



Investigation of activation cross sections for deuteron induced reactions on strontium up to 50 MeV



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ABSTRACT

Excitation functions were measured for the $^{nat}\text{Sr}(\text{d},\text{x})^{88,87\text{m},87\text{g},86\text{g},85\text{g}}\text{Y}$, $^{87\text{m},85\text{g},83\text{g},82}\text{Sr}$, $^{86\text{g},84\text{g},83,82\text{m},81}\text{Rb}$ reactions by the stacked foil activation technique and high-resolution gamma-spectrometry up to 50 MeV. We present the first experimental activation cross section data for all investigated reactions. Our experimental data are compared with the TALYS code results as available in the TENDL-2015 on-line library. Use of deuteron induced reactions on Sr for production of medical isotopes is discussed.

1. Introduction

We are performing a systematic study of activation cross-sections of light charged particle induced reactions. The aim is to get reliable experimental data for optimizing production of specific radionuclides for different applications and to consolidate the database allowing improvement of theoretical models. In the present work we report experimental cross section data for deuteron induced reactions on strontium. The investigated activation radio-products are of interest in different applications:

^{90}Y is widely used in brachytherapy and radioimmunotherapy but does not emit photons for imaging (Wiseman et al., 1999);

^{87}Y is a γ -emitting isotope that can be used to quantify the biodistribution of pharmaceuticals;

^{88}Y is a γ -emitting isotope used for calibration source.

^{86}Y is a positron emitting isotope that can be used to quantify regional kinetics of yttrium via positron emission tomography (PET).

$^{87\text{m}}\text{Sr}$ can be used in both diagnostic and therapeutic techniques for various skeletal diseases.

^{83}Sr is β^+ emitter, potentially useful radionuclide for therapy planning prior to the use of the longer lived β^+ emitter ^{89}Sr .

^{82}Sr is used exclusively to manufacture $^{82}\text{Sr}/^{82}\text{Rb}$ generators widely used in PET myocardial perfusion imaging.

^{85}Sr is a γ -emitting isotope used for calibration source.

^{86}Rb is used as tracer in isotopic measurements of field metabolic rate (Bradshaw et al., 2007).

^{84}Rb is used in real-time tumor tracking system using implanted

positron emission markers. Its half-life is comparable to the duration of a standard radiation therapy procedure.

$^{82\text{m}}\text{Rb}$ was proposed as a substitute for cardio-PET clinical studies due to its suitable physical properties (Rowshanfarzad et al., 2006).

^{81}Rb is used through the $^{81}\text{Rb}/^{81\text{m}}\text{Kr}$ generator where the short-lived eluent is used for investigation of pulmonary ventilation, or before surgical investigation.

In the frame of a systematic study of production routes of diagnostic and therapeutic radioisotopes we already made a series of measurements on excitation functions for the above mentioned activation products via proton, deuteron, alpha and ^3He -particle induced reactions on krypton, rubidium, strontium, yttrium and zirconium targets (Tárkányi et al., 1988a, 1988b, 1990, 2004, 2005, 2015; Qaim et al., 1988, 1994; Kovács et al., 1991a, 1991b, 1992; Fenyvesi et al., 1992; Blessing et al., 1997; Ido et al., 2002a, 2002b; Uddin et al., 2005, 2006, 2007; Spahn et al., 2006). In these works some comparisons of the different production routes are also discussed but recently, new experimental data and reviews were published (Qaim et al., 2001; Hermanne et al., 2001; Filosofov et al., 2010; Betak et al., 2011; Zaneb et al., 2016; Capote et al., 2017) and we hence compare only shortly the production routes of the most widely used isotopes.

2. Experimental procedure

The cross-section measurements were performed using the activation method combining a stacked foil irradiation and high resolution HPGe gamma-ray spectrometry. Cross-section data were deduced

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relative to the re-measured excitation functions of monitor reactions. The stack was irradiated in a short Faraday cup at an external beam line of the Cyclone 90 cyclotron of the Université Catholique in Louvain la Neuve (LLN) for 40 min with a 50 MeV, 50 nA deuteron beam.

