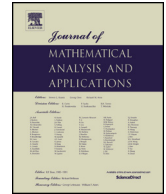




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Eigenvalue inequalities and absence of threshold resonances for waveguide junctions

Konstantin Pankrashkin

Laboratoire de Mathématiques d'Orsay, Univ. Paris-Sud, CNRS, Université Paris-Saclay, 91405 Orsay, France

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ABSTRACT

Let $\Lambda \subset \mathbb{R}^d$ be a domain consisting of several cylinders attached to a bounded center. One says that Λ admits a threshold resonance if there exists a non-trivial bounded function u solving $-\Delta u = \nu u$ in Λ and vanishing at the boundary, where ν is the bottom of the essential spectrum of the Dirichlet Laplacian in Λ . We give a sufficient condition for the absence of threshold resonances in terms of the Laplacian eigenvalues on the center. The proof is elementary and is based on the min–max principle. Some two- and three-dimensional examples and applications to the study of Laplacians on thin networks are discussed.

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1. Introduction

Let $\Lambda \subset \mathbb{R}^d$, $d \geq 2$, be a connected Lipschitz domain which can be represented as a family of several half-infinite cylinders attached to a bounded domain. More precisely, we assume that there exist bounded connected Lipschitz domains $\omega_j \subset \mathbb{R}^{d-1}$, called cross-sections, and n non-intersecting half-infinite cylinders $B_1, \dots, B_n \subset \Lambda$, isometric respectively to $\mathbb{R}_+ \times \omega_j$, $\mathbb{R}_+ := (0, +\infty)$, such that Λ coincides with the union $B_1 \cup \dots \cup B_n$ outside a compact set, see Fig. 1(a). The cylinders B_j will be called *branches*, the connected bounded domain $C := \Lambda \setminus \overline{B_1 \cup \dots \cup B_n}$ will be called *center*, and we assume that the boundary of C is Lipschitz too. We call such a domain Λ a *star waveguide*. Remark that the choice of a center is not unique: any center can be enlarged by including finite pieces of the branches, see Fig. 1(b).

In the present work, we would like to establish some elementary conditions guaranteeing the non-existence of non-trivial bounded solutions to

$$-\Delta u = \nu u \text{ in } \Lambda, \quad u = 0 \text{ at } \partial\Lambda, \tag{1}$$

E-mail address: konstantin.pankrashkin@math.u-psud.fr.

URL: <http://www.math.u-psud.fr/~pankrash/>.

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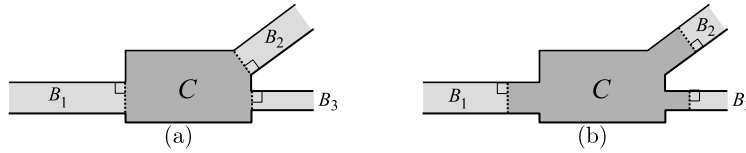


Fig. 1. (a) An example of a star waveguide Λ with three branches and a dark-shaded center. (b) An alternative choice of a center.

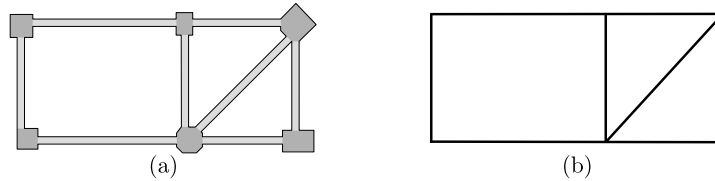


Fig. 2. (a) An example of a domain Ω_ε . The vertex parts are dark-shaded. (b) The associated one-dimensional skeleton X .

where ν is the bottom of the essential spectrum of the Dirichlet Laplacian $-\Delta_D^\Lambda$ acting in $L^2(\Lambda)$. It is standard to see that $\nu = \min \nu_j$, where ν_j is the lowest Dirichlet eigenvalue of the cross-section ω_j , and the spectrum of $-\Delta_D^\Lambda$ consists of the semi-axis $[\nu, \infty)$ and of a finite family of discrete eigenvalues $\lambda_j(-\Delta_D^\Lambda)$, $j \in \{1, \dots, N(\Lambda)\}$, while the case $N(\Lambda) = 0$ (no discrete eigenvalues) is possible. As shown e.g. in [19, Theorem 4], a non-trivial bounded solution of (1) exists iff the resolvent $z \mapsto (-\Delta_D^\Lambda - z)^{-1}$ has a pole at $z = \nu$, and in that case we say that Λ admits a threshold resonance.

The study of threshold resonances is motivated, in particular, by the analysis of Dirichlet Laplacians in systems of thin tubes collapsing onto a graph. Namely, for a small $\varepsilon > 0$, consider a domain $\Omega_\varepsilon \subset \mathbb{R}^d$ composed of finitely many cylinders (“edges”) $B_{j,\varepsilon}$ isometric to $I_j \times (\varepsilon\omega)$ with $I_j := (0, \ell_j)$, $\ell_j \in \mathbb{R}_+$, $j \in \{1, \dots, J\}$, connected to a “network” through some bounded Lipschitz domains (“vertices”) $C_{k,\varepsilon}$, see Fig. 2(a). (The case of nonidentical cross-sections is also possible but the formulations become more complicated.) We assume that the vertices $C_{k,\varepsilon}$ are isometric to εC_k , where the domains C_k are ε -independent, $k \in \{1, \dots, K\}$, and that the pieces are glued together in such a way that if one considers a vertex $C_{k,\varepsilon}$ and extends the attached edges to infinity, then one obtains a domain isometric to $\varepsilon\Lambda_k$, where Λ_k is an ε -independent star waveguide whose center is C_k .

Denote by $-\Delta_D^{\Omega_\varepsilon}$ the Dirichlet Laplacian in Ω_ε . In various applications, one is interested in the asymptotics of its eigenvalues $\lambda_m(-\Delta_D^{\Omega_\varepsilon})$ as ε tends to 0, see e.g. the monographs [11,25] and the reviews [15,17]. As the domain Ω_ε collapses onto its one-dimensional skeleton X composed from the intervals I_j coupled at the vertices, see Fig. 2(b), one may expect that the eigenvalue asymptotics might be determined by some effective operator acting on the functions defined on X . The results obtained by several authors, see e.g. [14,19], can be informally summarized as follows. Consider the star waveguides Λ_k associated to each vertex as described above, the associated Dirichlet Laplacians $-\Delta_D^{\Lambda_k}$ and their discrete eigenvalues $\lambda_j(-\Delta_D^{\Lambda_k})$, $j \in \{1, \dots, N(\Lambda_k)\}$, $k \in \{1, \dots, K\}$, then the bottom of the essential spectrum is exactly the first Dirichlet eigenvalue ν of the cross-section ω , and the following holds as ε tends to 0: there exists $N \geq N(\Lambda_1) + \dots + N(\Lambda_K)$ such that

- for $m \in \{1, \dots, N\}$ there holds

$$\lambda_m(-\Delta_D^{\Omega_\varepsilon}) = \frac{a_m}{\varepsilon^2} + \mathcal{O}(e^{-c_m/\varepsilon}) \text{ with } a_m \in (0, \nu] \text{ and } c_m > 0,$$

- for any $m \geq 1$ there holds

$$\lambda_{N+m}(-\Delta_D^{\Omega_\varepsilon}) = \frac{\nu}{\varepsilon^2} + \mu_m + \mathcal{O}(\varepsilon),$$

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