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## Radiation defect dynamics studied by pulsed ion beams

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## ABSTRACT

The formation of stable radiation damage in solids often proceeds via complex dynamic annealing processes, involving migration and interaction of ballistically-generated point defects. Our current understanding of the underlying physics is still not sufficient for predicting radiation damage even for Si, which is arguably the simplest and most extensively studied material. The complexity of radiation damage is closely related to radiation defect dynamics. Here, we demonstrate how defect interaction dynamics can be studied by pulsed beam irradiation when the total ion fluence is split into a train of equal square pulses. By varying the passive portion of the beam duty cycle, we measure a characteristic time constant of dynamic annealing and, hence, the defect relaxation rate. Measurements of stable lattice disorder as a function of the active portion of the beam duty cycle give an effective defect diffusion length. We illustrate the pulsed beam method with examples for Si bombarded at 100 °C with 500 keV Ar ions.

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## 1. Introduction

The formation of stable radiation damage in materials is a dynamic phenomenon. For most irradiation environments relevant to nuclear technologies and radiation-assisted processing of electronic materials, the damage buildup proceeds via complex dynamic annealing (DA) processes (see, for example, reviews in [1–3]). The DA involves migration and interaction of ballistically-generated point defects after the thermalization of collision cascades, over time scales  $\gtrsim 1$  ps. Such DA leads to point defect annihilation or, in turn, to clustering and the growth of various extended lattice defects, such as dislocations, stacking faults, planar defects, non-stoichiometric inclusions, amorphous zones, bubbles, and voids. As a result, lattice disorder experimentally observed after irradiation can significantly depart from predictions based only on collisional processes [1–3]. The DA manifests as a non-trivial dependence of radiation damage on irradiation conditions. This currently precludes the prediction and ultimately control of radiation damage.

Most previous studies of DA have focused on the measurements of the dependencies of damage on the beam flux (i.e., the dose rate) and on the density of collision cascades, determined by ion mass and energy for a given target material [1–6]. The flux effect is caused by the interaction of mobile defects originating in different collision cascades created in close proximity of each other. This

involves a convolution of both spatial (such as the defect diffusion length,  $L_d$ ) and temporal (i.e., the lifetime of mobile point defects,  $\tau$ ) parameters of DA. Attempts to understand DA from flux effect data have been somewhat cumbersome [3,5,7] since serious assumptions are required to extract the  $\tau$  and  $L_d$  from experimental data by modeling defect interaction processes.

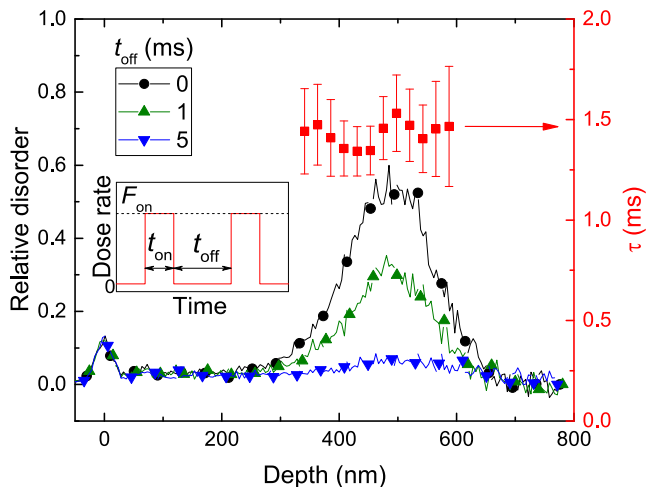
Here, we describe the development of a pulsed beam method (PBM) for accessing the dynamic regime of defect accumulation and measuring both  $\tau$  and  $L_d$  [3,8–13]. The inset in Fig. 1 shows a schematic of the time dependence of the beam flux during pulsed beam irradiation, defining the additional PBM parameters ( $t_{on}$ ,  $t_{off}$ , and  $F_{on}$ ). In such experiments, the total ion fluence is split into a train of equal pulses. These PBM experiments involve ion irradiation of a series of specimens with all except one of the irradiation parameters fixed (either  $t_{off}$  or  $t_{on}$ ). Irradiation is followed by the measurement of the level of stable disorder by, for example, ion channeling. Parameters  $\tau$  and  $L_d$  are evaluated based on the analysis of experimental dependencies of the level of stable disorder on  $t_{off}$  and  $t_{on}$ , respectively [3,8–13].

