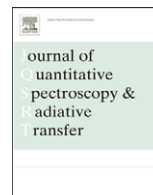




Contents lists available at ScienceDirect

Journal of Quantitative Spectroscopy & Radiative Transfer

journal homepage: www.elsevier.com/locate/jqsrt

Matrix formulations of radiative transfer including the polarization effect in a coupled atmosphere–ocean system

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ARTICLE INFO

Article history:

Received 23 June 2009

Received in revised form

5 November 2009

Accepted 22 November 2009

Keywords:

Radiative transfer

Polarization

Discrete ordinate method

Matrix operator method

GOSAT

ABSTRACT

A vector radiative transfer model has been developed for a coupled atmosphere–ocean system. The radiative transfer scheme is based on the discrete ordinate and matrix operator methods. The reflection/transmission matrices and source vectors are obtained for each atmospheric or oceanic layer through the discrete ordinate solution. The vertically inhomogeneous system is constructed using the matrix operator method, which combines the radiative interaction between the layers. This radiative transfer scheme is flexible for a vertically inhomogeneous system including the oceanic layers as well as the ocean surface. Compared with the benchmark results, the computational error attributable to the radiative transfer scheme has been less than 0.1% in the case of eight discrete ordinate directions. Furthermore, increasing the number of discrete ordinate directions has produced computations with higher accuracy. Based on our radiative transfer scheme, simulations of sun glint radiation have been presented for wavelengths of 670 nm and 1.6 μm. Results of simulations have shown reasonable characteristics of the sun glint radiation such as the strongly peaked, but slightly smoothed radiation by the rough ocean surface and depolarization through multiple scattering by the aerosol-loaded atmosphere. The radiative transfer scheme of this paper has been implemented to the numerical model named *Pstar* as one of the *OpenCLASTR/STAR* radiative transfer code systems, which are widely applied to many radiative transfer problems, including the polarization effect.

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1. Introduction

The top of the Earth's atmosphere is irradiated by solar radiation, emitting thermal infrared radiation to space. That radiation interacts with molecules and particles that are present in the terrestrial atmosphere, and reacts also with the ground surface. Especially, the solar radiation is unpolarized at the top of the atmosphere, but the state of polarization and the radiant intensity are changed because of the radiative transfer process. The process,

including the polarization effect, follows the radiative transfer equation [1].

The discrete ordinate approach is one technique to solve the radiative transfer equation and obtains the radiation field for given boundary conditions. By discretizing the angular coordinates of the phase function and radiance with respect to the zenith angle, the radiative transfer equation is reduced to a linear differential matrix equation. The ways of solving the matrix equation have been studied for the case of multilayered system with underlying ground surface [2–4]. Based on the solutions, the atmospheric radiative transfer (RT) models have been developed for several purposes such as the radiation budget and remote sensing studies. The scalar RT code

