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Determination of the limits of quasi-static/rigid and dynamic solutions for problems with frictional interfaces $\stackrel{\circ}{\approx}$

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ABSTRACT

Frictional interfaces exhibit complex, nonlinear behaviour, and are often sources of energy dissipation, wear, and failure mechanisms. High fidelity models of a system with frictional interfaces, however, can be computationally intensive due to the nonlinearity. Thus, numerous techniques exist that each requires different assumptions for an analysis. One categorical divide in techniques is between quasi-static and dynamic analyses. These two phenomenologically different methods are compared in order to ascertain the regimes over which each of these methods is valid. Understanding of the extent of the inertial dominated and stiffness dominated regimes offers insight into the contribution of wave propagation effects to the system's response at the frictional interface, and determines the limits of applicability of each type of analysis.

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1. Introduction

One of the greatest remaining problems in structural mechanics is the modeling of frictional interfaces and their dissipative properties. Energy dissipation and frictional contact, though, is a topic germane to many branches of engineering. In particular, energy dissipation is central to the mechanics of jointed interfaces. These commonly are found as two surfaces in frictional contact as part of a built-up structure, such as in engines [1,2], automotive brakes [3], aerospace structures [4,5] and blade-disk assemblies [6,7], or bridges [8] amongst many other applications. These interfaces exhibit complex, nonlinear behaviour, and are often sources of not only energy dissipation, but also wear and potential failure mechanisms. Consequently, it is important to develop accurate models of frictional interfaces in order to understand the stresses, dynamics, and responses of a system under various loadings.

Due to the nonlinearity introduced by the frictional interface, numerical computation of the response of a system containing interfaces can be prohibitively expensive. As a result, numerous techniques have been developed in order to create pragmatic models of built-up structures. These techniques, though, each require different sets of assumptions for an analysis. One categorical divide in

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techniques is between quasi-static analyses and dynamic analyses. Quasi-static analyses such as [5,9–12] are restricted to frequency ranges in which wave propagation effects are negligible, but allows for a high fidelity treatment of the localized stress and displacement fields. That is, it is assumed that in most practical applications, the accelerations (and consequently the inertial terms) are small enough to be neglected, which leads to the quasi-static formulation [13]. Conversely, dynamic analyses such as [14–17,7] necessitate either large, impractical simulations that model the frictional interface in high fidelity, or, in most cases, simplified simulations that regularize the application and effects of friction as occurring in a discrete manner in exchange for computational efficiency. A natural consequence of this categorical divide in modeling approaches is the set of questions:

- What effect does regularization of the contact interface have on the predicted energy dissipation rates?
- At what frequency is the quasi-static assumption of negligible acceleration terms no longer valid?

1.1. The effect of regularization on energy dissipation

From the study of interfacial mechanics, the actual mechanics of stick-slip behaviour over an interface have significant consequences for predicting dissipative properties such as the coefficient of friction [6,18]. Regularization of the contact interface, i.e. the simplification of the kinematics at a contact interface to a single degree of freedom, can lose the effects of micro-slip behaviour, which generally leads to severe under-predictions of energy





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 y_i

 γ_{in}

 η_n

μ

 ν_i

 $\rho_j \\
\Theta_j$

 θ_i

 ϕ_{jn}

 φ_j

 σ_n

 Ω

 Ω_{ς}

 $\omega_{\rm F}$

 ω_n

 au_{max} au_n

 α_{jn}, β_{jn}

 κ_{Aj}, κ_{Tj}

 κ_A, κ_T λ_n rod i

shape of rod *i*

tions of motion for rod *j*

coefficient of friction

density of rod j

shape of rod *i*

boundary excitation

excitation frequency

rod *i*

С

Poisson's ratio of rod *j*

the *n*th mode shape of rod *j*

Nomencl	ature
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A	cross-sectional area	of	rod

- c_j normalization constant for the *n*th axial mode shape of rod *j*
- E_i Young's modulus of rod j
- F_j compressive axial force applied to the boundary of rod j