## 2. Experimental

Float-zone grown (100) Si single crystals (with a resistivity of  $\sim 5 \Omega \text{ cm}$ ) were bombarded at 100 °C with 500 keV  $^{40}\text{Ar}^+$  ions at 7° off the [100] direction. To improve thermal contact, the samples were attached to the Cu sample holder with Ag paste. All irradiations were performed in a broad beam mode [8]. In each irradiation

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**Fig. 1.** Left axis: selected depth profiles of relative disorder in Si bombarded at 100 °C with a pulsed beam of 500 keV Ar ions with  $F_{on} = 1.8 \times 10^{13} \text{ cm}^{-2} \text{ s}^{-1}$ ,  $t_{on} = 1 \text{ ms}$ , a total ion fluence of  $8.75 \times 10^{14} \text{ cm}^{-2}$ , and different  $t_{off}$  values given in the legend. For clarity, only every 10th experimental point is depicted. Right axis: the depth dependence of the dynamic annealing time constant ( $\tau$ ) measured as described in the text. The inset is a schematic of the time dependence of the instantaneous beam flux for pulsed beam irradiation, defining  $t_{on}$ ,  $t_{off}$ , and  $F_{on}$ .

run, the total ion fluence was split into a train of square pulses, each with an instantaneous beam flux  $F_{on}$  and a fluence per pulse of  $F_{on}t_{on}$  (see the inset in Fig. 1). For  $\tau$  measurements [3,8,10,11,13],  $F_{on} = 1.8 \times 10^{13} \text{ cm}^{-2} \text{ s}^{-1}$ ,  $t_{on} = 1 \text{ ms}$ , the total fluence was  $8.8 \times 10^{14} \text{ cm}^{-2}$ , and adjacent pulses were separated by time  $t_{off}$ , which was varied between 0.2 and 10 ms. For  $L_d$  measurements [3,9,12],  $F_{on} = 10^{14} \text{ cm}^{-2} \text{ s}^{-1}$ , the total fluence was  $1.2 \times 10^{15} \text{ cm}^{-2}$ ,  $t_{on}$  was varied between 0.1 and 3 ms, and each pulse was separated by  $t_{off} = 20 \text{ ms}$ , which, as will be shown below, was much greater than the  $\tau$  value.

The dependence of stable lattice damage on  $t_{off}$  and  $t_{on}$  was studied *ex-situ* at room temperature by ion channeling. Depth profiles of lattice disorder were measured with 2 MeV  $^4\text{He}^+$  ions incident along the [100] direction and backscattered into a detector at 164° relative to the incident beam direction. Raw channeling spectra were analyzed with one of the conventional algorithms [14] for extracting depth profiles of relative disorder. Values of average relative bulk disorder ( $n$ ) were obtained by averaging depth profiles of relative disorder over 20 channels ( $\sim 38 \text{ nm}$ ) centered on the bulk damage peak maximum. Error bars of  $n$  are standard deviations. The total ion fluence was chosen such that, for continuous beam irradiation (i.e.,  $t_{off} = 0$ ),  $n$  was  $\sim 0.6$  (with  $n = 1$  corresponding to full amorphization). The 4 MV ion accelerator (National Electrostatics Corporation, model 4UH) at Lawrence Livermore National Laboratory was used for both ion irradiation and ion beam analysis. A more detailed description of the experimental arrangement can be found elsewhere [3,8–13]. The choice of the pulsing parameters for measurements of  $\tau$  and  $L_d$  was discussed in detail in [3].

