Food Hydrocolloids 62 (2017) 128-139

Contents lists available at ScienceDirect

Food Hydrocolloids

journal homepage: www.elsevier.com/locate/foodhyd

Review

### Gastrophysics: Statistical thermodynamics of biomolecular denaturation and gelation from the Kirkwood-Buff theory towards the understanding of tofu

Seishi Shimizu <sup>a, \*</sup>, Richard Stenner <sup>a, 1</sup>, Nobuyuki Matubayasi <sup>b, c</sup>

<sup>a</sup> York Structural Biology Laboratory, Department of Chemistry, University of York, Heslington, York YO10 5DD, United Kingdom
 <sup>b</sup> Division of Chemical Engineering, Graduate School of Engineering Science, Osaka University, Toyonaka, Osaka 560-8531, Japan

<sup>c</sup> Elements Strategy Initiative for Catalysts and Batteries, Kyoto University, Katsura, Kyoto 615-8520, Japan

#### A R T I C L E I N F O

Article history: Received 9 May 2016 Received in revised form 19 July 2016 Accepted 22 July 2016 Available online 25 July 2016

Keywords: Statistical thermodynamics Kirkwood-Buff theory Hydration Cosolvents Gelation Tofu

#### ABSTRACT

Sugars, alcohols, or salts, when added to food, affects the heat denaturation of proteins and the sol-gel transition of macromolecules. Such an effect of cosolvents has long been known and exploited; yet understanding how they work at a molecular level has been a matter of scientific debate for decades, because of the lack of a definitive theory which can provide a microscopic explanation. Here we show that a rigorous statistical thermodynamic theory, the Kirkwood-Buff (KB) theory, provides not only a long-awaited microscopic explanation but also a clear guideline on how to analyze experimental data. KB theory synthesizes the classical Wyman-Tanford formula and partial molar volume, and enables the determination of biomolecule-water and biomolecule-cosolvent interactions solely from experimental data. Nothing beyond the materials in introductory physical chemistry or chemical thermodynamics textbooks is necessary to follow the derivations presented in this review.

© 2016 Elsevier Ltd. All rights reserved.

#### Contents

1.	Statistical thermodynamics for food science: why necessary?	128
2.	Thermodynamics without statistical mechanics is prone to confusion	129
3.	Clarification comes from statistical thermodynamics	130
4.	Statistical thermodynamics (the Kirkwood-Buff theory) completes and fulfils the Wyman-Tanford formula	132
5.	The Wyman-Tanford formula can be simplified when cosolvents are dilute	133
6.	The Clausius-Clapeyron equations fill the remaining theory-experiment gap	133
7.	The Kirkwood-Buff theory can readily be applied to food science: tofu as an example	133
	7.1. Heat denaturation of soy protein isolates in the presence of salts	134
	7.2. Effects of polyols on the thermal denaturation of soy proteins	135
	7.3. Effects of polyols on the gelation of tofu	135
	7.4. What the Kirkwood-Buff theory cannot do in its present form	137
8.	Conclusion	137
	Acknowledgments	137
	Acknowledgments	137-138
	References	138

## 1. Statistical thermodynamics for food science: why necessary?

Our aim is to convince the readers that statistical thermodynamics is indeed a useful tool for food science. This is especially







<sup>\*</sup> Corresponding author.

E-mail address: seishi.shimizu@york.ac.uk (S. Shimizu).

<sup>&</sup>lt;sup>1</sup> Present address: School of Physics, University of Bristol, Bristol BS8 1 TL, United Kingdom.

true, when we try to understand what is really happening at a molecular scale. Molecular-based understanding is central to food science, because it attempts to elucidate the texture and taste of food based upon its microscopic behaviour, i.e., the structure and interaction of the constituent molecules (Belitz, Grosch, & Schieberle, 2009; de Man, 1999; Walstra, 2003; Nishinari & Fang, 2016). Statistical thermodynamics, then, is indispensable, because it is the only branch of science which can provide a link between the microscopic and macroscopic worlds (Ben-Naim, 2006; Hill, 1956).

