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# Quantum effects of hydrogen nuclei on a structure and a dynamical rearrangement of hydrogen-bond networks

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## ABSTRACT

To understand anomalous energetics and dynamics in water, nuclear quantum effects such as zero-point energy and delocalization of wave packets (WPs) representing water hydrogen atoms should be essential. However, since quantum calculations of many-body dynamics are highly complicated in general, none has yet directly viewed the quantum WP dynamics of hydrogen atoms in liquid water. Our semiquantum molecular dynamics simulation for the first time made it possible to observe the hydrogen WP dynamics in liquid water. Here, we demonstrate that the microscopic WP dynamics are closely correlated to and play key roles in the dynamical rearrangements in the hydrogen-bond network (HBN) of liquid water. Especially, the memory loss of the dipolar angle formed by two waters is accelerated with the nuclear quantum effect introduced by the nuclear WP, rationalizing the faster memory decay of the HBN in the quantum liquid water. We found the quantum effects of hydrogen atoms on mesoscopic liquid water behaviors such as intermittent collective motions associated with the rearrangement of HBN and the concomitant fluctuation and relaxation. Our results will provide new physical insights on HBN dynamics in water whose significance is not limited to the pure liquid dynamics but also for understanding chemical and biological reactions in liquid water.

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## 1. Introduction

Chemical reactions taking place in water are essential for many important processes in electrochemistry, environmental, pharmaceutical, and biological sciences. It is recognized that the unique and anomalous properties of water stem from its dynamics of hydrogen-bond network (HBN). Many attempts have been conducted to clarify how liquid water is dynamically organized via the HBN. Computer simulation has been a strategy of particular importance in this field. A number of empirical potentials have been developed toward better agreement with experiments, with the parameters calibrated for use in full classical water (FCW) molecular dynamics (MD) and Monte Carlo simulations [1–4]. The classical MD is nowadays very straightforward, and is routinely used to study complex chemical problems involving many thousands of particles [3]. In contrast, quantum mechanical simulations are yet unestablished for many-body dynamics in condensed phases, despite that the nuclear quantum effects such as zero-point energy (ZPE), delocalization of WPs describing water hydrogen atoms, etc. are envisaged to be essential in the HBN dynamics. Although considerable progress has been made in the development of exact quantum-mechanical methods, they are yet applicable to relatively small systems and short times, and only the static spatial and momentum distributions were obtained. In order to capture such quantum effects in many degrees of freedom, a number of approaches exploiting semiquantum and mixed quantum–classical simulations have been proposed and applied in recent years. The path-integral

Monte Carlo (PIMC) method was used to assess the effect of quantum-mechanical fluctuations on a liquid structure [5,6], and the centroid molecular dynamics (CMD) method was proposed to study water dynamics through the equilibrium correlation functions and mean square displacement [7]. However, none has yet directly viewed quantum dynamics of water hydrogen atoms in condensed bulk water and their effects on mesoscopic behaviors such as intermittent collective motions associated with a rearrangement of HBNs and the concomitant fluctuation and relaxation [1,8]. We for the first time report a direct simulation study of hydrogen WP dynamics in liquid water, and suggest its key roles in a mesoscopic dynamical rearrangement of HBNs.

## 2. Simulation method

In order to take account of the nuclear quantum aspects such as the WP delocalizing and the ZPE effects in the liquid water dynamics, the hydrogen atoms are described by  $N$ -independent symmetric 3D squeezed coherent state Gaussian WP basis functions,  $\Psi(\{x_i, y_i, z_i, t\}) = \prod_{i=1}^N \Psi_i(x_i, y_i, z_i, t)$ , where

$$\Psi_i(x_i, y_i, z_i, t) = N_i \exp \left[ A_i(t) \left\{ (x_i - q_{x_i}(t))^2 + (y_i - q_{y_i}(t))^2 + (z_i - q_{z_i}(t))^2 \right\} + ip_{x_i}(t) \{x_i - q_{x_i}(t)\} + ip_{y_i}(t) \{y_i - q_{y_i}(t)\} + ip_{z_i}(t) \{z_i - q_{z_i}(t)\} \right]. \quad (1)$$