The stack contained a sequence of 10 blocks of Al (10 μm), $\text{Sr}(\text{NO}_3)_2$ (9–55 μm), Al (50 μm), Er (32 μm), Al (10 μm), $\text{Ba}(\text{NO}_3)_2$ (2 μm , sedimented), Al (50 μm) and Ag (10 μm) foils followed by 15 blocks of Ti (10 μm), Al (10 μm), $\text{Ba}(\text{NO}_3)_2$ (2 μm , sedimented), Al (50 μm) Ti (10 μm), Al (10 μm), $\text{Ba}(\text{NO}_2)_2$ (2 μm , sedimented), Al (50 μm). The 25 $\text{Ba}(\text{NO}_3)_2$ targets covered the 49.4–9.2 MeV energy range. The thin foils were purchased from Goodfellow Ltd., the compounds used for sedimentation were spectroscopic grade chemical materials from our storage. The sedimented targets were prepared in such a way that for the calculated mass of compound powder a given amount of evopreme glue was dissolved in chloroform and the compound powder was added to this solution under continuous ultrasonic mixing. The solution was poured in a 11 mm spot with Al backing and incubated until the chloroform content was completely evaporated and the compound was sedimented. The thicknesses of the targets were determined by weight and area measurement, in the case of sedimented samples the stoichiometry was also taken into account.

Four series of gamma-ray spectra were measured to follow the decay and were started at different times after end of bombardment (EOB): 3.9–7.1 h, 25.7–45.6 h, 92.1–172.4 h and 1846.2–1282.6 h, respectively. Due to the time required for the transfer of the irradiated targets from LLN to VUB (where the gamma spectra measurements were done), the large number of target foils irradiated in our irradiation session (4 stacks) and the limited detector capacity, it was impossible to optimize the gamma spectra measurements for all targets regarding all produced radioisotopes.

Gamma spectra were evaluated by the automatic fitting algorithm included in the *Genie* (2000) package (*Genie*) or in an iterative process using the *Forgamma* (Szekely, 1985) codes.

The decay data were taken from NUDAT2.6 (NUDAT2.6) and are shown in Table 1 together with the Q values of the contributing reactions (Pritychenko and Sonzogni). The preliminary data of the primary beam energy were taken from the settings of the accelerator; the beam intensity was derived from the Faraday cup results. The recommended data for the simultaneously re-measured excitation functions of the $^{27}\text{Al}(\text{d},\text{x})^{22,24}\text{Na}$ and $^{\text{nat}}\text{Ti}(\text{d},\text{x})^{48}\text{V}$ monitor reactions were taken from the IAEA database (Tárkányi et al., 2001). The cross sections were determined using the well-known activation formula. Production cross sections on $^{\text{nat}}\text{Sr}$ were determined, not taking into account the abundance of the different naturally occurring Sr isotopes (^{84}Sr (0.56%), ^{86}Sr (9.86%), ^{87}Sr (7.00%), ^{88}Sr (82.58%).

Uncertainty on cross-sections was determined according to the recommendation in the ISO guide (Guide ISO, 1993) by taking the sum in quadrature of all individual linear contributions: beam current (7%), beam-loss corrections (max. 1.5%), target thickness (1%), detector efficiency (5%), photo peak area determination and counting statistics (1–20%). The overall uncertainty hence ranges from 9% to 25%.

The median energies in targets were initially (for stack preparation) obtained from a degradation calculation based on the incident energy and the Andersen (Andersen and Ziegler, 1977) polynomial approximation of stopping powers. The complete excitation functions of the $^{27}\text{Al}(\text{d},\text{x})^{22,24}\text{Na}$ (Al from cover and backing) and $^{\text{nat}}\text{Ti}(\text{d},\text{x})^{48}\text{V}$ monitor reactions were re-measured simultaneously with the excitation functions of reactions induced on Ba targets to get final values for the beam energy and intensity in the Sr targets. If needed a corrections were applied based on the results of the fitted monitor reactions (final energy and intensity) (Tárkányi et al., 1991). Uncertainty of energy was estimated by following the cumulative effects during the energy degradation of possible contributing uncertainties (primary energy, target thickness, energy straggling, correction for monitor reaction).