Applying statistical thermodynamics to complex systems such as food is far from being straightforward. Most commonly, two strategies have been adopted: (i) computer simulation (Barker & Grimson, 1989; Euston, Ur-Rehman, & Costello, 2007; Fundo, Quintas, & Silva, 2015) and (ii) development of simple models (van der Sman, 2016; van der Sman, Paudel, Voda, & Khalloufi, 2013). Simulations (such as molecular dynamics and Monte Carlo) implement statistical thermodynamics numerically. Simple modelbased approaches are drawn chiefly from the models of polymers, surfaces, and colloids. The crux of the both approaches lies in the elegance of approximations, aimed at grasping the essence of molecular structure and interactions out of the overwhelming complexity of food systems.

In contrast to the above, we take an alternative approach: (iii) rigorous theory as a tool to extract molecular-level information from thermodynamic data. This approach is distinct from (i) and (ii) in that certainty, credibility and clarity of interpretation are guaranteed by the rigorous nature of the theory, because the theory comes directly from the Laws of Physics (Booth, Abbott & Shimizu, 2012; Booth, Omar, Abbott & Shimizu, 2015; Shimizu, 2004; Shimizu & Abbott, 2016; Shimizu & Boon, 2004; Shimizu & Matubayasi, 2014c; Shimizu, 2015; Stenner, Matubayasi, & Shimizu, 2016).

The aim of this review is to demonstrate how useful statistical thermodynamics is. We will derive all the necessary formulae from scratch. The derivation is quite straightforward; no background knowledge is required beyond introductory chemical thermodynamics, such as Gibbs-Duhem and Clausius-Clapeyron equations (Atkins & de Paula, 2014, pp. 680–682).

### 2. Thermodynamics without statistical mechanics is prone to confusion

Our proposal to apply rigorous statistical mechanics to food science does not mean in any way that we are advocating the abolition of the current thermodynamic and calorimetric approaches. On the contrary, statistical mechanics fulfils the full potential of thermodynamic analysis, by bringing in an unprecedented interpretive clarity at a molecular level. (Such a combination of thermodynamics and statistical mechanics is commonly called statistical thermodynamics; our standpoint is to pursue food(-related) science within the framework of statistical thermodynamics.) What have instead been abolished are confusion and ambiguity caused by the lack of an explicit molecular basis (Shimizu, 2004; Shimizu & Boon, 2004; Shimizu & Matubayasi, 2014a).

**Question:** Consider the addition of extra molecular component(s) – such as sugars, salts, amino acid derivatives, or macromolecules – to food. Such an addition of cosolvents affects gelation, solubility, denaturation, and aggregation. How do cosolvents modulate such equilibria? (Note that such extra components are referred to in many different names, such as cosolvents, cosolutes, additives, or solutes; "cosolvents" will be used throughout this paper.)

A thermodynamic answer: Consider a transition  $\alpha \rightarrow \beta$  of the solute (referred to as species *u*), such as sol  $\rightarrow$  gel, solute in pure phase  $\rightarrow$  solute dissolved in solvent, folded  $\rightarrow$  unfolded, or monomers  $\rightarrow$  aggregate, and the accompanying the standard Gibbs free energy  $\Delta \mu_{u}^*$ . (Throughout this paper,  $\Delta$  signifies the change that accompanies a transition  $\alpha \rightarrow \beta$ .) The addition of cosolvents (species 2) into water (species 1) changes the water activity, and therefore the chemical potential of water  $\mu_1$ . How  $\Delta \mu_u^*$  changes with  $\mu_1$  can be expressed as a *competition* between the change in number of water and cosolvent molecules bound to the biomolecules,  $\Delta N_{u1}$  and  $\Delta N_{u2}$  (Fig. 1) that accompany the transition:

$$-\left(\frac{\partial\Delta\mu_{u}^{*}}{\partial\mu_{1}}\right)_{T,P,n_{u}\to0} = \Delta N_{u1} - \frac{n_{1}}{n_{2}}\Delta N_{u2}$$
(1)

where  $n_1$  and  $n_2$  are bulk concentrations of water and cosolvent (Casassa & Eisenberg, 1964; Parsegian, Rand, & Rau, 1995; Parsegian, Rand, & Rau, 2000; Schellman, 1987; Tanford, 1968, 1969, 1970; Timasheff, 1998, 2002a, 2002b; Wyman, 1948; 1964). For example,  $\Delta N_{u1}$  in the context of protein unfolding signifies  $\Delta N_{ui} = N_{ui}^u - N_{ui}^f$ , which corresponds to the difference in the number of bound water between the unfolded state (u) and the folded state (f).

Eq. (1) is commonly referred to as the Wyman-Tanford formula, whose interpretation owes to the seminal contributions by Wyman, Eisenberg, Tanford, Schellman, Timasheff, and Parsegian (Casassa & Eisenberg, 1964; Parsegian et al., 1995, 2000; Schellman, 1987; Tanford, 1968, 1969, 1970; Timasheff, 1998, 2002a, 2002b; Wyman, 1948; 1964; ). Eq. (1) can readily be, and has indeed been, applied to interpret thermodynamic data (Baier, Decker, & McClements, 2004; Miyawaki & Tatsuno, 2010; Miyawaki, Dozen, & Nomura, 2013; Miyawaki, Omote & Matsuhira, 2015). All one should do is to plot  $\Delta \mu_u^*$  against  $\mu_1$  (or equivalently against *RT* ln  $a_1$ where  $a_1$  is the water activity (Atkins & de Paula, 2014, pp. 680–682)) in order to obtain  $\Delta N_{u1} - \frac{n_1}{n_2} \Delta N_{u2}$  (Fig. 2). The food science applications of this formula include the effects of sugars, salts and alcohols on the thermal denaturation of proteins, as well as on gelation (Baier et al., 2004; Miyawaki & Tatsuno, 2010; Miyawaki et al., 2013; Miyawaki et al., 2015).

**Further simplification:** Which is the dominant contribution to the equilibrium shift, the change of water binding ( $\Delta N_{u1}$ ) or cosolvent binding ( $\Delta N_{u2}$ )? The Wyman-Tanford formula (Eq. (1)) on its own does not provide a definitive answer to this question. Yet dialysis measurements show that sugars, polyols, and "kosmotropic" salts are preferentially excluded from biomolecular surfaces (Fig. 1(b)) (Timasheff, 1998, 2002a, 2002b). Consequently they are not bound to biomolecules; hence the change of the bound number is zero, i.e.,  $\Delta N_{u2} = 0$  (Parsegian et al., 1995, 2000). This renders Eq. (1) a powerful tool by simplifying it to

$$-\left(\frac{\partial\Delta\mu_{u}^{*}}{\partial\mu_{1}}\right)_{T,P,n_{u}\to0} = \Delta N_{u1}$$
<sup>(2)</sup>

The change in number of bound water molecules  $\Delta N_{u1}$ , which accompany folding, gelation, solubilisation, aggregation, can therefore be measured directly from the Wyman-Tanford plot (Fig. 2) (Parsegian et al., 1995, 2000).

**Controversy:** Eq. (2) was the focus of the intense controversy, because of the unrealistically large numbers of water molecules have been estimated to be released from protein-ligand binding and allosteric transitions (Timasheff, 1998, 2002a). In the course of controversy, the following doubts have been raised about Eq. (1) (Parsegian et al., 2000; Timasheff, 1998, 2002a):

Download English Version:

# https://daneshyari.com/en/article/604007

Download Persian Version:

https://daneshyari.com/article/604007

Daneshyari.com