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# Fourier transform spectral study of $B^2\Sigma^+ - X^2\Sigma^+$ system of AlO

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#### Abstract

The spectrum of  $B^2\Sigma^+-X^2\Sigma^+$  system of AlO has been recorded on BOMEM DA8 Fourier transform spectrometer at an apodized resolution of 0.05 cm<sup>-1</sup>. Nineteen bands of the  $\Delta v = 1$ , 0, -1, and -2 sequences of this band system have been analyzed for the rotational structure. Out of which seven bands, viz. 3–2, 4–3, 2–3, 3–4, 4–5, 5–6 and 6–7 have been analyzed for the first time. The rotational lines of these 19 bands along with 20 earlier analyzed bands, a total of 7200 lines, have been fitted in a simultaneous least squares fit. The study has resulted in determining more precise vibrational and rotational constants of the two states. Because of the high resolution employed it became necessary to invoke  $H_0$  and  $H_1$  coefficients, and a fifth order term to explain the anomalous spin-doubling observed in the v'' = 5, 6 and 7 levels of the  $X^2\Sigma^+$  state. © 2007 Elsevier Inc. All rights reserved.

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#### 1. Introduction

The blue-green system,  $B^2\Sigma^+ - X^2\Sigma^+$  transition, of AlO has been known for over 100 years and extensively studied both experimentally and theoretically [1–11,17–19,23,24]. The astrophysical importance of AlO and its aeronomical significance is well known [12–16,21,23]. The existence of AlO molecule in cool stellar atmospheres has been shown through the identification of B-X system in spectrum of some normal Mira giants and Mira variables [12,16]. In stars of lower temperature the absorption of the head of 0-0 band  $(\lambda 4842 \text{ Å})$  of the B-X system strengthens. In normal Mira giants because of higher temperature the band is seen in emission. In Mira variables the AlO absorption is greater than in normal giants and the maximum intensities far exceed those obtained in the normal M giants. The strengthening of AlO ( $\lambda$  4842 Å) intensity, in these stars is quite variable from cycle to cycle. The possibility of some sort of

interference could not be ruled out. It was suggested to look for some other system of AlO, such as  $A^2\Pi_i - X^2\Sigma^+$  system. Temperatures and densities in the upper atmosphere, as well as differential excitation of solar radiation have been inferred from observations on B-X system of AlO, following release of organoaluminium compounds [13–15].

Over the years several workers have tried to give a vibrational expression involving the observed bands of the B–X system [1–11]. Roy [5] tried to give an improved vibrational expression compared to Mecke [1]. But as it has been recently confirmed that, since some of the band heads of the system are shifted due to interaction of  $A^2\Pi_i$  state with the ground  $X^2\Sigma^+$  state [28], it would never have been possible to find a perfect expression with minimum obs.–cal. differences. A sudden drop of intensity in the bands 4–7, 6–7, 5–7, 10–7 and those involving  $v'' \ge 5$  and 8 was noticed by Rosen [7]. Making use of the extrapolated  $B_v$  values he concluded that in the vibrational levels  $v'' \ge 6$ , 7 and 8 predissociation sets in at N = 60, 44 and 18, respectively. He also concluded that all the bands of  $v'' \ge 9$  progressions were predissociated. On the basis of this

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observation Rosen proposed a low dissociation energy for the ground state. Lagerqvist et al. [9] observed some more bands of the same transition involving v'' = 5. He stated that the predissociations observed by Rosen might have been illusory; which later turned out to be really true. In the bands with  $v'' \ge 6$  the rotational lines go off because of perturbations and not because of predissociations. Shimauchi [10] photographed the spectra in air, nitrogen, oxygen and argon using high-grade aluminium rods as electrodes. She was able to record the bands with v' = 16and v'' = 12 and for the first time indicated that v'' = 9 level of  $X^2\Sigma^+$  state is raised by 10 cm<sup>-1</sup> on account of perturbation by certain state. Shimauchi's vibrational analysis was most extensive. Goodlett and Innes [11], Mahieu et al. [20], Singh et al. [25], Coxon and Naxakis [27] and Saksena et al. [28] tried to determine the sign and magnitude of  $\gamma_0''$ on some consideration or other but it was Yamada et al. [29] who gave most accurate value of  $\gamma_0''$  using microwave data. An attempt was made by Bernard and Gravina [23] to calculate the rotational line positions of six bands of  $A^2\Pi_i - X^2\Sigma^+$  transition. These authors have taken the spin-doubling constants of v = 0, 1 and 2 levels of the  $\dot{X}^2\Sigma^+$  states as  $+0.00073 \text{ cm}^{-1}$ ,  $-0.00022 \text{ cm}^{-1}$  and -0.00134 cm<sup>-1</sup>, respectively (see Table 1, Ref. [23]). How these values were arrived at is not very clear since it has already been known that the spin-doubling constant  $\gamma_n^{\prime\prime}$ increases with v and has a positive sign [25,26,28].