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$q_{\alpha_i}(t)$  and  $p_{\alpha_i}(t)$  with  $\alpha_i = x_i, y_i$  or  $z_i$  are the 3D Cartesian coordinates of the WP center and their momenta, respectively. (For simplicity,  $\hbar = 1$  and all the coordinates are mass-scaled.)  $A_i$  is defined by  $A_i(t) = (-1 + 2i\rho_i(t)\Lambda_i(t))/4\rho_i^2(t)$ , where  $\rho_i(t)$  and  $\Lambda_i(t)$  correspond to the WP width and its momentum, respectively.  $N_i = (2\pi\rho_i^2(t))^{-3/4}$  is the normalization factor. The WP function (1) is specified by the time-dependent parameters  $\{q_{\alpha_i}(t), p_{\alpha_i}(t), \rho_i(t), \Lambda_i(t)\}$  to be determined through the time-dependent variational principle of the action integral [9–14]. The resulting equations of motions (EOM) which we have to solve finally become

$$\dot{q}_{\alpha_i} = \frac{p_{\alpha_i}}{m_i}, \quad \dot{p}_{\alpha_i} = -\frac{\partial \sum_{i<j}^N \langle V_{ij}(x_{ij}, y_{ij}, z_{ij}) \rangle}{\partial q_{\alpha_i}} \quad (2)$$

$$\dot{\rho}_i = \frac{\Lambda_i}{m_i} \lambda_i = -\frac{\partial \sum_{i<j}^N \langle V_{ij}(x_{ij}, y_{ij}, z_{ij}) \rangle}{3\partial \rho_i} + \frac{\hbar^2}{4m_i\rho_i^3}, \quad (3)$$

for  $i = 1, 2, \dots, N$ . Here, the intra- and intermolecular potential functions  $V_{ij}(x_{ij}, y_{ij}, z_{ij})$  are averaged as  $\langle X \rangle \equiv \langle \Psi, t | X | \Psi, t \rangle$  (and the mass  $m_i$  and  $\hbar$  have been retrieved). Our semiquantum water (SQW) program requires only  $2N$  additional EOM for the auxiliary coordinates and momenta representing the WP widths of hydrogen atoms. The 3D SQW MD simulation was developed by expanding the potentials  $V_{ij}(x_{ij}, y_{ij}, z_{ij})$  around  $\langle x_i \rangle$ ,  $\langle y_i \rangle$  and  $\langle z_i \rangle$  for hydrogen atoms [9,10]. The dynamics of oxygen atoms are also influenced by the additional non-classical forces from the semiquantum hydrogen atoms. With the explicit extended potential functions, we can perform the SQW dynamics simulations of 1024 water molecules solving the EOM (2) and (3). The extra computational costs for SQW over full classical water (FCW) are only marginal, about 20%. Nevertheless, our SQW simulation reproduced the major properties of the PIMC and CMD methods as was reported in our previous paper [9]. Our SQW method could be especially effective in biosimulations where a large number of solvating water molecules are often needed to realize stable simulations.

### 3. Results and discussions

Fig. 1 shows power spectra of kinetic and potential energies of WPs describing the semiquantum hydrogen atoms (the left two panels). The power spectrum is defined as the absolute square of Fourier transformed energy dynamics,  $|\int E(t) \exp(i\omega t) dt|^2$ . The power spectra of the low frequency region are both well-fitted by  $1/f^2$ , while the higher frequency spectra exhibit the complicated behavior. This is quite different from the classical parts of kinetic, potential and intramolecular energies in the same semiquantum water (SQW) simulation, where the entire spectra can be fitted by single functions  $1/f^\alpha$  with the exponents  $0.5 < \alpha < 1.5$  (the right three panels). The power spectra demonstrate that, although the classical parts exhibit multi-timescale dynamics as was reported by previous works [8], the low frequency energy dynamics of WPs are rather of single-timescale and not strongly influenced by the classical energy fluctuations.

To study dynamics of the structure, we define the time-dependent radial distribution functions (RDFs) as

$$g(r, t) = \frac{n(r, t)}{4\pi r^2 dr n_0}, \quad (4)$$

where  $n(r, t)$  represents the number of oxygen atoms between shells of radii  $r$  and  $r + dr$  from a representative oxygen atom (to be called  $i$ th oxygen atom), and  $n_0$  is the oxygen number density of the whole system. The RDF dynamics directly reflects the liquid structure rearrangement around the oxygen atom. Fig. 2 shows the time-dependent maximum peak positions of the RDFs for pairs of the  $i$ th oxygen atom and the surrounding other 1023 oxygen atoms. The panels are different in the

