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New atomic data for Ge XX

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ABSTRACT

We have performed large-scale configuration interaction (CI) calculations using CIV3for the lowest (in energy) 155 fine-structure levels of aluminum-like germanium ion. We have calculated the energies and lifetimes of these levels, together with oscillator strengths, and transition probabilities for the electric-dipole (optically allowed and intercombination) transitions among the levels of ground state 3s²3p (²P) and higher energy levels of states 3s³3p², 3s²3d, 3p³, 3s³93d, 3p²3d, 3s³d², 3d³, 3s² (4s, 4p, 4d, 4f) of Ge XX in the LSJ coupling scheme. The present results include relativistic effects through the Breit–Pauli operator. Improvements to the *ab initio* calculations are made by applying a "fine-tuning" technique in which small adjustments are made to some diagonal elements of the Hamiltonian matrix so as to bring the calculated energy separations into as close agreement as possible with experimental values. Comparisons are made with other available experimental and theoretical results and the accuracy of the present results is assessed.

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1. Introduction

The spectra of Al-like ions with Z > 30 have received a great deal of attention both experimentally and theoretically. In previous experimental works, which mainly make use of high-energy lasers or tokomak discharges, the spectra of some Al-like ions in plasmas have been measured, though for the heavier ions such as germanium, selenium, molybdenum and silver only the spectra from the transitions between $3s^2p^2$ and $3s^23d$ configurations have been identified by Hinnov et al. [1] and by Sugar et al. [2] for copper through molybdenum. Transitions in Ge XX between the $3s^23p$, $3s^23p$, $3s^23d$, $3p^3$ and $3s^2p3d$ configurations were identified in the extreme-ultraviolet spectra emitted from linear plasmas, produced using line-focused laser beams [3].

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http://dx.doi.org/10.1016/j.jqsrt.2015.07.023 0022-4073/© 2015 Elsevier Ltd. All rights reserved. Large-scale theoretical studies of Al-like ions have been performed using different methods: systematic studies of oscillator strengths for fine-structure transitions in the aluminum isoelectronic sequence [4]; relativistic manybody calculations of electric-dipole lifetimes, transitions rates and oscillator strengths for n=3 states in Al-like ions [5].

Critical data compilations based on available theoretical and experimental sources are given in National Institute of Standards and Technology (NIST), Atomic Spectra Database (ASD) [6].

In this work, we have chosen to use the CIV3 program [7,8] to calculate energy levels and oscillator strengths. In a wide range of previous calculations this code has demonstrated its ability to obtain accurate results, especially when strong mixing occurs. The present study involves extensive calculations and gives results for optically allowed and intercombination electric-dipole transitions amongst the lowest 155 energy levels of the Ge XX ion by including configuration interaction (CI) to represent valence shell correlation effects.

2. Theoretical method

The wavefunctions describing the atomic states included in these calculations were obtained using configuration interaction code CIV3 of Hibbert [7,8]. The CI

| Table 1 | | | | | | |
|-----------------|--------------|--------|----------|----------|------------|--------|
| Radial function | parameters f | for of | otimized | orbitals | of Al-like | Ge XX. |

| Orbitals (<i>nl</i>) | Expansion coefficients (<i>C_{jnl}</i>) | Power of <i>r</i> (<i>I_{jnl}</i>) | Exponents (ξ _{jnl}) |
|---------------------------|---------------------------------------------------|-------------------------------------------------|----------------------------------|
| 3d | 0.0768888 | 3 | 13.6044828 |
| | 0.9416767 | 3 | 7.2784949 |
| | 0.0033590 | 3 | 3.3769297 |
| 4s | - 1.7810545 | 4 | 5.2772495 |
| | 0.0071013 | 4 | 3.2983472 |
| | 2.2740362 | 3 | 7.3552125 |
| | 0.0112894 | 3 | 2.7572943 |
| | - 1.1918070 | 2 | 8.2033091 |
| | 0.1380578 | 1 | 23.9226036 |
| 4p | 3.4588210 | 4 | 5.4838356 |
| | 0.1294863 | 4 | 13.6449050 |
| | -3.2431628 | 3 | 5.2874429 |
| | 0.1214499 | 3 | 3.9884308 |
| | 0.3424522 | 2 | 13.9180855 |
| 4d | -1.3406807 | 4 | 4.6181499 |
| | 0.1747193 | 4 | 3.1520884 |
| | -0.0311369 | 3 | 1.7274663 |
| | 0.7434741 | 3 | 7.8346912 |
| 4f | 0.9417413 | 4 | 5.2748413 |
| | 0.0677806 | 4 | 3.6034363 |
| 5s | 5.7226171 | 1 | 2.7557600 |
| | - 17.0442927 | 2 | 5.0267621 |
| | 8.2829726 | 3 | 3.4616527 |
| | 8.3735165 | 4 | 11.4040566 |
| | -5.0226538 | 5 | 4.2541872 |
| 5p | 16.0477928 | 2 | 4.2665146 |
| | - 18.4549978 | 3 | 7.1073204 |
| | 8.9776521 | 4 | 7.1349628 |
| | -8.4608734 | 5 | 6.0800546 |
| 5d | 1.3594932 | 3 | 9.6533726 |
| | 3.2152389 | 4 | 6.0685123 |
| | -4.3486283 | 5 | 8.7436015 |
| 5f | 2.9310063 | 4 | 6.5130178 |
| | -2.7708161 | 5 | 6.3994632 |

