



# Lifetimes of long-lived states in inhomogeneous magnetic fields



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## ABSTRACT

Long-lived states (LLS), also known as singlet states, have been widely studied in the last decade. So far, LLS have only been observed in homogeneous magnetic fields, which preclude applications to many biological samples that are inherently inhomogeneous. We present a method to measure the lifetimes  $T_{LLS}$  of long-lived states in inhomogeneous magnetic fields, which combines established sequences for the excitation of LLS with their conversion into long-lived coherences (LLC) that can be detected by windowed acquisition. The method is applied to a pair of diastereotopic scalar-coupled protons of glycine in the dipeptide Alanine-Glycine (Ala-Gly).

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## 1. Introduction

If the static magnetic field is not perfectly homogeneous, NMR signals suffer from inhomogeneous broadening that makes it difficult to determine spectral parameters like chemical shifts, coupling constants, and the underlying homogeneous linewidths. It is possible to extract high resolution NMR information by refocusing inhomogeneous line broadening using some form of spin-echo spectroscopy [1], but the task of recovering the full spectral information often remains an elusive objective.

Recently, long-lived states (LLS) have attracted considerable attention because of their long spin memory [2]. If we consider two coupled protons, it is easy to prepare LLS by a simple sequence of pulses such as described by Sarkar et al. [3]. The population of the resulting LLS can then be 'sustained' using sequences of radio-frequency pulses with or without windows [4,5]. The difference in the decay rates of the populations of LLS and ordinary Zeeman states is due to the formers' immunity to the dominant intramolecular dipole-dipole relaxation mechanism. LLS can be utilized to study various slow dynamical processes that are normally difficult to characterize because of short longitudinal relaxation times [3]. LLS are not necessarily confined to pairs of coupled

protons, but can also be excited in multi-spin systems, as occur in many molecules of biological importance [6–10]. Furthermore, LLS signals can be enhanced by dissolution dynamic nuclear polarization (DNP) [11], thus allowing the study of substrates in very low concentrations [12]. Although applications of LLS are increasingly popular in various fields of chemistry and biology, a major challenge is their detection in inhomogeneous magnetic fields. Since LLS are not directly observable, they are usually converted into detectable single-quantum coherences that are sensitive to magnetic field inhomogeneities. Thus the application of LLS is challenging when the magnetic field homogeneity is poor as for *in-cellulo* or *in vivo* NMR.

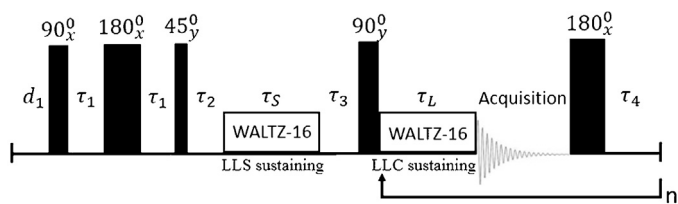
Zero-quantum coherences (ZQCs) between the singlet and central triplet states are known as long-lived coherences (LLCs) [13–15]. Their lifetimes  $T_{LLC}$  can be much longer than those of single quantum coherences (SQCs), but shorter than those of LLS, since LLCs are not immune to intramolecular dipole-dipole relaxation. We have shown that LLCs are immune to spatial and temporal variations of external static magnetic fields, so that they can provide spectra with very high resolution even in moderately inhomogeneous fields [16]. We have also shown that LLCs can be used to record high-resolution two-dimensional NMR spectra in inhomogeneous magnetic fields to determine scalar couplings and differences of chemical shifts [17]. However, studies in inhomogeneous magnetic fields have not been extended from LLCs to the realm of LLS so far.

The lifetime  $T_{LLS}$  can normally be determined after conversion into SQCs by monitoring the latter's signal intensities as a

Abbreviations: LLC, Long-lived coherences; LLS, Long-lived states; SID, Sustained induction decay; SQC, Single quantum coherence.

