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# The resolution of the Yang–Mills Plateau problem in super-critical dimensions



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## ABSTRACT

We study the minimization problem for the Yang–Mills energy under fixed boundary connection in supercritical dimension  $n \geq 5$ . We define the natural function space  $\mathcal{A}_G$  in which to formulate this problem in analogy to the space of integral currents used for the classical Plateau problem. The space  $\mathcal{A}_G$  can be also interpreted as a space of weak connections on a “real measure theoretic version” of reflexive sheaves from complex geometry.

We prove the existence of weak solutions to the Yang–Mills Plateau problem in the space  $\mathcal{A}_G$ .

We then prove the optimal regularity result for solutions of this Plateau problem. On the way to prove this result we establish a Coulomb gauge extraction theorem for weak curvatures with small Yang–Mills density. This generalizes to the general framework of weak  $L^2$  curvatures previous works of Meyer–Rivière and Tao–Tian in which respectively a strong approximability property and an admissibility property were assumed in addition.

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## 1. Introduction

### 1.1. A nonintegrable Plateau problem

Consider a smooth compact Riemannian  $n$ -manifold  $M$  with boundary and let  $G$  be a compact connected simply connected non-abelian Lie group with Lie algebra  $\mathfrak{g}$ . We assume that a principal  $G$ -bundle  $P \rightarrow \partial M$  is fixed over the boundary of  $M$ . On  $P$  we consider a  $G$ -invariant connection  $\omega$ , which corresponds to an equivariant horizontal  $n$ -plane distribution  $Q$  (see [26] for notations and definitions).

Analogously to the Plateau problem, we may then ask which is the “most integrable” extension of  $P, Q$  to a horizontal distribution on a principal  $G$ -bundle over  $M$ . By Frobenius’ theorem, the condition for integrability in this case is that for any two horizontal  $G$ -invariant vector fields  $X, Y$ , their lie bracket  $[X, Y]$  be again horizontal. The  $L^2$ -error to integrability of an extension of  $Q$  over  $M$  can be measured by taking vertical projections  $\mathcal{V}$  of  $[X_i, X_j]$  for  $X_i, X_j$  varying in an orthonormal basis of  $Q$ :

$$\int_M \sum_{i,j} |\mathcal{V}([X_i, X_j])|^2. \quad (1.1)$$

Note that  $F(X, Y) = \mathcal{V}([X, Y])$  is known to be a tensor, and  $F$  is nothing but the curvature of the connection.

From now on we will work on the vector bundle  $E \rightarrow M$  associated to the principal bundle  $G$  corresponding to a representation of  $G$ . The covariant derivative  $\nabla$  on  $E$  is identified, in a trivialization, and via the implicit action of the representation, with the local expression

$$\nabla \stackrel{loc}{=} d + A,$$

where  $A$  is a  $\mathfrak{g}$ -valued 1-form on a given chart of  $M$ . The structure equation relating curvature to connection takes the form

$$F \stackrel{loc}{=} dA + A \wedge A \quad (1.2)$$

in a trivialization. Here  $\wedge$  represents a tensorization of the usual exterior product of forms with the Lie bracket on  $\mathfrak{g}$ . In this setting the  $L^2$ -error in integrability (1.1) is identified with the *Yang–Mills energy*, which we consider as being a functional of the connection  $\nabla$ :

$$\mathcal{YM}(\nabla) := \int_M |F_\nabla|^2. \quad (1.3)$$

We observe that, similarly to the area functional in the Plateau problem,  $\mathcal{YM}$  has a *large invariance group* given by changing coordinates in the fibers via  $G$ . The gauge group can,

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