

# On $SL(2, R)$ symmetry in nonlinear electrodynamics theories

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## Abstract

Recently, it has been observed that the Noether–Gaillard–Zumino (NGZ) identity holds order by order in  $\alpha'$  expansion in nonlinear electrodynamics theories as Born–Infeld (BI) and Bossard–Nicolai (BN). The nonlinear electrodynamics theory that couples to an axion field is invariant under the  $SL(2, R)$  duality in all orders of  $\alpha'$  expansion in the Einstein frame. In this paper we show that there are the  $SL(2, R)$  invariant forms of the energy momentum tensors of axion-nonlinear electrodynamics theories in the Einstein frame. These  $SL(2, R)$  invariant structures appear in the energy momentum tensors of BI and BN theories at all orders of  $\alpha'$  expansion. The  $SL(2, R)$  symmetry appears in the BI and BN Lagrangians as a multiplication of Maxwell Lagrangian and a series of  $SL(2, R)$  invariant structures.

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## 1. Introduction

Duality transformations in nonlinear electrodynamics has been studied in [1]. Classical electromagnetism is the most familiar duality-invariant theory. The Maxwell's Hamiltonian and the equations of motion are invariant under rotations. Note that the Lagrangian, however, is not

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invariant. This brings that nonlinear deformations of Lagrangian will require modifications which are also non-invariant.

Commonly, duality transformations may be found out in the path integral as a Legendre transform. Given some Lagrangian  $\mathcal{L}(F)$  depending only on the field strength of a vector field, one can construct [2]

$$\mathcal{L}(\tilde{F}, G) = \mathcal{L}(F) - \frac{1}{2} \epsilon^{abcd} F_{ab} \partial_c \tilde{A}_d \quad (1)$$

in which  $F$  is treated as a fundamental field. The classical equations of motion for  $F$  require that  $G_{ab} = \partial_a \tilde{A}_b - \partial_b \tilde{A}_a$  is related to  $F$  by

$$G_{ab} = -2 \frac{\partial \mathcal{L}(F)}{\partial F^{ab}} \quad (2)$$

through  $G_{ab} = -\frac{1}{2} \epsilon_{abcd} \tilde{G}^{cd}$  and  $\tilde{G}^{ab} = \frac{1}{2} \epsilon^{abcd} G_{cd}$ .

Consistency of the duality constraint can be expressed as a requirement in which the Lagrangian must transform under duality in a particular way, defined by the Noether–Gaillard–Zumino (NGZ) identity [3].

We consider  $\mathcal{L}_{inv}$  that is invariant under the following transformation:

$$\delta \begin{pmatrix} \tilde{F} \\ G \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} \tilde{F} \\ G \end{pmatrix} \quad (3)$$

where  $A^T = -D$ ,  $B^T = B$  and  $C^T = C$  are the infinitesimal parameters of the transformations [2]. By applying the duality symmetry, the NGZ identity could be written as following [4]:

$$\mathcal{L} = \mathcal{L}_{inv} - \frac{1}{4} G_{ab} F^{ab}. \quad (4)$$

When the theory only has linear duality (e.g. only  $F^2$  terms in the action)  $\delta \mathcal{L}_{inv} / \delta F$  vanishes. So it could be found that any higher order dependence ( $F^4$ ,  $F^6$ , ...) must be part of  $\mathcal{L}_{inv}$ . The NGZ identity, along with (2) can be solved to find  $G(F)$  and various Lagrangians which provide a duality symmetry between equations of motion and Bianchi identities. We will discuss two cases of nonlinear deformations of the Maxwell theory that depend only on  $F$ 's without derivatives.

The NGZ identity has a significant consequence for the duality rotations properties of the energy-momentum tensor. The energy-momentum tensor, which can be obtained as the variational derivative of the Lagrangian with respect to the gravitational field, is invariant under duality transformation [5]. It has been found as [6]

$$T_{ab} = g_{ab} \mathcal{L} + G_a{}^c F_{bc} \quad (5)$$

Considering this relation and (4), it could be found that  $\mathcal{L}_{inv} = 1/4 T_a{}^a$ .

The Born–Infeld Lagrangian density can be written intelligently in term of the square root of a determinant [7]:

$$\mathcal{L}_{BI} = \sqrt{-\det(\eta_{ab} + F_{ab})} - 1 \quad (6)$$

where the fundamental (*scale*)<sup>2</sup> ( $= T^{-1} = 2\pi\alpha'$  in the string theory context) has been set equal to 1, has many considerable aspects, including electro-magnetic duality symmetry [8]. Since in four dimensions:

$$-\det(\eta_{ab} + F_{ab}) = \frac{1}{2} F_{ab} F^{ab} - \frac{1}{16} (F_{ab} \tilde{F}^{ab})^2, \quad \tilde{F}^{ab} = \frac{1}{2} \epsilon^{abcd} F_{cd}, \quad (7)$$

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