

## Microstructures, surface states and field emission mechanism of graphene–tin/tin oxide hybrids

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### ABSTRACT

The effects of microstructures and surface states on the field emission, which are important to a good understanding of the field emission mechanism, are unclear. In this paper, the microstructures and surface states of graphene–Sn/SnO<sub>2</sub> hybrids were analyzed, and the field emission mechanism was explored. Raman spectra and images revealed that SnO<sub>2</sub>/Sn droplets are strongly bound on graphene surface, and there exist oxygen vacancies at the surface of graphene–Sn/SnO<sub>2</sub> hybrids. Among X-ray photoelectron spectroscopy spectra, the peak of O 1s shifts 1.6 eV toward higher binding energies in the 5 min sample with the best field emission properties, which indicates that the field emission improvement in graphene–Sn/SnO<sub>2</sub> hybrids arises from the band-bending effect and a lower work function.

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### 1. Introduction

Graphene as a two-dimensional macromolecular sheet of carbon atoms has superior electrical conductivity and mechanical properties [1,2], which would make it an excellent electron transport material, even more appropriate than C<sub>60</sub>, polyaniline, graphite-like carbon [3]. Zhang et al. [4] reported that graphene has a high level of emission current at a low external electric field. Li et al. [5] reported temperature dependence of the field emission from the few-layer graphene films. Stratakis et al. [6] found that field emission properties of microstructured silicon vertices can be enhanced by depositing free-standing graphene. These results suggest that graphene might also be an attractive material for electron transport in field emission displays.

Field emission properties of SnO<sub>2</sub> have been explored extensively [7–9]. Researchers have also attempted to further improve the field emission properties of SnO<sub>2</sub>. Such as, Wang et al. [10] reported that highly oriented SnO<sub>2</sub> nanorod arrays have more excellent field emission than SnO<sub>2</sub> films. Wu [11] fabricated Sb doped SnO<sub>2</sub> nanowires, he found that the field emission of SnO<sub>2</sub> nanowires was improved by doping Sb. Jang et al. [12] concluded that

posttreatment of H<sub>2</sub> exposure can be an available process for improving the field emission properties of either SnO<sub>2</sub> nanowire or nanorod.

Graphene–SnO<sub>2</sub> hybrids have also been prepared by various methods [13–15]. Such as, Lian et al. [16] prepared graphene–SnO<sub>2</sub> hybrids by a gas–liquid interfacial synthesis approach. Zhang et al. [17] synthesized SnO<sub>2</sub> nanorods on graphene by a straightforward nanocrystal-seeds-directing hydrothermal method. Wang et al. [18] prepared graphene–SnO<sub>2</sub> hybrids involving solution-based chemical synthesis, spray drying, and annealing. Graphene–SnO<sub>2</sub> hybrids can also be formed by in situ reduction/oxidation [19]. Our previous work [20] has also shown that few-layer graphene and SnO<sub>2</sub> (G–Sn/SnO<sub>2</sub>) hybrids are good electron emitters. However, the field emission mechanisms of G–Sn/SnO<sub>2</sub> hybrids are still controversial. Furthermore, the effects of microstructures and surface states on the field emission of hybrids, which are important to a good understanding of the electron emission mechanism, are unclear. In this paper, G–Sn/SnO<sub>2</sub> hybrids were prepared on n-Si(100) substrates. The crystal structures, microstructures, optical properties, surface states of G–Sn/SnO<sub>2</sub> hybrids were analyzed by using X-ray diffraction (XRD) pattern, transmission electron microscopy (TEM) images, laser Raman spectra, X-ray photoelectron spectroscopy (XPS) spectra. The electron transport mechanism during field emission in G–Sn/SnO<sub>2</sub> hybrids is formulated based on the experimental results.

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## 2. Experimental

Schematic illustration of fabricating process of the G-Sn/SnO<sub>2</sub> hybrids is shown in Fig. 1. The detail methods are given in our previous report [20]. XRD patterns were obtained on a Rigaku D/max-2400 X-ray diffractometer with Cu K $\alpha$  radiation ( $\lambda = 0.15406$  nm). TEM images were characterized by a JEOL JEM-2100 transmission electron microscope operated at 300 kV. Laser Raman spectra were investigated using a JY LabRAM HR800 laser Raman spectrometer from 100 to 3000 cm<sup>-1</sup> at room temperature. The 514.5 nm line of the laser was used as the excitation source. XPS spectra were analyzed by a VG Scientific Corp MK-II X-ray photoelectron spectroscope.

## 3. Results and discussion

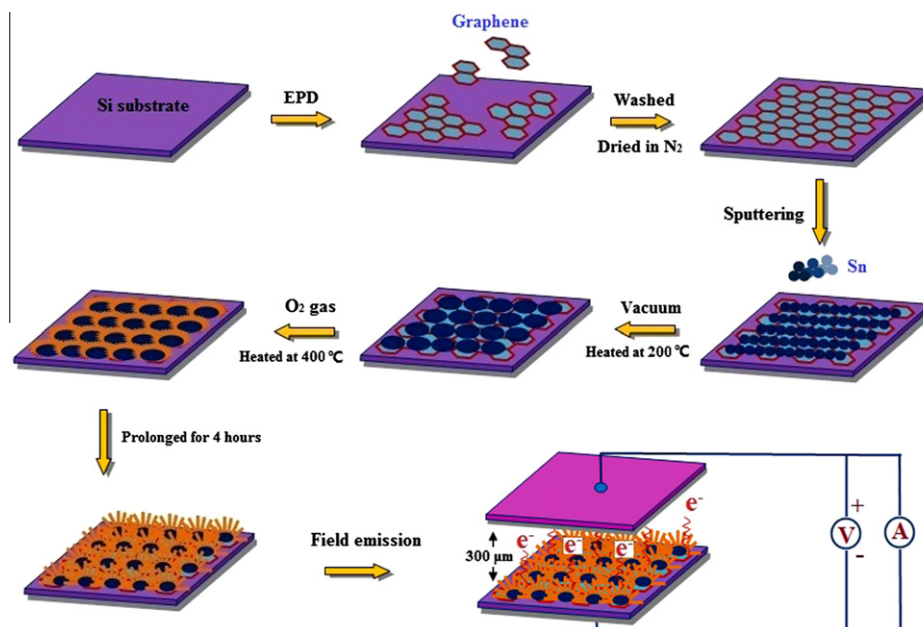
### 3.1. Microstructures of G-Sn/SnO<sub>2</sub> hybrids

