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Third order spectral analysis robust to mixing artifacts for mapping cross-frequency interactions in EEG/MEG

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ABSTRACT

We present a novel approach to the third order spectral analysis, commonly called bispectral analysis, of electro-22 encephalographic (EEG) and magnetoencephalographic (MEG) data for studying cross-frequency functional 23 brain connectivity. The main obstacle in estimating functional connectivity from EEG and MEG measurements 24 lies in the signals being a largely unknown mixture of the activities of the underlying brain sources. This often 25 constitutes a severe confounder and heavily affects the detection of brain source interactions. To overcome this 26 problem, we previously developed metrics based on the properties of the imaginary part of coherency. Here, 27 we generalize these properties from the linear to the nonlinear case. Specifically, we propose a metric based 28 on an antisymmetric combination of cross-bispectra, which we demonstrate to be robust to mixing artifacts. 29 Moreover, our metric provides complex-valued quantities that give the opportunity to study phase relationships 30 between brain sources. 31

The effectiveness of the method is first demonstrated on simulated EEG data. The proposed approach shows a 32 reduced sensitivity to mixing artifacts when compared with a traditional bispectral metric. It also exhibits a better 33 performance in extracting phase relationships between sources than the imaginary part of the cross-spectrum for 34 delayed interactions. The method is then applied to real EEG data recorded during resting state. A cross- 35 frequency interaction is observed between brain sources at 10 Hz and 20 Hz, i.e., for alpha and beta rhythms. 36 This interaction is then projected from signal to source level by using a fit-based procedure. This approach high-37 lights a 10–20 Hz dominant interaction localized in an occipito-parieto-central network. 38

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Introduction

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Electroencephalography (EEG) and magnetoencephalography (MEG) 45 46 are noninvasive techniques which provide the opportunity to directly measure ongoing brain activity with very high temporal but relatively 47 low spatial resolution. While in the past decades the main focus of 48 EEG/MEG studies was on the analysis of event related potentials, i.e. the 49 50average brain response to a given stimulus, more recently the variability of brain activity has attracted many researchers. The recent interest in 51this field reflects the understanding that a mere localization of specific 5253brain activities is far from sufficient to understand how the brain operates, but that it is necessary to study the brain as a network. In this 54 framework, the analysis of brain rhythms has been recognized as a prom-5556ising approach since coherent neuronal activity has been hypothesized to 57serve as a mechanism for neuronal communication (Fries, 2009; Gross 58et al., 2006; Miller et al., 2009; Tallon-Baudry et al., 1996; Womelsdorf and Fries, 2006). 59

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The study of brain connectivity using noninvasive electrophysiolog- 60 ical measurements like EEG or MEG also presents some problems which 61 still need to be faced. Most notably, the fact that the data are a largely 62 unknown mixture of the activities of the actual brain sources constitutes 63 a severe confounder. For instance, two sensors can record from the 64 same neural populations, opening the possibility for spurious interac- 65 tions between sensors in the absence of true brain interactions. Though 66 the problem of mixing artifacts is well known (Nunez et al., 1997), it is 67 increasingly acknowledged and studied not only for channel data (often 68 referred to as volume conduction or field spread) (Srinivasan et al., 69 2007; Winter et al., 2007) but also at the source level, i.e., after source 70 activities have been estimated from channel data using an inverse calcu-71 lation (Schoffelen and Gross, 2009). Indeed, almost all the linear and 72 nonlinear methods used to analyze multivariate data for neuroscientific 73 applications (an excellent overview can be found in Pereda et al., 2005) 74 are highly sensitive to mixing artifacts. 75

To overcome the problem of volume conduction it was suggested to 76 exploit the fact that the propagation of electromagnetic fields is much 77 faster than neural communication: while phase shifts between electric 78 scalp potentials (EEG) or neuromagnetic fields (MEG) and the underly-79 ing source activity are too small to be observable (Stinstra and Peters, 80 1998), the temporal resolution of the data is still sufficient to capture 81

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phase shifts of neuronal signal propagation. This observation has been 82 83 exploited in the imaginary part of coherency (ImCoh), a measure of brain connectivity which cannot be caused by mixtures of independent 84 85 sources (Nolte et al., 2004). However, in the presence of interacting sources, the actual value still depends on how sources are mapped 86 into sensors (Nolte et al., 2004) or on source space (Sekihara et al., 87 2011). This drawback was addressed for two sources using pairwise 88 measures with the lagged phase coherence (Pascual-Marqui, 2007b; 89 90 Pascual-Marqui et al., 2011) or with the weighted phase lag index 91 (Vinck et al., 2011).

Nonlinear methods to estimate correlations between power 92addressing the issue of artifacts of volume conduction have also been 93 94recently suggested. In Brookes et al. (2012), the authors address the 95problem of field spread which generates spurious source space connectivity results. Using a seed based approach, the linear projection of the 96 seed voxel is first regressed out from the signals at the test voxel and, 97 then, power correlations are assessed both within and across multiple 98 frequency bands. Similarly, in Hipp et al. (2012), sensor signals are 99 orthogonalized before computing power envelope correlations at the 100 same or different frequencies, thus removing signal components that 101 share the same phase. 102

A noteworthy approach to the problem of volume conduction is 103 104 proposed in Gómez-Herrero et al. (2008). Here, after an initial principal component analysis (PCA), the authors propose to subtract a linear 105 multivariate autoregressive model from sensor data to suppress all 106 time-delayed correlations, with the idea that all neural interactions 107 require a minimum delay. An independent component analysis (ICA) 108 109 is then applied to the residuals and the ICA mixing matrix is used to model the effects of volume conduction (see also Hyvärinen et al., 110 2010). This approach takes note of the fact that a direct application 111 of ICA to the data would be a conceptual contradiction to the objec-112 113tive of the research, namely studying causality relationships between 114 sources.