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Nomenclature

α	scattering angle defined in Eq. (36)	\mathbf{I}_1	single scattering component of the Stokes vector
$\alpha_l, \beta_l, \gamma_l, \delta_l, \varepsilon_l, \zeta_l$	Greek constants, which are expansion coefficients of $\hat{\mathbf{P}}$	$\tilde{\mathbf{I}}_1$	accurate single scattering solution using the complete phase matrix
α, β	integral constants of homogeneous solution	\mathbf{I}_{TMS}	Stokes vector, whose single scattering component are corrected
$a_1, a_2, a_3, a_4, b_1, b_2$	elements of scattering phase matrix	\mathbf{J}	total source vector
$A_{m,n}^l$	coefficient of the generalized spherical function	Λ^2	Eigenvalue of matrix \mathbf{G}
$\mathbf{A}^{(m)}$	matrix defined in Eq. (16)	l, m, n	integral numbers
$\mathbf{A}^\pm, \mathbf{B}^\pm$	base function for the scaled Stokes vector defined in Eqs. (82a) and (82b)	\mathbf{L}	rotation matrix
β, β_t	incidence and scattering angles with respect to the surface normal vector	μ	cosine of polar angle
$b_n, \mathbf{b}_n, \mathbf{B}_n$	expansion coefficients for Planck's function defined in Eqs. (74), (77b), and (77a)	$\tilde{\mu}$	cosine of the polar angle in the ocean
B	Planck function	μ_n	cosine of zenith angle for surface normal vector
\mathbf{B}	thermal emission vector	μ_i, w_i	double-Gaussian quadrature points and weights
\mathbf{B}_l	coefficient matrix for $\mathbf{A}^{(m)}$, which contains the Greek constants	m_a, m_o	real parts of the refractive index of the air and the seawater
γ	coefficients of particular solution ξ_S defined in Eq. (73)	\tilde{m}	relative refractive index
\mathbf{c}_n	expansion coefficients for Planck function defined in Eq. (78)	M	number of the solar zenith angles
\mathbf{C}	base function defined in Eq. (69a)	\mathbf{M}	matrix of quadrature points defined in Eq. (43)
\mathbf{C}_0	inner product of \mathbf{U} and \mathbf{V}	ν	auxiliary parameter defined in Eq. (39)
$\mathbf{C}_1, \mathbf{C}_2$	Auxiliary matrix defined in Eq. (87)	N	number of quadrature points (discrete ordinate directions)
δ	Dirac delta function	ξ	Stokes vector in eigenspace of \mathbf{G}
δ_{ij}	Kronecker's delta	ξ_{H}	homogeneous solution of transfer Eq. (68)
\mathbf{D}	diagonal matrix for the relation of mirror symmetry	ξ_S, ξ_B	particular solution of transfer Eq. (68)
$\hat{\mathbf{D}}$	\mathbf{D} matrix for all discrete ordinate directions	p	probability density function of the wave slope
$\mathbf{D}_1, \mathbf{D}_2$	auxiliary matrices defined in Eqs. (26a) and (26b)	P_l^m	associated Legendre function
\mathbf{D}_n	coefficient matrix defined in Eq. (84c)	$P_{m,n}^l$	generalized spherical function
ε	source vector of a single layer	\mathbf{P}	phase matrix
\mathbf{E}	exponential functions defined in Eq. (70)	$\hat{\mathbf{P}}$	scattering phase matrix
$\mathbf{E}_{0\pm}$	exponential functions defined in Eqs. (45a) and (45b)	$\mathbf{P}_c^{(m)}, \mathbf{P}_s^{(m)}$	cosine and sine series of phase matrix
\mathbf{E}_4	four-order unit matrix	\mathbf{P}_l^m	expansion matrix for $\mathbf{A}^{(m)}$
f	delta-M truncation factor	\mathbf{P}^*	truncated phase matrix
F	auxiliary function defined in Eq. (38)	R_l^m, T_l^m	auxiliary functions defined in Eqs. (21a) and (21b)
F_0	extraterrestrial solar flux	$\tilde{\mathbf{r}}$	Fresnel reflection matrix for the flat ocean surface
\mathbf{F}	solar flux vector for several solar zenith angles	\mathbf{R}	reflection matrix of a single layer
\mathbf{F}_0	solar flux vector	\mathbf{R}_S	reflection matrix of the ocean surface
\mathbf{g}	source vector defined in Eq. (63) and used in transfer Eq. (61)	$\tilde{\mathbf{R}}$	reflection matrix of the rough ocean surface
G	shadowing factor	σ^2	mean square slope of wave facets
\mathbf{G}	coefficient matrix of transfer Eq. (61)	σ_S, σ_B	source vector defined in Eqs. (59) and (60) and used in transfer Eq. (57)
H	Heaviside's step function	\mathbf{S}	base function defined in Eq. (69b)
$\mathbf{H}, \tilde{\mathbf{H}}$	coefficient matrices defined in Eqs. (98a) and (98b) and used in source vector (97)	\mathbf{S}_S	source vector of single scattering
θ	polar angle	\mathbf{S}_B	source vector of thermal emission
Θ	scattering angle	$\mathbf{S}_k^{(m)}$	azimuthally separated source vector defined in Eq. (30)
I, Q, U, V	Stokes parameters	$\mathbf{S}_{S+,k}^{(m)}, \mathbf{S}_{S-,k}^{(m)}, \mathbf{S}_{B,k}^{(m)}$	source terms defined in Eqs. (31a)–(31e)
I_p	polarized radiance	τ	optical depth
\mathbf{I}	Stokes vector	τ_a	optical depth from the top of the atmosphere to the ocean surface
$\tilde{\mathbf{I}}$	diffused component of the Stokes vector	τ_c	optical thickness of a single layer
$\mathbf{I}_1^{(m)}, \mathbf{I}_2^{(m)}$	azimuthally separated Stokes vector defined in Eqs. (28a) and (28b)	$\tilde{\mathbf{t}}$	Fresnel transmission matrix for the flat ocean surface
		\mathbf{T}	transmission matrix of a single layer
		\mathbf{T}_S	transmission matrix of the ocean surface

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