In the light of the above we decided to record the high resolution spectra of AlO molecule using F.T. spectrometer to understand the behaviour of spin-doubling constant,  $\gamma_v''$  of the ground  $X^2\Sigma^+$  state. In turn, the study also helped in understanding the reasons due to which, so far, it was not possible to find an unique vibrational expression for the  $B^2\Sigma^+-X^2\Sigma^+$  system of AlO by all the previous workers.

### 2. Experimental

The spectrum of AlO was excited in a microwave discharge (2450 MHz, 180 W) through a flowing mixture of AlCl<sub>3</sub> vapours, buffer gas argon and a small amount of oxygen. The AlCl<sub>3</sub> kept in a small quartz boat of side arm of a discharge tube was heated by an electrical furnace and vapours, along with argon and oxygen, were let-in the discharge zone. The gas pressures were so optimized as to give very intense characteristic glow of AlO [22]. A spherical lens was used to focus the emission signal into the interferometer. The spectra in the region 18000-22000 cm<sup>-1</sup> were recorded with BOMEM DA8 Fourier transform spectrometer with an apodized resolution of 0.05 cm<sup>-1</sup> using quartz UV beam splitter and silicon detector. The emission signal being strong no filter was required. Fifty scans (integration time  $\sim$ 75 min.) were coadded to obtain an improved signal-to-noise ratio. The estimated rotational temperature from the rotational structure was ~1700 K. The accuracy of the measurements is expected to be better than  $0.005 \,\mathrm{cm}^{-1}$ .

Table 1 Head measurements of the bands of  $B^2\Sigma^+ - X^2\Sigma^+$  transition (in cm<sup>-1</sup>)