Table 2Method of determining the radial functions.

wavefunctions are represented as

$$\Psi(J) = \sum_{j=1}^{M} a_{ij} \phi_j(\infty_j L_j S_j J), \tag{1}$$

where (ϕ_j) denotes a set of single-configuration wavefunctions (configuration state functions: CSFs); (α_j) defines the coupling of the angular momenta of the electrons; and the orbital L_j and spin S_j angular momenta are coupled to give the total angular momentum *J*. The mixing coefficients (a_{ij}) are obtained by diagonalizing the Breit–Pauli Hamiltonian with respect to the basis (ϕ_i) .

The Hamiltonian used in this work consists of the nonrelativistic electrostatic terms plus the one-body mass correction, the Darwin term, and the spin–orbit, spinother-orbit, and spin–spin operators of the Breit–Pauli Hamiltonian. The inclusion of the mass correction and Darwin terms shifts the energy of a configuration as a whole, while the spin–orbit, spin–other-orbit and spin– spin terms give rise to the fine-structure splitting. The general configuration-interaction code CIV3 uses orthonormal orbitals. Their radial functions $P_{nl}(r)$ are expressed as superpositions of normalized Slater-type Orbitals (STOs) of the form

$$P_{nl}(r) = \sum_{i=1}^{k} C_{jnl}\chi_{jnl}(r), \qquad (2)$$

where C_{inl} are the Clementi-type [9] coefficients and

$$\chi_{jnl}(r) = \frac{(2\xi_{jnl})^{\mathbf{I}_{jnl} + \frac{1}{2}}}{[(2\mathbf{I}_{jnl})!]^{\frac{1}{2}}} r^{\mathbf{I}_{jnl}} \exp(-\xi_{jnl}r)$$
(3)

with the integer

 $I_{jnl} \ge l+1$

The parameters C_{jnl} and ξ_{jnl} are determined variationally while the parameters I_{jnl} are normally kept fixed in any optimization process. The radial functions are chosen to satisfy the orthonormality condition

$$\int_{0}^{\infty} P_{nl}(r) P_{n'l}(r) dr = \delta_{nn'}; \ l+1 < n' \le n$$
(4)

| Orbital | Process of optimization | | | | |
|-----------------------|-----------------------------------|---------------------------------------------------------------------------------------------------------------------------------------------------------|--|--|--|
| 1s, 2s, 2p, 3s, 3p | Hartree–Fock orbit | als of 3s ² 3p ² P ⁰ of Ge XX (Clementi and Roetti [9]) | | | |
| | Eigenvalue minimized | Configurations | | | |
| | Real spectroscopic | | | | |
| 3d | 3s ² 3d ² D | 3s ² 3d ² D | | | |
| 4s | 3s ² 4s ² S | 3s ² 4s ² S | | | |
| 4p | 3s ² 4p ² P | 3s ² 4p ² P, 3s ² 3p ² P ⁰ | | | |
| 4d | 3s ² 4d ² D | 3s ² 4d ² D, 3s ² 3d ² D | | | |
| 4f | 3s ² 4f ² F | 3s ² 4f ² F | | | |
| | Valence correlation | 1 | | | |
| 5s | 3s ² 5s ² S | $3s3pnd (3 \le n \le 5)$, $3s3dnp (n = 4, 5)$, $3s3dnf (n = 4, 5)$, $3s4pnd (n = 4, 5)$, $3s4dnf (n = 4, 5)$, $3s4d5p$, $3s4f5d$, $3s5p5d$, | | | |
| 5p | 3s ² 5p ² P | 3s5d5f, $3p3dns$ ($n = 4, 5$), $3p4snp$ ($n = 4, 5$), $3p4snd$ ($n = 4, 5$), $3p4d5s$, $3p5s5d$, $3d4p5s$, $3d5s5p$, $4s4pnd$ ($n = 4, 5$), | | | |
| 5d | 3s ² 5d ² D | 4s4d5p, 4s5p5d, 4p4d5s, 4p5s5d, 4d5s5p, 5s5p5d | | | |
| 5f | 3s ² 5f ² F | | | | |

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