Contents lists available at ScienceDirect

Journal of Geometry and Physics

journal homepage: www.elsevier.com/locate/jgp

Rational sphere valued supercocycles in *M*-theory and type IIA string theory

Domenico Fiorenza^a, Hisham Sati^{b,c,*}, Urs Schreiber^d

^a Dipartimento di Matematica, La Sapienza Università di Roma Piazzale Aldo Moro 2, 00185 Rome, Italy

^b University of Pittsburgh, Pittsburgh, PA 15260, USA

^c New York University, Abu Dhabi, United Arab Emirates

^d Mathematics Institute of the Academy, Žitna 25, 115 67 Praha 1, Czech Republic¹

ARTICLE INFO

Article history: Received 26 July 2016 Received in revised form 20 November 2016 Accepted 25 November 2016 Available online 8 December 2016

Keywords: Homotopy Lie algebras Supersymmetry Branes Rational homotopy theory

1. Introduction

ABSTRACT

We show that supercocycles on super L_{∞} -algebras capture, at the rational level, the twisted cohomological charge structure of the fields of *M*-theory and of type IIA string theory. We show that rational 4-sphere-valued supercocycles for M-branes in M-theory descend to supercocycles in type IIA string theory with coefficients in the free loop space of the 4-sphere, to yield the Ramond–Ramond fields in the rational image of twisted K-theory, with the twist given by the B-field. In particular, we derive the M2/M5 \leftrightarrow F1/Dp/NS5 correspondence via dimensional reduction of sphere-valued super- L_{∞} -cocycles.

© 2016 Elsevier B.V. All rights reserved.

(Super)cocycles play an important role in the study of the geometric and topological structures associated with physical theories (see [1] for an earlier survey). In [2] we discussed cocycles of super L_{∞} -algebras (super Lie *n*-algebras for arbitrary n) forming the brane bouquet that gives the WZW terms of all the Green–Schwarz sigma models for all the branes in string theory and *M*-theory. This includes those with gauge fields on their worldvolume, the D-branes and the M5-brane, which were missing in the classical brane scan.

In [3] we had shown that this approach allows deriving the rational image of a twisted cohomology theory that unifies the M2-brane charges and the M5-brane charges (this is recalled in Section 2). Rationally this cohomology theory turns out to be represented by the 4-sphere, hence is cohomotopy in degree 4. This is in higher analogy to the familiar statement that the unification of Dp-brane charges with the F1-brane charge ought to be in twisted K-cohomology theory. (That the fields of *M*-theory should take values in the 4-sphere was first suggested in [4.5].)

In Section 3 we show, at the rational level, that indeed the twisted M2/M5 charges in degree-4 cohomotopy in 11 dimensions dimensionally reduce to the twisted K-theory of the F1/Dp/NS5-brane charges in 10 dimensions (for $p \in$ $\{0, 2, 4\}$, where the dimensionally reduced cohomology theory is represented by the rationalization of the homotopy quotient $\mathcal{L}S^4//S^1$ of the free loop space of the 4-sphere. In particular this exhibits a purely L_{∞} -theoretic derivation, at the rational level, of twisted K-theory as the home of the brane charges in type II string theory. The lift of this twisted charge

http://dx.doi.org/10.1016/j.geomphys.2016.11.024 0393-0440/© 2016 Elsevier B.V. All rights reserved.





CrossMark

Corresponding author at: University of Pittsburgh, Pittsburgh, PA 15260, USA. E-mail address: hsati@nyu.edu (H. Sati).

¹ On leave at MPI Bonn.

structure to M-theory has been an open problem. This may be viewed as one confirmation at the rational level of the proposal in [4.5] on the description of *M*-theory via twisted generalized cohomology.

In the existing literature, the cocycles for the WZW terms of the Dp-branes are instead constructed separately as independent cocycles on extended super-Minkowski spacetime (see [3] for references and for the super L_{∞} -algebraic formulation). In Section 4 we show that the same L_{∞} -descent mechanism which unifies the M2- and M5-brane charges also applies to the separate Dp-brane cocycles, and they descend to again a single cocycle with coefficients in (the rational image of) the relevant truncation of twisted *K*-theory.

The techniques that we use are from geometric homotopy theory [6], cast in computationally powerful algebraic language. Lecture notes accompanying the discussion here may be found in [7]. We consider super L_{∞} -algebras as in [8,2]. These are generalizations of super Lie algebras to super Lie n-algebras, for arbitrary n, where instead of just a super Lie bracket we have brackets of all arities with the Lie bracket being the binary one. More precisely, our construction takes place in the homotopy category of super L_{∞} -algebras, so that a morphism from a super L_{∞} -algebra g to a super L_{∞} -algebra h will actually be a span of morphism

$$\mathfrak{g} \leftarrow \tilde{\mathfrak{g}} \rightarrow \mathfrak{h}$$

where $\tilde{\mathfrak{g}} \to \mathfrak{g}$ is a quasi-isomorphism, i.e., an L_{∞} -morphism inducing an isomorphism of graded vector spaces at the level of cohomology from $H^{\bullet}(\tilde{\mathfrak{g}})$ to $H^{\bullet}(\mathfrak{g})$. Passing from \mathfrak{g} to $\tilde{\mathfrak{g}}$ is an example of resolution. This concept has many incarnations, depending on the context (homotopic, fibrant, cofibrant, projective, injective). For us, what is important is that is a concept of equivalence within a category between the object at hand and another (or a combination of such) that generally behaves in a more utilizable way within the same category.