F_{yn}	effective axial force due to the superposition function	
$F_{\Theta n}$	axial force due to the coupling with rotational	
	displacement	
G_j	shear modulus of rod j	
J_j	second polar moment of area of rod <i>j</i>	
Ĺ	length of rod <i>j</i>	
Ň	number of modes used in the dynamic solution	
p_0	applied axial pressure	
r	radius of both rods in the quasi-static model	
r _j	radius of rod j	
T_j	internal torque in rod <i>j</i>	
T_{Dmax}	the maximum interfacial torque for the	
	dynamic model	
T _{QSmax}	the maximum interfacial torque for the quasi-	
	static model	
T_{Un}	torque due to the coupling with axial displacement	
$T_{\varphi n}$, $T_{\varphi in}$ modal torque for the rotational displacement		
t	time	
U_j	total axial displacement of bar j	
U _i	elastic displacement of bar <i>i</i>	

 u_j elastic displacement of bar j ϖ_n natural frequency of the *n*th axial mode x_i axial position of bar j ξ_n temporal coefficient for the *n*th axial mode

 $(\cdot)^*$ denotes non-dimensional value of preceding variable

natural frequency of the *n*th torsional mode

superposition function for the axial displacement of

normalization constant for the *n*th torsional mode

coupling coefficients in the axial and torsional equa-

superposition function for the radial displacement of

normalization constant for the *n*th torsional mode

amplitude of boundary excitations with S = E, A, B, and

temporal coefficient for the *n*th torsional mode

coefficients for the torsional mode shapes

non-dimensional coupling coefficients

wave speed of the *n*th torsional mode

total rotational displacement of rod *j*

wave speed of the *n*th axial mode

elastic rotational displacement of rod *j*

maximum shear stress at the interface

dissipation rates. Further, the distribution of pressure across the contact area is non-uniform as the real contact surfaces are rough and tend to develop clusters of concentrated contact patches [13] that are the primary source of interfacial friction [19]. For instance, even when the shear force is insufficient to cause sliding, the presence of partial slip regions can lead to the bulk motion of a body [20]. A conclusion, then, of quasi-static analyses is that modeling the contact interface in a discrete manner with a single Coulomb friction model is inadequate to model accurately all of the frictional effects in a real joint [19] (even though it may well describe the global behaviour [4]) since the partial slip behaviour that is not captured accurately is fundamental to precise predictions of energy dissipation and localized behaviour.

Quasi-static simulations, though, have several limitations. The principle limiting factor associated with modeling joints is the high computation cost required for simulations [17], which leads to the conclusion that if the full description of the joint interface is included in a dynamic simulation, then the computational cost would be prohibitively expensive. Additionally, dynamic effects due to wave propagation and resonances are observed to influence significantly the energy dissipation characteristics of a joint at high frequencies [8]. Thus, a key question in the dynamic analysis of jointed structures is: how can joints be modeled accurately and efficiently without resorting to a full representation of the interfacial mechanics? There have been a number of different methods proposed over the past century to address this question [21]. Historically, these methods have been somewhat ad hoc, and generally employ tuned linear models that are valid only for a specific range of excitation frequencies. In order to develop a higher fidelity approach to modeling joints, recent research focuses on developing reduced order constitutive models to model accurately the hypothesized behaviour of a jointed interface; these models, however, still regularize the interface to be a single discrete point. One model in particular, the Iwan model [22], shows promise as a more general representation of joint friction [19]. Furthermore, the Iwan model is able to represent accurately the energy dissipation characteristics calculated from a quasi-static model of an elastic rod on a frictional foundation [1], which implies that one potential avenue for research is the further development of these models to allow for high accuracy constitutive modeling in dynamic simulations without the high computational cost of faithfully modeling the entire interface. One major constraint of this approach, though, is that the model parameters still require tuning with experimental data in order to be predictive [23], which can be a challenging task in built-up structures [2]. Thus, with regard to the question of what is the effect of regularization of the contact interface on the predictions of the system's dynamics, energy dissipation, wear, and potential failure mechanisms, the answer is dependent on the validity of the constitutive model used to represent the contact interface [24].

1.2. The frequency range of validity for quasi-static models

In the absence of a predictive, reduced order, high fidelity constitutive model of joint interfaces, analysts are left to decide between using high fidelity quasi-static models or regularized dynamic models to study the dynamics, energy dissipation, wear, and potential failure mechanisms in a built-up structure. Central to this decision is the question: at what frequency is a quasi-static assumption no longer valid? This question is addressed by studying the system of Fig. 1, in which two cylinders are arranged end-to-end, in Section 2. The quasi-static solution for the inception of slip is developed in Section 2.1. From the dynamic perspective, this system is first considered for purely torsional motion in Download English Version:

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