

# Promoting defect formation and microwave loss properties in $\delta$ -MnO<sub>2</sub> via Co doping: A first-principles study



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## ABSTRACT

The  $\delta$ -MnO<sub>2</sub> has a layered structure and is expected to be a good absorbing material. Here, CASTEP was used for a theoretical study of the microwave absorption properties of  $\delta$ -MnO<sub>2</sub> with defects (interstitial and substitutional cobalt atoms, vacancies). The results show that by analyzing the density of states (DOS) and partial density of states (PDOS), the defects change the charge distribution and increase the magnetic moment (increased by about 19 orders of magnitude). The bond length increased from 1.979 Å to 1.980 Å after substitutional defects formed to enhance the displacement polarizability. The changes in the charge distribution increase the atomic polarizability. The presence of crystal defects enhances both the magnetic loss and dielectric loss. In addition, the calculated defect formation energies show that the Co atoms tend to form interstitial atoms in MnO<sub>2</sub> (−16.90 eV), and the oxygen vacancy defects (−0.77 eV) are more easily formed than the manganese vacancy (33.14 eV).

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## 1. Introduction

Manganese dioxide has been widely studied for many years due to its low cost and lack of toxicity. Manganese dioxide has great potential as a catalyst [1], electrode [2], absorber [3], etc. The formation mechanism of MnO<sub>2</sub> has been extensively researched. Manganese dioxide have several different crystalline phases. There are many research results into the morphological and electrochemical performance of MnO<sub>2</sub> [4–6]. When studied as a microwave absorbent, MnO<sub>2</sub> shows great performance in dielectric loss properties via Fe, Co, or Ni doping [7,8]. While the spontaneous magnetization of MnO<sub>2</sub> is almost zero, the magnetic loss is low. Experimental results show that the magnetic moment increased after magnetic ionic doping, but these studies are mostly based on  $\alpha$ -MnO<sub>2</sub> and  $\beta$ -MnO<sub>2</sub>.

Recently, theoretical calculations have become increasingly important. They are a credible way to predict materials properties and can predict electrochemical performance [9–11], chemical reactions [12], and other physical and chemical properties [13,14]. Pure elements [15], metal compounds [16–18], and organic matter [19] have all been studied. First-principles are also used in MnO<sub>2</sub> studies.

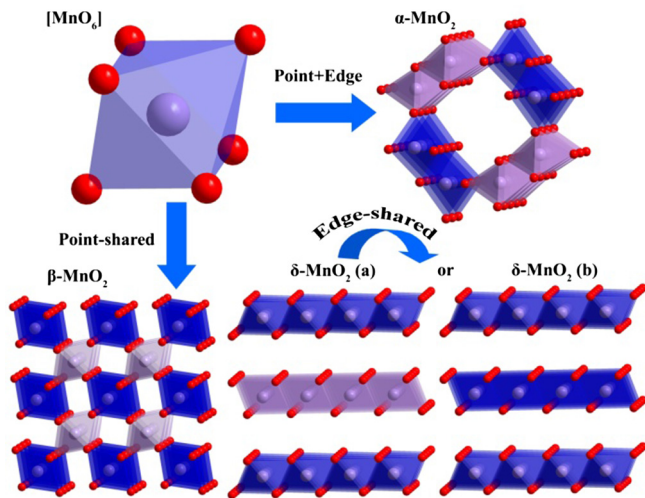
Tompson et al. [20] studied the surface properties of  $\beta$ -MnO<sub>2</sub> by PBE + U to prove that it has potential as battery cathodes, catalysts and supercapacitors. Duan et al. studied the dielectric and electromagnetic properties of  $\alpha$ -MnO<sub>2</sub> [21] and  $\beta$ -MnO<sub>2</sub> [22] after doping and vacancy formation. Yusuke et al. [23] researched the electronic states of four kinds of MnO<sub>2</sub> crystals ( $\alpha$ -,  $\beta$ -  $\delta$ - and  $\lambda$ -MnO<sub>2</sub>), but did not consider the effect of the crystal defect. Kwon et al. [24] studied  $\delta$ -MnO<sub>2</sub> with Ruetschi defects and identified the theoretical basis for enhanced photoconductivity. Similarly, most studies use  $\alpha$ -MnO<sub>2</sub> and  $\beta$ -MnO<sub>2</sub>. Thus far, there are few theoretical studies aiming at the microwave absorption properties of  $\delta$ -MnO<sub>2</sub>.