Furthermore, we will make constant use of the duality between (finite type) super L_{∞} -algebras and differential gradedcommutative super-algebras, identifying a super L_{∞} -algebra g with its Chevalley–Eilenberg algebra CE(g) as in [8]. These Chevalley–Eilenberg algebras of super L_{∞} -algebras are what are called *FDAs* in the supergravity literature (going back to [9]). The point of identifying these as dual to super L_{∞} -algebras is to make manifest their higher gauge theoretic nature and the relevant homotopy theory, which is crucial for the results we present here. For instance, for $p \in \mathbb{N}$, the line (p + 2)-algebra $b^{p+1}\mathbb{R}$, i.e., the chain complex with \mathbb{R} in degree p+1 and zeros everywhere else, corresponds to the Chevalley-Eilenberg algebra

$$\mathsf{CE}(b^{p+1}\mathbb{R}) := \big(\mathbb{R}[g_{p+2}]; \ dg_{p+2} = 0\big),$$

where the generator g_{p+2} has degree p + 2. Notice that $CE(b^{p+1}\mathbb{R})$ is the minimal Sullivan model for the rational space $B^{p+2}\mathbb{R}$, reflecting the fact that $b^{p+1}\mathbb{R}$ is the L_{∞} -algebra corresponding to the ∞ -group $B^{p+1}\mathbb{R} \simeq \Omega B^{p+2}\mathbb{R}$. In order to amplify this relation between L_{∞} -algebras and rational homotopy theory, we also write I(X), or simply IX, for the L_{∞} -algebra whose CE-algebra is a given Sullivan model of finite type for some rational space X:

 $\mathfrak{l}(X) = L_{\infty}$ -algebra dual to given Sullivan model (A_X, d_X) for rationalization of X

i.e.

$$\mathsf{CE}(\mathfrak{l}(X)) := (A_X, d_X)$$

See Appendix A for more details on rational homotopy theory and Sullivan models. For example, with this notation then the rationalized spheres S^n are incarnated as

$$CE(IS^{n}) = \begin{cases} (\mathbb{R}[g_{n}], dg_{n} = 0) & \text{for } n \text{ odd} \\ (\mathbb{R}[g_{n}, g_{2n-1}], dg_{n} = 0, \ dg_{2n-1} = g_{n} \wedge g_{n}) & \text{for } n > 0 \text{ even.} \end{cases}$$

A convenient feature of the dual picture is the following: if $CE(\mathfrak{h}) \rightarrow CE(\mathfrak{g})$ is a relative Sullivan algebra, that is, a cofibration in the standard model structure on differential graded commutative algebras (DGCAs), then the corresponding L_{∞} -morphism $\mathfrak{g} \to \mathfrak{h}$ is a fibration in the model structure whose fibrant objects are L_{∞} -algebras, due to [10, prop. 4.36, prop. 4.42]. Although relative Sullivan algebras do not exhaust fibrations of L_{∞} -algebras, they are flexible enough to allow us to realize all the fibrations we will need in the present article as relative Sullivan algebras. See [11] for more on the homotopy theory of L_{∞} -algebras as a category of fibrant objects.

The model structure whose fibrant objects are L_{∞} -algebras in [10] is for ordinary L_{∞} -algebras, not for super L_{∞} -algebras that we consider here. Nevertheless, the result is readily adapted: A super L_{∞} -algebra g determines a functor $\Lambda \mapsto (\mathfrak{g} \otimes \Lambda)_{\text{even}}$ with values in ordinary L_{∞} -algebras on the category of finitely generated Grassmann algebras Λ , and this construction embeds super L_{∞} -algebras into this functor category. (For super Lie algebras this was observed in [12], see [13] and [14, Cor. 3.3].) Now, by [10, Theorem 4.35], the opposite model structure for ordinary L_{∞} -algebras is cofibrantly generated, and so a standard argument [15, section 11.6] gives that this functor category inherits the corresponding projective model structure. That is the model structure in which the computations in this paper take place. However, we need to invoke only a bare minimum of model category theory; all we use is the computation of homotopy fibers as ordinary fibers of fibration resolutions. In the following we will find it very useful to succinctly capture results via (commuting) diagrams. We will use the notation $hofib(\phi)$ to indicate the homotopy fiber of a morphism ϕ .

Download English Version:

https://daneshyari.com/en/article/5500079

Download Persian Version:

https://daneshyari.com/article/5500079

Daneshyari.com