

Short Note

A primal formulation for the Helmholtz decomposition

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Available online 11 April 2007**Abstract**

In 1999, Jean-Paul Caltagirone and Jérôme Breil have developed in their paper [Caltagirone, J. Breil, Sur une méthode de projection vectorielle pour la résolution des équations de Navier–Stokes, C.R. Acad. Sci. Paris 327(Série II b) (1999) 1179–1184] a new method to compute a divergence-free velocity. They have used the **grad**(div) operator to extract the solenoidal part of a given vector field. In this contribution we explain how this method can be considered as a real Helmholtz decomposition and we present a stable approximation in the framework of spectral methods. Numerical results are presented to illustrate the efficiency of this approach.

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1. Introduction

The approximation of the **grad**(div) operator pervades many applied physics domains. Besides the ideal ocean wave problem without Coriolis force and no friction [15], it arises in the Maxwell equations [11] and in the Navier–Stokes equations for fluid flow problems when using a penalty formulation for the incompressibility condition [14]. The problem also arises in the ideal linear magneto hydrodynamics equations when computing the stability behavior of a fusion plasma device [16]. Another original application of this operator was introduced by J.P. Caltagirone and J. Breil in their paper [13] where they used this operator to extract from a given velocity field its solenoidal part. These authors had christened it *vector projection* which consists in solving the following problem: Let $\mathbf{u}^*$ be a non-divergence free velocity field, find a couple of vector fields $(\mathbf{u}, \mathbf{v})$ such that

$$-\nabla(\nabla \cdot \mathbf{u}) = \nabla(\nabla \cdot \mathbf{u}^*), \quad \text{in } \Omega, \quad (1.1)$$

$$\mathbf{u} \cdot \mathbf{n} = 0, \quad \text{on } \partial\Omega, \quad (1.2)$$

$$\mathbf{v} = \mathbf{u} + \mathbf{u}^*, \quad \text{in } \Omega, \quad (1.3)$$

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where $\mathbf{v}$ and $\mathbf{u}$ are respectively divergence-free and curl-free. Here $\Omega \subset \mathbb{R}^d$ ($d = 2, 3$) is a simply connected and bounded domain with Lipschitzian border. $\mathbf{n}$ denotes the outer unit normal along the boundary.

The objective of this note is on the one hand to explain how the previous system can be considered as a Helmholtz decomposition step and on the other hand to present a stable discretization in the framework of spectral methods. We end this note by presenting a relevant numerical experiment.

Some notations – The symbol $L^2(\Omega)$ stands for the usual Lebesgue space and $H^1(\Omega)$, the Sobolev space, involves all the functions that are, together with their gradient, in $L^2(\Omega)$. The $\mathcal{C}(\Omega)$ denotes the space of continuous functions defined in Ω .

2. Continuous problems and their variational formulations

In order to write the continuous problem in its variational form we introduce the relevant spaces of functions.

Let $H(\text{div}, \Omega)$ denote the space (see [12])

$$H(\text{div}, \Omega) = \{\mathbf{w} \in (L^2(\Omega))^d; \text{div } \mathbf{w} \in L^2(\Omega)\},$$

endowed with the natural norm

$$\|\mathbf{w}\|_{H(\text{div}, \Omega)} = (\|\mathbf{w}\|_{(L^2(\Omega))^d}^2 + \|\text{div } \mathbf{w}\|_{L^2(\Omega)}^2)^{1/2}.$$

The continuous problem we consider reads: *Find $\mathbf{u}$ in $H(\text{div}, \Omega)$ it such that:*

$$-\nabla(\nabla \cdot \mathbf{u}) = \mathbf{f}, \quad \text{in } \Omega, \quad (2.4)$$

$$\mathbf{u} \cdot \mathbf{n} = 0, \quad \text{on } \partial\Omega, \quad (2.5)$$

where $\mathbf{f}$ is a given data.

Since $\text{curl}(\mathbf{grad} \cdot) \equiv \mathbf{0}$ we notice that a necessary condition for the existence of a solution to problems (2.4) and (2.5) is that $\text{curl } \mathbf{f} = \mathbf{0}$ and by consequence we can state the existence of a function $\varphi(x, y)$ such that

$$\mathbf{f} = \mathbf{grad} \varphi.$$

This leads to restate the basic problem as: *For a given $\varphi \in L_0^2(\Omega)$, find $\mathbf{u} \in H(\text{div}, \Omega)$ such that*

$$-\nabla(\nabla \cdot \mathbf{u}) = \nabla \varphi, \quad \text{in } \Omega, \quad (2.6)$$

$$\mathbf{u} \cdot \mathbf{n} = 0, \quad \text{on } \partial\Omega, \quad (2.7)$$

where $L_0^2(\Omega)$ denotes the $L^2(\Omega)$ subspace of functions having zero average values. This formulation is equivalent to the dual one that reads: *For a given $\varphi \in L_0^2(\Omega)$ find $\mathbf{u} \in X(\Omega)$ and $\psi \in L_0^2(\Omega)$ such that:*

$$\mathbf{u} - \nabla \psi = 0, \quad \text{in } \Omega, \quad (2.8)$$

$$-\nabla \cdot \mathbf{u} = \varphi, \quad \text{in } \Omega, \quad (2.9)$$

$$\mathbf{u} \cdot \mathbf{n} = 0, \quad \text{on } \partial\Omega, \quad (2.10)$$

where

$$X(\Omega) = \{\mathbf{w} \in H(\text{div}, \Omega); \mathbf{w} \cdot \mathbf{n} = 0 \text{ on } \partial\Omega\}.$$

This dual formulation can be rewritten as a classical Helmholtz decomposition, indeed: Let $\mathbf{u}^*$ and $\mathbf{v}$ be two vector fields such that $\nabla \cdot \mathbf{u}^* = \varphi$, $\mathbf{u}^* \cdot \mathbf{n} = 0$ and $\mathbf{v} = \mathbf{u} + \mathbf{u}^*$. The problem (2.8)–(2.10) then becomes: *Find $\mathbf{v} \in X(\Omega)$ and ψ in $L_0^2(\Omega)$ such that*

$$\mathbf{v} - \nabla \psi = \mathbf{u}^*, \quad \text{in } \Omega, \quad (2.11)$$

$$\nabla \cdot \mathbf{v} = 0, \quad \text{in } \Omega, \quad (2.12)$$

$$\mathbf{v} \cdot \mathbf{n} = 0, \quad \text{on } \partial\Omega. \quad (2.13)$$

Consequently the Helmholtz decomposition of the vector field $\mathbf{u}^*$ can be achieved using either the primal formulation (1.1)–(1.3) or its equivalent dual one (2.11)–(2.13).

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