$\delta$ -MnO<sub>2</sub> is a MnO<sub>2</sub> with a layered structure that can be synthesized via a one-pot soft chemical reaction [25]. There are K<sup>+</sup> ions and H<sub>2</sub>O molecules in the interlayer and this stabilizes the structure. Because of the special structure,  $\delta$ -MnO<sub>2</sub> has good absorption capacity, and it can be used to absorb Cu<sup>2+</sup>, Ni<sup>2+</sup>, Zn<sup>2+</sup>, Cd<sup>2+</sup>, etc. These features can facilitate microwave absorption when the magnetic ions are contained in the interlayer.

Mn is a transition metal from the 4th period. Its electronic configuration is 3d<sup>5</sup>4s<sup>2</sup>. Mn in manganese dioxide is tetravalent, and its outermost electron shells are 3d shells that are partially filled. Mn has a permanent magnetic moment. Different configuration of [MnO<sub>6</sub>] octahedras has different spin configuration, shown in Fig. 1. In microwave frequency, main magnetic loss specific mechanism are natural resonance and eddy current effect [26]. Eddy current effect ( $W_e$ ) can be described as  $W_e = CB_m^2$ , where C is a

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**Fig. 1.**  $[\text{MnO}_6]$  octahedra and the magnetic structure of  $\alpha$ -,  $\beta$ -, and  $\delta$ - $\text{MnO}_2$ . The Mn atoms are purple and O are red. The blue and lavender octahedra represent the different spin directions. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

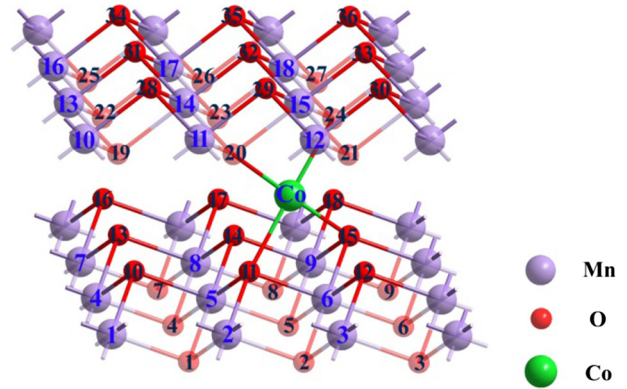
value related with the frequency of electric field, the sample size and the resistivity.  $B_m$  is the magnetic induction intensity.  $B_m$  is proportional to applied magnetic field and magnetization intensity. Magnetization intensity is proportional to net magnetic moment. The net magnetic moment can be found from calculation method. The initial setup of spin configuration will influence the calculation results. According to the Goodenough-Kanamori rule, when two  $[\text{MnO}_6]$  octahedra share the same ridge and these two Mn atoms have the same spin directions—this Mn-Mn coupling is ferromagnetic (FM). When two  $[\text{MnO}_6]$  octahedra share the same point and these two Mn atoms have different spin directions, then this is called antiferromagnetic (AFM) [23]. The  $[\text{MnO}_6]$  octahedra are edge-shared in each layer of the  $\delta$ - $\text{MnO}_2$ . There are two kinds of spin configuration schemes, shown as Fig. 1  $\delta$ - $\text{MnO}_2$  (a) and (b). There is almost no energy difference between the FM and AFM configurations [23]. In the experimental situation, pure  $\text{MnO}_2$  exhibited nonmagnetic character. So in this paper, spin configuration followed as  $\delta$ - $\text{MnO}_2$  (a): each layer of  $\delta$ - $\text{MnO}_2$  should be FM configured, and adjacent layers show AFM configuration.

In this paper, the dielectric and magnetic properties of  $\delta$ - $\text{MnO}_2$  are discussed by first-principles calculations. Interionic electrostatic interactions are analyzed to explain the electromagnetic wave absorption mechanisms of  $\delta$ - $\text{MnO}_2$ . The influence of crystal defects on the microwave loss properties is discussed. This work complements previous theoretical studies on  $\text{MnO}_2$ .

