



Wave breaking for the Fornberg–Whitham equation

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Abstract

In this paper we consider the Fornberg–Whitham (FW) equation and we give sufficient conditions on the initial data to lead to wave-breaking.

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1. Introduction

The Fornberg–Whitham equation

$$u_{xxt} - u_t + \frac{9}{2}u_x u_{xx} + \frac{3}{2}uu_{xxx} - \frac{3}{2}uu_x + u_x = 0 \quad (1)$$

$$u(x, 0) = u_0(x)$$

was proposed as a model for shallow water waves (see [8]). For classical solutions decaying far out in space, e.g. $u(., t) \in H^4(\mathbb{R}) \subset C^3(\mathbb{R})$ with a continuously differentiable dependence on time t , one can write (1) in the equivalent nonlocal conservation law form

$$u_t + \frac{3}{2}uu_x = (1 - \partial_x^2)^{-1} \partial_x u. \quad (2)$$

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For $u(\cdot, 0) \in H^2(\mathbb{R})$, Kato's semi-group approach for quasi-linear hyperbolic equations [10] yields the existence of a unique solution $u \in C([0, T], H^2(\mathbb{R})) \cap C^1([0, T], L^2(\mathbb{R}))$, and if $T < \infty$, then $\limsup_{t \nearrow T} \|u(\cdot, t)\|_{H^2(\mathbb{R})} = \infty$. Indeed, setting $X = L^2(\mathbb{R})$, $Y = H^2(\mathbb{R})$, and choosing the pseudo-differential operator to be $Q = (1 - \partial_x^2)$, we can rewrite (2) as

$$u_t(x, t) = A(u)u + F(u, t)$$

with $A(u) = -\frac{3}{2}u\partial_x$, where the domain $\mathcal{D}(A(u)) = Y \subset X$, and $F(u, t) = Q^{-1}\partial_x u$. The requirements for the theorem for quasilinear equations of the general form

$$u_t + a(u)u_x = b(u)u, \quad t \geq 0,$$

are shown by Kato in [10].

Recently, Thompson and Holmes [12] proved the following blow-up criterion:

Theorem A ([12]). *Let T be the maximal time of existence of the solution to (2) with initial data $u(\cdot, 0) \in H^2(\mathbb{R})$. If $T < \infty$ then*

$$\int_0^T \|u_x(x, t)\|_{L^\infty(\mathbb{R})} dt = +\infty.$$

The aim of this paper is to sharpen this blow-up criterion and to give sufficient conditions for the development of singularities in finite time in a solution to the (FW) equation.

2. Blow-up scenario

In this section, we will refine the blow-up criterion. Before we begin, we would like to recall that wave-breaking occurs when, if the maximal time of existence $T > 0$ is finite, the $\sup_{(t,x) \in [0,T) \times \mathbb{R}} |u(t, x)| < \infty$ whilst the $\sup_{x \in \mathbb{R}} |u_x(t, x)| \rightarrow \infty$ as $t \nearrow T$. We refer the reader to the discussions in [14] and [7].

Proposition 1. *If for some smooth initial data $u_0 \in H^\infty(\mathbb{R})$ the maximal existence time $T > 0$ is finite, then (2) blows-up in finite time if and only if*

$$\liminf_{t \nearrow T} [u_x(t, x)] = -\infty$$

Proof. For simplicity of notation, we denote the operator $(1 - \partial_x^2)^{-1}$ by Θ . Note that this is an isomorphism between $H^s(\mathbb{R})$ and $H^{s+2}(\mathbb{R})$ for every $s \in \mathbb{R}$.

Kato's semigroup theory ensure that $u \in C^1([0, T], H^s(\mathbb{R}))$ for every $s \geq 2$.

We will use the classical energy method in this proof. We begin by multiplying (2) by u and integrating over \mathbb{R} . We get:

$$\partial_t \left\{ \frac{1}{2} \int_{\mathbb{R}} u^2 dx \right\} = \int_{\mathbb{R}} u \Theta_x u dx \quad (3)$$

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