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**Figure 1.** Pulse sequence designed to determine  $T_{LLS}$  in inhomogeneous magnetic fields. Typically, the total number of loops may be  $n = 1024$ , and the recovery delay  $d_1$  may be 3–5 times  $T_1$ ,  $\tau_1 = 1/(4J)$ ,  $\tau_2 = 1/(2\Delta\nu_{IS})$ ,  $\tau_3 = 1/(4\Delta\nu_{IS})$ , where  $J$  is the scalar coupling constant and  $\Delta\nu_{IS}$  is the chemical shift difference of the two involved spins,  $\tau_S$  is a variable delay where the LLS are sustained,  $\tau_L$  is a fixed delay where the LLC are sustained (in our example 24 ms), the acquisition time and  $\tau_4$  have the same duration of 492.35  $\mu\text{s}$ . All 246 points acquired at intervals of 2  $\mu\text{s}$  within each acquisition block are averaged to make up a single point of the sustained induction decay (SID) [16]. The sustaining or spin locking pulses can be modulated in the manner of WALTZ-16 sequences.

function of the interval  $\tau_S$  where the LLS are sustained, and fitting their decay to an exponential function  $\exp\{-\tau_S/T_{LLS}\}$  [18,19]. Unfortunately, this method is limited to homogeneous magnetic fields. We demonstrate here the measurement of LLS lifetimes in inhomogeneous magnetic fields by converting the LLS for detection into LLC that are insensitive to inhomogeneities. This can be achieved by combining sequences for the excitation and sustaining of LLS with sequences for the detection of LLC, as shown in Figure 1.

## 2. Experimental

All experiments were performed at room temperature using an ultra-shielded 500 MHz Avance Bruker spectrometer equipped with an inverse 5 mm cryoprobe. A 0.5 M solution of Ala-Gly was prepared in 0.5 ml  $D_2O$ . The chemical shift difference of the two diastereotopic protons of Gly is  $\Delta\nu_{IS} = 57$  Hz and the scalar coupling constant is  $J = -17.24$  Hz. The longitudinal relaxation time of these protons is  $T_1 = 1.2$  s and the transverse relaxation time of their SQCs is  $T_2 \approx 0.9$  s. The LLC sustaining time  $\tau_L$  and the acquisition time (which must be equal to the refocusing interval  $\tau_4$ ) depend on the spectral width in the  $\omega_2$  domain. For the detection of LLC's in inhomogeneous fields, we chose a spectral width of 40 Hz, hence a duration of the detection loop  $\tau_{loop} = 1/40$  Hz = 25 ms (this plays the role of a dwell time for the detection of LLC's), a spin-locking interval where the LLC are sustained by rf irradiation  $\tau_L = 24$  ms, and a refocusing pulse length  $\tau_\pi = 15.3$   $\mu\text{s}$ . The acquisition time and the refocusing interval must have the same duration and obey the relation  $\tau_4 = (1/2)(\tau_{loop} - \tau_L - \tau_\pi) = 492.35$   $\mu\text{s}$ . The amplitude of the WALTZ-16 irradiation applied in both intervals  $\tau_S$  and  $\tau_L$  where the LLS and LLC are sustained was 1 kHz. Each LLC spectrum was acquired with  $n = 1024$  data points. A total of 13 experiments were carried out while incrementing the sustaining interval  $\tau_S$  to monitor the decays of the integrals of the LLC signals which were fitted to an exponential function  $\exp\{-\tau_S/T_{LLS}\}$ .

## 3. Methods, results and discussion

The pulse sequence consists of three parts (Figure 1): (i) broadband excitation and sustaining of LLS as described by Sarkar et al. [3]; (ii) conversion of LLS into LLC, and (iii) acquisition of signals derived from LLS as shown by Bornet et al. [16]. As seen from the Figure 1, starting from equilibrium magnetization ( $I_z + S_z$ ), the initial  $90_x$  pulse is followed by a spin echo sequence to create terms ( $2I_xS_z + 2I_zS_x$ ) due to the evolution under the scalar coupling  $J_{IS}$ , regardless of the homogeneity of the static magnetic field. At the top of the echo, a  $45_y$  pulse converts ( $2I_xS_z + 2I_zS_x$ ) into ( $2I_xS_x - 2I_zS_z$ )  $\equiv$  ( $ZQ_x + DQ_x - 2I_zS_z$ ), i.e., into a superposition of

multiple-quantum coherences and  $zz$  order. The ( $ZQ_x - 2I_zS_z$ ) terms are converted into ( $-ZQ_x - 2I_zS_z$ ) under free evolution in the interval  $\tau_2 = 1/(2\Delta\nu_{IS})$ , since both are immune to the inhomogeneity of the static field, whereas  $DQ_x$  is dephased. The terms  $ZQ_x$  and  $2I_zS_z$  are then sustained by a WALTZ-16 sequence for a variable time  $\tau_S$ . This rf irradiation has the effect of suppressing the chemical shifts to that the spins  $I$  and  $S$  become magnetically equivalent. In the process, the ( $-ZQ_x - 2I_zS_z$ ) terms are transformed into a long-lived state (LLS), also known as “triplet-singlet imbalance” or “T/S imbalance”, in analogy to “A/E imbalance” in methyl groups [20]. This amounts to a non-vanishing difference between the average population of the three triplet states that are symmetric under spin permutation and the population of the antisymmetric singlet state:

$$\overline{|T\rangle\langle T|} - \overline{|S\rangle\langle S|} = \frac{4}{3} \langle I_1 \cdot I_2 \rangle \neq 0$$

$$\text{where } \overline{|T\rangle\langle T|} = 1/3 [ |T_{+1}\rangle\langle T_{+1}| + |T_0\rangle\langle T_0| + |T_{-1}\rangle\langle T_{-1}| ]$$

$$\text{and } \overline{|S\rangle\langle S|} = [ |S_0\rangle\langle S_0| ].$$

After the sustaining interval  $\tau_S$ , the rf field is switched off, so that the LLS is converted back into ( $-ZQ_x - 2I_zS_z$ ) terms. The delay  $\tau_3 = 1/(2\Delta\nu_{IS})$  then allows the conversion of ( $-2I_xS_x - 2I_yS_y - 2I_zS_z$ )  $\equiv$  ( $-ZQ_x - 2I_zS_z$ ) into ( $2I_xS_y - 2I_yS_x - 2I_zS_z$ )  $\equiv$  ( $-ZQ_y - 2I_zS_z$ ). Again, neither  $ZQ_y$  nor  $zz$ -order are affected by the magnetic field inhomogeneity. Now a  $90_y$  pulse transforms ( $-ZQ_y - 2I_zS_z$ ) into ( $-2I_yS_z + 2I_zS_y - 2I_zS_z$ ), that is, into anti-phase SQC terms in addition to unobservable  $zz$ -order. These anti-phase terms, if one attempted to detect them in an inhomogeneous static field, would be obscured by the linewidth. In order to alleviate this problem, the signal has to be acquired by some procedure that is not influenced by static field inhomogeneities. One such method is to convert the order from LLS to LLC, and then to detect the signal using an “on the fly” or “windowed” LLC acquisition scheme which is immune to static field inhomogeneities.

In the product operator formalism, an LLC can be represented by terms ( $I_x - S_x$ ) or ( $2I_yS_z - 2I_zS_y$ ) which slowly decay when they are spin-locked [15]. The state ( $-2I_yS_z + 2I_zS_y - 2I_zS_z$ ) remaining after the  $90_y$  pulse contains anti-phase SQC terms which can be observed in the form of narrow spectral lines if acquired by spin locking alternating with windowed acquisition. In the observation windows, there is a partial conversion under the effect of the scalar coupling  $J_{IS}$  so that the Fourier transformation of the signal detected in the  $t_2$  interval gives rise to a doublet at  $+2\pi J_{IS}$  and  $-2\pi J_{IS}$  in the  $\omega_2$  dimension of which only one peak is shown here.

In order to measure the relaxation time  $T_{LLS}$ , the duration of the sustaining interval  $\tau_S$  has to be varied while detecting the corresponding LLC signals (Figure 3). In this way we can simultaneously measure LLS as well as LLC lifetimes. The  $T_{LLS}$  value can be obtained by fitting the decaying signals to an exponential function  $\exp\{-\tau_S/T_{LLS}\}$  (Figure 4). The LLC lifetime is obtained from  $T_{LLC} = (\pi\Delta\nu_{LLC})^{-1}$  where  $\Delta\nu_{LLC}$  is the linewidth at half height of the LLC signal in the  $\omega_2$  dimension, expressed in Hz.

In high resolution solution-state NMR, the required magnetic field homogeneity must typically be on the order of  $10^{-9}$  or 1 ppb [21]. In our experiments, we deliberately made the magnetic field inhomogeneous by missetting various shim currents ( $z_1, z_2, z_3, x, y, z_0x, z_0y$ ) to broaden the linewidth as shown in Figure 2 for a homogeneous field with a linewidth  $\Delta\nu = 1.8$  Hz (red line) and for an inhomogeneous field with a linewidth  $\Delta\nu = 17$  Hz (blue line). Such changes in the homogeneity of the static magnetic field have little effect on the linewidths of LLCs [16].

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