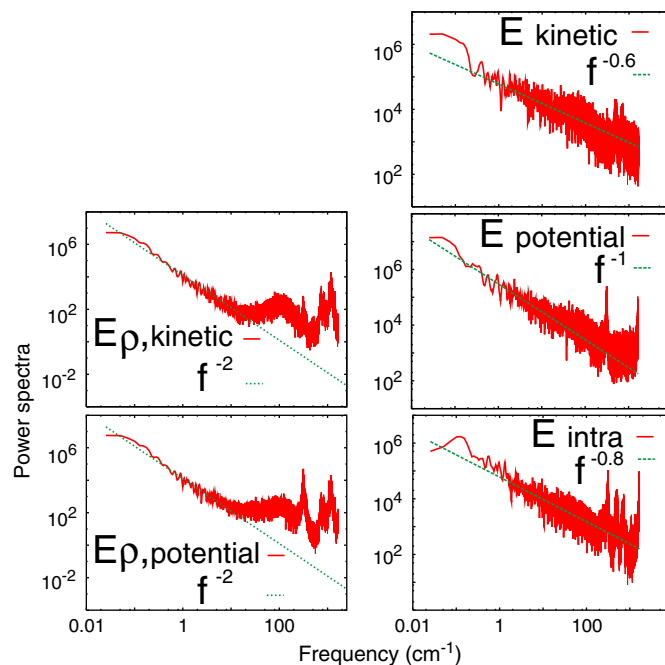


Fig. 1. Power spectra of the kinetic and potential energies of WPs describing semiquantum hydrogen atoms, and the classical parts of kinetic, potential and intramolecular energies. The power spectra of the WP kinetic and potential energies show a simple decay  $1/f^2$  in the lower frequency region, whereas those of the classical energy parts are all reproduced by the functions  $1/f^\alpha$  with the exponents  $0.5 < \alpha < 1.5$ . Although the classical energy dynamics follow the multi-timescale dynamics, the WP energy dynamics are not strongly coupled to them and exhibit the single-timescale dynamics.

time-resolution for averaging as indicated beside each panel. With the finer time-resolution, the maximum peak positions for SQW and FCW both fluctuate randomly reflecting the short-time thermal modulations. When the time-resolution is coarse-grained, the random thermal fluctuations are smeared out and the time-dependent maximum peak position exhibits a behavior of transitions between well-structured and less-structured states. This implies that the water structure in the meso-timescale have two typical configurations at fully coordinated and distorted local HBNs. It supports the mixture model recently reported by X-ray experiments [15–17]. We also mention that the distorted structure for SQW is less structured than that for FCW, reflecting the smoothed free energy landscape due to the ZPE of the semiquantum hydrogen atoms [9–12].

Fig. 3 demonstrates that the hydrogen quantum effects influence the hydrogen bond dynamics. We calculated the number of water molecules hydrogen-bonded to the  $i$ th water molecule as a function of time,  $N_{\text{HB}}(t)$ . The common geometrical criterion for hydrogen bonds was used [3], i.e., if the distance between two oxygen atoms,  $r_{\text{OO}}(t)$ , is less than 3.5 Å and the OH–O angle is greater than 150°, the water pair is considered as hydrogen bonded. The evolution of  $N_{\text{HB}}(t)$  directly reflects the HBN dynamics [1]. The upper panel of Fig. 3 shows the normalized time correlation functions (TCFs)  $\langle N_{\text{HB}}(t)N_{\text{HB}}(0) \rangle / \langle N_{\text{HB}}^2(0) \rangle$ . The TCF for SQW decreases faster than the TCF for FCW, indicating that the quantum effect of hydrogen atoms accelerates the memory decay of the HBN. We further show the TCF of the angular coordinate  $\phi(t)$  which is defined as an angle formed by dipole vectors of two water molecules (dipolar angle). The angular TCF exhibits the similar tendency: The memory of the dipolar angle decays faster with the nuclear quantum effects than without the nuclear quantization, rationalizing the faster memory decay of the HBN in the quantum liquid water. These accelerated memory decays are in accord with the experimentally observed ultrafast memory loss in liquid water [18]. Fig. 3 also shows that the kink-like ultrafast (2 fs) dynamics appears both in the HBN and dipolar angle dynamics; the disappearance of the HBN kink in the SQW case partially depends on the definition of a hydrogen bond since the

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