| Head measurements of the bands of $B^2\Sigma^+ - X^2\Sigma^+$ transition (in cm <sup>-1</sup> ) |            |           |                        |                  |                        |                  |            |
|---|------------|-----------|------------------------|------------------|------------------------|------------------|------------|
| Sr. no.   | $\Delta v$ | v'– $v''$ | $R_2$ head             | $N^{\mathrm{b}}$ | R <sub>1</sub> head    | $N^{\mathrm{b}}$ | $\Delta v$ |
| 1   | 4          | 5-1       | 23919.005 <sup>a</sup> | 10               | 23918.846 <sup>a</sup> | 10               | 0.159      |
| 2   | 4          | 6–2       | 23794.177 <sup>a</sup> | 11               | 23793.977 <sup>a</sup> | 11               | 0.200      |
| 3   | 4          | 7–3       | 23676.038 <sup>a</sup> | 11               | 23675.806 <sup>a</sup> | 11               | 0.232      |
| 4   | 4          | 8-4       | 23564.653 <sup>a</sup> | 11               | 23564.369 <sup>a</sup> | 11               | 0.284      |
| 5   | 4          | 9-5       | 23460.083 <sup>a</sup> | 12               | 23459.717 <sup>a</sup> | 11               | 0.366      |
| 6   | 4          | 10-6      | 23362.057 <sup>a</sup> | 12               | 23361.600 <sup>a</sup> | 12               | 0.457      |
| 7   | 4          | 11 - 7    | 23270.735 <sup>a</sup> | 12               | 23270.145 <sup>a</sup> | 12               | 0.590      |
| 8   | 3          | 3-0       | 23210.161 <sup>a</sup> | 11               | 23210.004 <sup>a</sup> | 11               | 0.157      |
| 9   | 3          | 4–1       | 23085.894 <sup>a</sup> | 12               | 23085.717 <sup>a</sup> | 12               | 0.177      |
| 10  | 3          | 5–2       | 22968.421 <sup>a</sup> | 12               | 22968.207 <sup>a</sup> | 12               | 0.214      |
| 11  | 3          | 6–3       | 22857.673 <sup>a</sup> | 12               | 22857.420 <sup>a</sup> | 12               | 0.253      |
| 12  | 3          | 7–4       | 22753.627 <sup>a</sup> | 13               | 22753.321 <sup>a</sup> | 12               | 0.306      |
| 13  | 3          | 8-5       | 22656.430 <sup>a</sup> | 12               | 22656.045 <sup>a</sup> | 12               | 0.385      |
| 14  | 3          | 9–6       | 22565.902 <sup>a</sup> | 13               | 22565.402 <sup>a</sup> | 13               | 0.500      |
| 15  | 3          | 10 - 7    | 22482.120 <sup>a</sup> | 14               | 22481.439 <sup>a</sup> | 14               | 0.681      |
| 16  | 2          | 2-0       | 22362.536 <sup>a</sup> | 13               | 22362.365 <sup>a</sup> | 12               | 0.171      |
| 17  | 2          | 3-1       | 22245.623              | 13               | 22245.421              | 13               | 0.202      |
| 18  | 2          | 4–2       | 22135.484              | 13               | 22135.253              | 13               | 0.231      |
| 19  | 2          | 5-3       | 22032.095              | 14               | 22031.819              | 13               | 0.276      |
| 20  | 2          | 6-4       | 21935.480              | 14               | 21935.127              | 14               | 0.353      |
| 21  | 1          | 1-0       | 21507.742              | 14               | 21507.546              | 14               | 0.196      |
| 22  | 1          | 2-1       | 21398.166              | 14               | 21397.975              | 14               | 0.191      |
| 23  | 1          | 3-2       | 21295.440              | 15               | 21295.176              | 15               | 0.264      |
| 24  | 1          | 4-3       | 21199.368              | 15               | 21199.006              | 15               | 0.362      |
| 25  | 0          | 0-0       | 20645.852              | 16               | 20645.585              | 16               | 0.267      |
| 26  | 0          | 1 - 1     | 20543.706              | 16               | 20543.500              | 16               | 0.206      |
| 27  | -1         | 0 - 1     | 19682.250              | 19               | 19682.021              | 19               | 0.229      |
| 28  | -1         | 1-2       | 19594.361              | 20               | 19594.018              | 19               | 0.343      |
| 29  | -1         | 2-3       | 19513.194              | 21               | 19512.785              | 20               | 0.409      |
| 30  | -1         | 3-4       | 19438.842              | 22               | 19438.318              | 21               | 0.524      |
| 31  | -1         | 4–5       | 19371.420              | 23               | 19370.738              | 22               | 0.682      |
| 32  | -1         | 5–6       | 19310.854              | 25               | 19309.900              | 24               | 0.954      |
| 33  | -1         | 6–7       | 19257.550              | 33               | 19256.224              | 31               | 1.326      |
| 34  | -2         | 0-2       | 18733.664              | 24               | 18733.287              | 23               | 0.377      |
| 35  | -2         | 1-3       | 18660.020              | 25               | 18659.530              | 24               | 0.490      |
| 36  | -2         | 2-4       | 18593.223              | 26               | 18592.635              | 26               | 0.588      |
| 37  | -2         | 3-5       | 18533.528              | 28               | 18532.655              | 27               | 0.873      |
| 38  | -2         | 4–6       | 18480.752              | 31               | 18479.588              | 30               | 1.164      |
| 39  | -2         | 5–7       | 18435.915              | 37               | 18434.105              | 34               | 1.810      |

*Note:* The accuracy of the  $R_2/R_1$  heads is  $\sim \pm 0.03$  cm<sup>-1</sup>.

### 3. Discussion

In the case of B–X band system of AlO, the efforts to obtain unique expression using head measurements could not succeed earlier because of two reasons. One, as the origin to vertex separation [32] differs vastly from bands of positive to negative sequences, resulting in large obs.–cal. differences and secondly, due to rapid increase in the spin-doubling constant  $\gamma_v''$ , in the  $v'' \ge 4$ –7 the separation of the two heads  $R_2$  and  $R_1$ , involving these levels also increases very rapidly (see Table 1). The band head positions shown in Table 1 are either the observed band heads in the F.T. spectrum recorded in the present study or the interpolated values from the least squares fit of rotational lines of all the bands.

Being a  ${}^2\Sigma^{-2}\Sigma$  transition the  $\gamma'_v$  and  $\gamma''_v$  are highly corelated and cannot be estimated independently using the

<sup>&</sup>lt;sup>a</sup> These are calculated positions.

<sup>&</sup>lt;sup>b</sup> Rotational number N, where head is formed.

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