3. Results

The measured cross-sections for the production of $^{88,87\text{m},87\text{g},86\text{g},85\text{g}}\text{Y}$, $^{87\text{m},85\text{m},83\text{g},82}\text{Sr}$, $^{86\text{g},84\text{g},83,82\text{m},81}\text{Rb}$ are shown in Tables 2–4 and Figs. 1–15. The figures also show the theoretical results available in the TALYS 1.6 (Koning et al., 2012) based TENDL-2015 library (Koning et al., 2015). We are presenting cumulative results in several cases, when the half-life of the isomeric state decaying into the ground state is much shorter then that of the ground state. In this case we use the results of a late measurement, when the isomeric state is completely decayed out. The TENDL results are constructed in this case by adding the direct isomeric cross section to the direct ground state cross section.

3.1. The $^{\text{nat}}\text{Sr}(\text{d},\text{x})^{88}\text{Y}$ reaction

The excitation function for production of ^{88}Y ($T_{1/2}=106.627$ d) is shown in Fig. 1. The agreement in magnitude with the theory is acceptable, but the threshold is shifted.

3.2. The $^{\text{nat}}\text{Sr}(\text{d},\text{x})^{87\text{m}}\text{Y}$ reaction

The independent cross sections for formation of the metastable state of the $^{87\text{m}}\text{Y}$ ($T_{1/2}=13.37$ h) are shown in Fig. 2 in comparison with TENDL-2015. The agreement with the theory is good, except between 5 and 13 MeV, where the experiment does not show so large low energy contribution.

3.3. The $^{\text{nat}}\text{Sr}(\text{d},\text{x})^{87\text{g}}\text{Y}$ reaction

The excitation function for production of the ground state of ^{87}Y ($T_{1/2}=79.8$ h) was measured after complete decay of isomeric state ($T_{1/2}=13.37$ h, IT:98.43%) and is hence cumulative. The theory underestimates the experimental data (Fig. 3).

3.4. The $^{\text{nat}}\text{Sr}(\text{d},\text{x})^{86\text{g}}\text{Y}$ reaction

The cross sections for production of ground state of ^{86}Y ($T_{1/2}=14.74$ h) were obtained after the complete decay of shorter lived isomeric state ($T_{1/2}=47.4$ min). The description by TENDL-2015 is good (Fig. 4).

3.5. The $^{\text{nat}}\text{Sr}(\text{d},\text{xn})^{85\text{m}}\text{Y}$ reaction

The radionuclide ^{85}Y has two isomeric states: the shorter-lived higher lying state ($T_{1/2}=4.86$ h) and the ground state ($T_{1/2}=2.68$ h). They are decaying independently. We could obtain independent cross-section data for production of both states. Only a few cross-section points could be derived for the isomeric state. According to Fig. 5 the theory overestimates the experiment.

3.6. The $^{\text{nat}}\text{Sr}(\text{d},\text{xn})^{85\text{g}}\text{Y}$ reaction

The cross-sections for the $^{85\text{g}}\text{Y}$ ground state ($T_{1/2}=2.68$ h) were measured with acceptable statistics at all energy points (Fig. 6). The theory underestimates the experiment.

3.7. The $^{\text{nat}}\text{Sr}(\text{d},\text{x})^{87\text{m}}\text{Sr}$ reaction

The cross-sections for $^{87\text{m}}\text{Sr}$ ($T_{1/2}=2.815$ h) are shown in Fig. 7. Some shift in the shape can be observed in comparison with the TENDL-2015 results. The metastable $^{87\text{m}}\text{Sr}$ is produced directly and through the decay of $^{87\text{m}}\text{Y}$ (13.37 h, ϵ : 1.57%) and $^{87\text{g}}\text{Y}$ (79.8 h, ϵ : 100%) parents. These indirect production paths were not taken into account in the presented TENDL curve, the good agreement means that the contribution of them is very low.

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