In this paper, we address the problem of mixing artifacts in relation 115to the use of nonlinear methods for studying cross-frequency phase-116 synchronization between neuronal populations. Specifically we refer to 117 bispectral measures, which were developed and applied on EEG/MEG in 118 abundance (Darvas et al., 2009a, 2009b; Dumermuth et al., 1971; 119 Helbig et al., 2006; Jirsa and Müller, 2013; Schwilden, 2006; Wang et al., 1202007), and we examine the question of what information can be derived 121 from such measures that estimate true functional connectivity between 122 123 brain regions as opposed to mixing artifacts. Our new contribution is, essentially, the generalization to nonlinear methods of the concepts based 124 on the imaginary part of coherency to solve the problem of volume 125conduction (Marzetti et al., 2008; Nolte et al., 2004, 2008, 2009). As will 126 be shown below, for linear measures (e.g., cross-spectra), the imaginary 127128 part equals the antisymmetric part (apart from a factor ι , i.e., the imaginary unit) and the antisymmetry property is the more general principle 129from which measures robust to artifacts of volume conduction can be de-130rived for second-order (linear) and for third-order (nonlinear) moments. 131 In this way, antisymmetrized cross-bispectra can be used along with the 132133 imaginary part of cross-spectra for identifying phase-locked brain areas 134without being confounded by mixing artifacts, but with the important difference that the former reflects the presence of brain rhythms locked 135together at different frequencies, while the latter focuses on interactions 136at the same frequency. Moreover, the proposed approach has also the ad-137138 vantage of improving, for a certain class of interactions, a limitation of the imaginary part of the cross-spectrum, which cannot provide information 139about relative phases, i.e. the phase difference of the activities of two 140 brain sources, in a way which is robust to artifacts of volume conduction. 141 Indeed, the imaginary part of the cross-spectrum is itself a real and not a 142 complex valued quantity, and real values do not contain information 143 about relative phases. Hence, the dilemma of linear measures is the fact 144 that possibly interesting quantities cannot be estimated in a way which 145is robust to artifacts of volume conduction. On the contrary, the antisym-146 147 metric part of third order moments (cross-bispectra) is itself complex and hence contains phase information which is not corrupted by noninteract- 148 ing sources. 149

The paper is organized as follows. In the Material and methods section, we present the theory for cross-bispectral measures robust to mixing artifacts. Specifically, we first recall the basic principles of the imaginary part of cross-spectrum and, then, we introduce the antisymmetric part of the cross-bispectrum, discussing its properties with regards to mixing artifacts. We describe a strategy to project the interaction from channel to source level by using a fit-based procedure. We also discuss some examples of interpretation of the phase of crossbispectral measures. In the Result section, we first analyze the perforbispectral measures. In the Result section, we first analyze the perforulated EEG data. We then describe an example of an application of the method to real EEG data. Finally, the Discussion section provides remarks on the method features and on its ability to give an insight on cross-frequency functional connectivity.

Material and methods

Theory for cross-bispectral measures robust to mixing artifacts

Cross-spectra and mixing artifacts

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We first recall some principles of second order statistical analysis in 167 the frequency domain. The respective statistical moments, the elements 168 of the cross-spectral matrix *S*, are defined as 169

$$S_{ij}(f) = \left\langle X_i(f)X_j^*(f) \right\rangle \tag{1}$$

where $X_i(f)$ and $X_j(f)$ are the Fourier coefficients of (eventually windowed) segments of data in channel *i* and channel *j* at frequency *f*, * denotes complex conjugation, and $\langle \cdot \rangle$ denotes taking the expectation value, i.e. taking the hypothetical average over an infinite number of rsegments. Of course, the expectation value is unknown and will in genrse estimated by a finite average over segments. Since $S = S^{\dagger}$, where $(\cdot)^{\dagger}$ denotes transpose and complex conjugation, *S* is an hermitian matrix. Complex coherency, *C*, is defined as the cross-spectrum normalized by power, i.e. the diagonal elements of it:

$$C_{ij}(f) = \frac{S_{ij}(f)}{\left(S_{ii}(f)S_{jj}(f)\right)^{1/2}}.$$
(2)

It was argued that the imaginary part of the coherency is a useful 182 quantity to study brain interaction because it cannot be generated 183 from a superposition of independent sources (Nolte et al., 2004). For 184 later use we rederive this result assuming that the data have zero 185 mean which, if not vanishing, has to be subtracted from the raw data. 186 We now assume that all sources $s_k(f)$ are mapped instantaneously into 187 channels as 188

$$X_i(f) = \sum_k a_{ik} s_k(f) \tag{3}$$

with a_{ik} being real valued coefficients corresponding to the forward 190 mapping of the *k*th source to the *i*th channel. Then the cross-spectrum 191 can be written as 192

$$S_{ij}(f) = \sum_{k} a_{ik} a_{jk} \langle |s_k(f)|^2 \rangle + \sum_{k \neq k'} a_{ik} a_{jk'} \langle s_k(f) s_{k'}^*(f) \rangle.$$
(4)

If we now assume that all sources are independent, the second term 195 on the right hand side in the above equation vanishes because for $k \neq k'$ 196

$$\left\langle s_{k}(f)s_{k'}^{*}(f)\right\rangle = \left\langle s_{k}(f)\right\rangle \left\langle s_{k'}^{*}(f)\right\rangle = 0. \tag{5}$$
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Since the first term in Eq. (4) is real valued, a non-vanishing imagi- 199 nary part of *S* must arise from interacting sources and can be used to 200

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