3.7] that if  $u_h(0) = u_{0h}$ , for an appropriate  $u_{h0} \approx u_0$ , then

$$\|u_h(t) - u(t)\| + h \|\nabla(u_h - u)\| \le Ct^{\nu(r-1)} \times \begin{cases} h^2 \ell_h \|A^r u_0\|, & r \in \{0, 1/2\} \\ h^2 \|A^r u_0\|, & r = 1, \end{cases}$$

where  $\ell_h = \max(1, \log h^{-1})$ . These estimates for the spatial error complement our bounds for the error in a time discretization.

We define a piecewise-constant approximation  $U(t) \approx u(t)$  by applying the DG method [11,13],

$$U^{n} - U^{n-1} + \int_{t_{n-1}}^{t_{n}} \partial_{t}^{1-\nu} AU(t) dt = 0 \quad \text{for } n \ge 1, \text{ with } U^{0} = u_{0},$$
(5)

where  $U^n = U(t_n^-) = \lim_{t \to t_n^-} U(t)$  denotes the one-sided limit from below at the *n*th time level. Thus,  $U(t) = U^n$  for  $t_{n-1} < t \le t_n$ . Since we do not consider any spatial discretization, U is a semidiscrete solution with values in  $\mathcal{H}$ . A short calculation reveals that

$$\int_{t_{n-1}}^{t_n} \partial_t^{1-\nu} AU(t) dt = \Delta t^{\nu} \sum_{j=1}^n \beta_{n-j} AU^j,$$

with

$$\beta_0 = \Delta t^{-\nu} \int_{t_{n-1}}^{t_n} \frac{(t_n - t)^{\nu - 1}}{\Gamma(\nu)} dt = \frac{1}{\Gamma(1 + \nu)}$$

and, for  $j \ge 1$ ,

$$\beta_j = \Delta t^{-\nu} \int_{t_{n-j-1}}^{t_{n-j}} \frac{(t_n - t)^{\nu - 1} - (t_{n-1} - t)^{\nu - 1}}{\Gamma(\nu)} dt = \frac{(j+1)^{\nu} - 2j^{\nu} + (j-1)^{\nu}}{\Gamma(1+\nu)}$$

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