### 3. Results and discussion

#### 3.1. Measurement of $\tau$

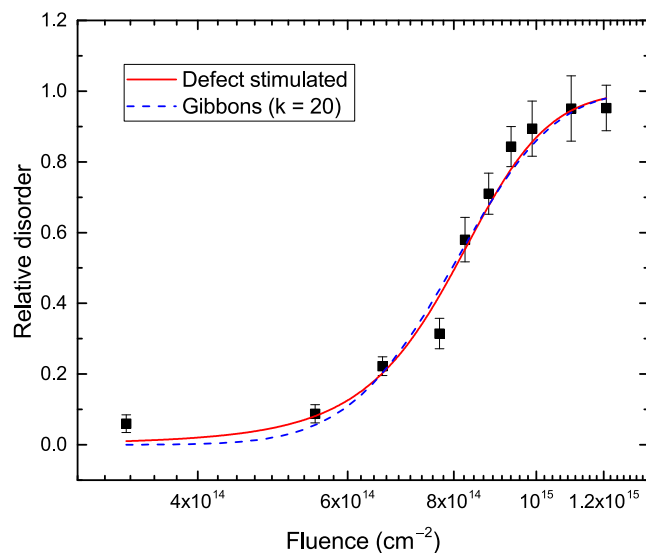
The measurement of either  $\tau$  or  $L_d$  for any given combination of the target material, temperature, ion mass, energy, and  $F_{on}$  starts with the selection of the total ion fluence resulting in the damage level of interest. Hence, in the absence of previous experimental data or predictive models, the first experiment in the PBM is the

measurement of the traditional damage buildup curve. This is shown in Fig. 2 for Si bombarded at 100 °C with a continuous beam of 500 keV Ar ions with  $F_{on} = 1.8 \times 10^{13} \text{ cm}^{-2} \text{ s}^{-1}$ . The damage buildup revealed by Fig. 2 is highly sigmoidal. Data can be readily fitted with a phenomenological defect stimulated and direct amorphization model [15] (with fitting parameters  $\sigma_a = 4.4 \times 10^{-14} \text{ cm}^{-2}$  and  $\sigma_s = 9.6 \times 10^{-15} \text{ cm}^{-2}$ ) and Gibbons overlap model [16] (with fitting parameters  $k = 20$ ,  $A = 2.6 \times 10^{-14} \text{ cm}^{-2}$ ). The Gibbons overlap model required 20 overlaps of disordered regions to fit the data, which suggests that the physics of stable damage formation for these irradiation conditions is beyond what this model is capable of describing.

Fig. 1 (left axis) shows representative depth profiles of relative disorder in Si bombarded with Ar ions with three different  $t_{off}$  values (given in the legend) and all the other parameters kept constant. It is seen that, with increasing  $t_{off}$ , while the damage level in the first  $\sim 250 \text{ nm}$  from the sample surface remains unchanged, the intensity of the bulk disorder peak ( $n$ ) dramatically decreases. This observation suggests different dynamic mechanisms of bulk and surface disordering. It is consistent with our previous depth-resolved studies of defect interaction dynamics in Si [3,8,9,13].

Fig. 3 (bottom axis) shows the  $n(t_{off})$  dependence for Si under Ar ion bombardment. The monotonic  $n(t_{off})$  decay, revealed by Fig. 3 (bottom axis) is related to the interaction of defects generated in different pulses (and, hence, in different cascades). As the beam is pulsed off the target, the concentration of mobile defects decreases via DA with a characteristic time constant  $\tau$ . For irradiation with  $t_{off} \gg \tau$ , DA processes have essentially decayed in time intervals between individual ion pulses.

The time constant of the dominant DA process ( $\tau$ ) can be quantitatively evaluated by analyzing such experimental  $n(t_{off})$  dependencies. For example,  $\tau$  can be obtained by fitting  $n(t_{off})$  dependencies with either the first order ( $n(t_{off}) = n_{\infty} + (n(0) - n_{\infty}) \exp(-t_{off}/\tau_1)$ ) or the second order ( $n(t_{off}) = n_{\infty} + \frac{n(0) - n_{\infty}}{1 + \frac{t_{off}}{\tau_2}}$ ) decay equations. Here,  $n_{\infty}$  is relative disorder for  $t_{off} \gg \tau_{1,2}$ . The “1” and “2” subscripts refer to whether the



**Fig. 2.** Dose dependence of relative disorder at the maximum of the bulk defect peak for Si bombarded at 100 °C with a continuous beam of 500 keV Ar ions with a beam flux of  $1.8 \times 10^{13} \text{ cm}^{-2} \text{ s}^{-1}$ . Results of fitting the data with the defect stimulated amorphization model [15] and the Gibbons overlap model [16] with  $k = 20$  are shown by solid and dashed lines, respectively.

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