Our previous work [20] reported that there are some small and compact droplets in scanning electron microscopy (SEM) image of SnO<sub>2</sub> films, while sphere flowers are formed after depositing graphene layer. In order to further investigate the microstructures of G-Sn/SnO<sub>2</sub> hybrids, it is necessary to obtain its TEM images. Fig. 2a shows TEM images of G-Sn/SnO<sub>2</sub> hybrids. One can see that transparent graphene is present over a large area, and the surfaces of the graphene are not perfectly flat, with wrinkles over the surface giving different levels of transparency. On the other hand, SnO<sub>2</sub> and/or Sn spherical droplets were attached to the surfaces of the graphene dispersedly. It should be noted that no droplets were found scattered outside graphene, even after sonication, which indicates that SnO<sub>2</sub> and/or Sn are strongly bound on graphene surface. This also confirms that G-Sn/SnO<sub>2</sub> hybrids could be successfully prepared.

Fig. 2b shows partial high-resolution TEM (HRTEM) image. The upper right inset shows distinct interplanar  $d$ -spacing of graphene, Sn and SnO<sub>2</sub>, while the lower left inset shows more general HRTEM images of G-Sn/SnO<sub>2</sub> hybrids. The measured  $d$ -spacing is 0.337 nm in accordance with graphene (002) plane. The spherical droplets

region show different interplanar  $d$ -spacing corresponding to the Sn and SnO<sub>2</sub> lattice plane, which was marked in the HRTEM images.

Fig. 3 shows Raman spectra of graphene, SnO<sub>2</sub> films, and G-Sn/SnO<sub>2</sub> hybrids. In the graphene spectrum, three sharp peaks at 1350, 1583, and 2695 cm<sup>-1</sup> correspond to the D, G, and 2D peaks, respectively. Elias et al. [21] reported that the main features in the Raman spectra of carbon based materials are the D and G peaks that lie at around 1350 and 1580 cm<sup>-1</sup>, respectively. The G peak corresponds to optical  $E_{2g}$  phonons at the Brillouin zone center, whereas the D peak is caused by breathing-like modes. The ratio of the intensity of the G-band to the D-band is related to the in-plane crystallite size [22]. Besides, the intensity ratio of  $I_G/I_D$  is widely used for characterizing the defect quantity in graphene, and a small ratio ( $I_G/I_D = 2.26$  in this work [23]) indicates a large disorder arising from structural defects [24]. The ratio of  $I_G/I_{2D}$  can be used to determine the layers number of graphene, and the small ratio ( $I_G/I_{2D} = 0.71$  in this work) shows the few number of layers. In addition, the full width at half-maximum (FWHM) of the 2D band is around 47.1 cm<sup>-1</sup>, which lies within the range of 35–50 cm<sup>-1</sup> observed for few-layered graphene systems [25]. For both SnO<sub>2</sub> films and G-Sn/SnO<sub>2</sub> hybrids, two strong peaks at 109.7 and 201.2 cm<sup>-1</sup> and a wide band at 290–720 cm<sup>-1</sup> were observed in the Raman spectra. Liu et al. reported a weak peak at 113.2 cm<sup>-1</sup> and assigned it to SnO<sub>2</sub>  $B_{1g}$  mode [26]. In addition, Diéguez et al. [27] concluded that  $B_{1g}$  mode appears quite often with smaller nanoparticles but is difficult to locate due to its very low intensity with respect to the other modes. McGuire et al. [28] confirmed that the peaks at 112 and 210 cm<sup>-1</sup> were stoichiometric composition as SnO, and the peak at 171 cm<sup>-1</sup> could be Sn<sub>2</sub>O<sub>3</sub>/Sn<sub>3</sub>O<sub>4</sub> or some other unknown phase. However, in this study, two sharp peaks at 109.7 and 201.2 cm<sup>-1</sup> with small FWHM were detected. So, the peak at 109.7 cm<sup>-1</sup> can be assigned to the SnO peak. A new Raman peak appeared at 201.2 cm<sup>-1</sup> was assigned to other tin oxides. In addition, a wide band at 290–720 cm<sup>-1</sup> containing classical  $E_g$  and  $A_{1g}$  modes was observed in both SnO<sub>2</sub> films and G-Sn/SnO<sub>2</sub> hybrids because of structural distortions induced by disorders such as local lattice disorder, oxygen deficiencies arising from Sn at the interface and surface [26]. In our G-Sn/SnO<sub>2</sub> hybrids, its XPS indicated that



**Fig. 1.** Schematic illustration of fabricating process of the G-Sn/SnO<sub>2</sub> hybrids: (i–iii) graphene was deposited on Si substrate, (iv) Sn nanoparticles were sputtered on graphene-coated Si substrate, (v) formation of aggregated Sn droplets at 200 °C for 1 h, (vi) formation of SnO<sub>2</sub> nuclei on the surface of Sn droplets at 400 °C, (vii) prolonged for 4 h to form G-Sn/SnO<sub>2</sub> hybrids, and (viii) field emission measurement.

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