## 2. Models and computational methods

The P- $\text{MnO}_2$  model is an ideal  $\delta$ - $\text{MnO}_2$  supercell ( $3 \times 3 \times 1$ ,  $\text{Mn}_8\text{O}_{16}$ ). The model of  $\delta$ - $\text{MnO}_2$  is determined as in the literature [27] (space group  $P6_3/mmc$ ,  $a = b = 2.840 \text{ \AA}$ ,  $c = 14.031 \text{ \AA}$ ), which was layered  $\text{MnO}_2$ . There is no pure  $\delta$ - $\text{MnO}_2$  in nature or even in the laboratory. Pure  $\delta$ - $\text{MnO}_2$  is an ideal product. In order to simplify the calculation process, the interlayer  $\text{K}^+$  and interlayer  $\text{H}_2\text{O}$  were moved from the crystal. The frame structure of  $\delta$ - $\text{MnO}_2$  was retained. Although the model is different from the experimental results, it can also show the properties of  $\delta$ - $\text{MnO}_2$  [24,28]. The atoms have been numbered as seen in Fig. 2.

This study contains five different models. Pure  $\delta$ - $\text{MnO}_2$  ( $\text{Mn}_{18}\text{O}_{36}$ ) is marked as P- $\text{MnO}_2$ , and  $\delta$ - $\text{MnO}_2$  with an extra Co atom (interstitials,  $\text{Mn}_{18}\text{CoO}_{36}$ ) is marked as I- $\text{MnO}_2$ . The  $\delta$ - $\text{MnO}_2$  with a Co atom substituted for the Mn atom ( $\text{Mn}_{17}\text{CoO}_{36}$ ) is marked as S-



**Fig. 2.** Numbered atoms. The Mn atoms are purple, O are red, and Co are green. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

$\text{MnO}_2$ , and the vacancy-O  $\delta$ - $\text{MnO}_2$  ( $\text{Mn}_{18}\text{O}_{35}$ ) is marked as v-O and vacancy-Mn  $\delta$ - $\text{MnO}_2$  ( $\text{Mn}_{17}\text{O}_{36}$ ) is marked as v-Mn. These models are shown in Fig. 3. The I- $\text{MnO}_2$  is formed by putting one Co atom in the P- $\text{MnO}_2$  crystal, and the doping site is one of the 6 h sites (Wyckoff position), which has been tested as shown in Fig. 3(b). The S- $\text{MnO}_2$  model (Fig. 3(c)) is built by replacing the No. 5 Mn atom with a Co atom, which based on P- $\text{MnO}_2$ . The doping concentrations are 5.5 mol%. The initial spin states of the Co atom follow the states of the Mn atom that are replaced. The v-O (Fig. 3(d)) and v-Mn (Fig. 3(e)) are constructed by removing one O atom No. 14 and Mn atom No. 5 from P- $\text{MnO}_2$ , respectively.

In this study, the CASTEP code [29]—which is based on the density functional theory (DFT) framework to probe the interactions and correlations of the series of solids based on the  $\delta$ - $\text{MnO}_2$ —and the plane wave basis is set; the ultrasoft pseudopotential was used. The Perdew Burke Ernzerhof (PBE) and generalized-gradient approximation (GGA) were also used. The GGA usually underestimates the band gaps of semiconductors. The local density approximation (LDA) was used with Hubbard U corrections (U-J) of 2.5 eV to 3d electrons (LSDA + U) [21]. The DFT + U was used to correct the van der Waals. After testing, a  $4 \times 4 \times 2$  Monkhorst-Pack grid in the Brillouin zone and a 500 eV cutoff energy were used. Mulliken population analysis have been calculated to obtain the information about overlap population and Mulliken charges. In CASTEP, related calculation method has been described in previous papers [30]. The entire calculation process is carried out in the reciprocal space.

## 3. Results and discussion

After geometrical optimization, the lattice constants were changed. The relaxed lattice constants of the supercell ( $3 \times 3 \times 1$ , P- $\text{MnO}_2$ ) are  $a = b = 8.934 \text{ \AA}$  and  $c = 9.791 \text{ \AA}$ . The lattice constants of theoretical samples are different from the experimental samples. The value of c in particular is much smaller than the experimental one (14.03  $\text{\AA}$ ). The differences may be due to the experimental samples that contain interlayer  $\text{K}^+$  and interlayer  $\text{H}_2\text{O}$ . versus pure  $\text{MnO}_2$ , the lattice constants of defective structures have slight differences in the a-axis and b-axis, while the changes in c-axis are obvious. All defective models have different degrees of decrease in the c-axis. The contractions may be caused by the Jahn-Teller distortion [31].

### 3.1. Electronic structures and magnetic loss

Almost all properties of materials are caused by behavior of electrons [32]. The density of states (DOS) and partial density of

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