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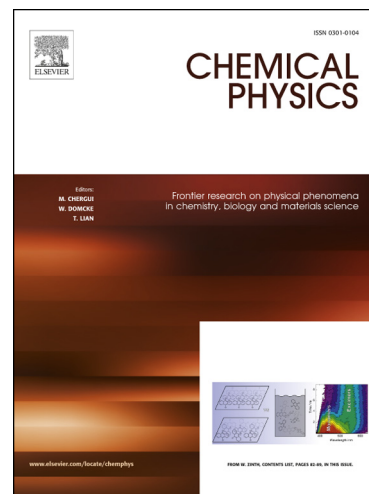
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Kinetics of Oxygen Species in an Electrically Driven Singlet Oxygen Generator

V. N. Azyazov, A. P. Torbin, A. A. Pershin, P. A. Mikheyev and M. C. Heaven

Abstract— The kinetics of oxygen species in the gaseous medium of a discharge singlet oxygen generator has been revisited. Vibrationally excited ozone $O_3(v)$ formed in $O+O_2$ recombination is thought to be a significant agent in the deactivation of singlet oxygen $O_2(a^1\Delta)$, oxygen atom removal and ozone formation. It is shown that the process $O_3(v \geq 2) + O_2(a^1\Delta) \rightarrow 2O_2 + O$ is the main $O_2(a^1\Delta)$ deactivation channel in the post-discharge zone. If no measures are taken to decrease the oxygen atom concentration, the contribution of this process to the overall $O_2(a^1\Delta)$ removal is significant, even in the discharge zone. A simplified model for the kinetics of vibrationally excited ozone is proposed. Calculations based on this model yield results that are in good agreement with the experimental data.

Index Terms— discharge driven oxygen iodine laser, singlet oxygen, ozone, vibration, quenching, recombination, $O(^3P)$, $O_2(a^1\Delta)$, O_3 .

I. INTRODUCTION

Gas phase processes involving the reactive oxygen species $O_2(a^1\Delta)$, $O(^3P)$ and O_3 have been studied extensively due to their importance in the terrestrial atmosphere [1-6], oxygen-containing plasmas [7-10], the oxygen-iodine laser [11-14] and combustion [15-16]. There has been a long standing interest in the possibility of developing a discharge driven oxygen iodine laser (DOIL), where $O_2(a^1\Delta)$ is produced via electric discharge [7-13]. In the DOIL devices investigated to date [12], the low small-signal gain of the active medium has prevented the efficient extraction the energy stored in singlet oxygen $O_2(a^1\Delta)$ [17]. To increase the gain, the $O_2(a^1\Delta)$ concentration must be increased [8]. Vasiljeva *et al.* [18] found that the $O_2(a^1\Delta)$ quenching rate in the discharge afterglow increases with oxygen pressure much faster than predictions based on the two-body quenching due to the components of an oxygen containing plasma (O , O_2 , O_3 etc.). They suggested that the faster deactivation of singlet oxygen at higher pressures was due to a three-body deactivation (3BD) process (reaction 1 in Table 1). With this assumption the rate of $O_2(a^1\Delta)$ deactivation increases with increasing concentrations of oxygen atoms $[O]$, molecular oxygen $[O_2]$ and the buffer gas $[M]$. To explain the singlet oxygen

dynamics in the discharge afterglow in the O_2/Ar mixture, Vasiljeva *et al.* [18] proposed a rate constant for reaction (1) that was $k_1^{O_2} = 1 \times 10^{-32}$ cm⁶/s for $M=O_2$ and $k_1^{Ar} = 0.63 \times 10^{-32}$ cm⁶/s for $M=Ar$.

The authors of Ref's [19, 20] attempted to measure the rate constant of reaction 1 using pulsed laser photolysis of O_2/O_3 mixtures at 248 nm to produce oxygen atoms and $O_2(a^1\Delta)$ molecules. They observed rapid quenching of $O_2(a^1\Delta)$ in $O(^3P)/O_2/O_3$ mixtures. It was found that the addition of Ar had no effect on the rate of $O_2(a^1\Delta)$ deactivation, whereas the addition of He and CO_2 decreased the quenching rate [19, 20], in contradiction with expectations based on the 3BD mechanism. Attempts to model the fast decay using the 3BD mechanism yielded a rate constant of $k_1^{O_2} = (1.1 \pm 0.1) \times 10^{-31}$ cm⁶/s, which was inconsistent with the results of the flowing afterglow experiments.

As an alternative to the 3BD mechanism, the authors of Ref. [21] proposed that vibrationally-excited ozone ($O_3(v)$) formed in the recombination process (Table 1, reaction 2) is responsible for the fast $O_2(a^1\Delta)$ quenching via reaction 3 (2BD mechanism). $v = v_1 + v_2 + v_3$ is the total number of excited vibrational quanta for the ozone, where the labels indicate the symmetric stretch (v_1), bending (v_2) and asymmetric stretch (v_3) modes. The reaction set that is used here to model the singlet oxygen kinetics is presented in Table 1. It is well-known [22-25] that $O_3(v)$ forms in process 2. Although the rate constant for quenching of $O_2(a^1\Delta)$ by vibrationally cold O_3 is small (4×10^{-15} cm³ s⁻¹), Kurylo *et al.* [26] found that the rate constant for quenching by O_3 that has one quantum of vibrational energy is faster by a factor of 38 ± 20 . Rawlins *et al.* [27] reported an estimated rate constant for quenching of $O_2(a^1\Delta)$ by ozone that has two or more quanta of the stretch modes to be in the range 10^{-11} - 10^{-10} cm³ s⁻¹. Recently [4, 19-21] additional experimental evidence for the critical role of vibrationally excited ozone in the quenching of $O_2(a^1\Delta)$ in $O(^3P)/O_2/O_3$ systems was reported.

There are two pathways that determine the fate of $O_3(v)$ formed by recombination process 2. Stabilization of $O_3(v)$ is achieved by process 4. Process 3 results in the dissociation of $O_3(v)$ and regeneration of oxygen atoms, $O(^3P)$. Two odd oxygen species, $O_3(v)$ and $O(^3P)$, are removed by reaction 5. Processes involving vibrationally-excited ozone may have a significant impact on the oxygen species kinetics in a DOIL. The objectives of the present study were to examine the influence of processes with $O_3(v)$ on the rates of $O_2(a^1\Delta)$ quenching, $O(^3P)$ removal and O_3 formation under conditions relevant to DOIL systems.

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