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Ice formation in the Arctic during summer: False-bottoms

Phan Thanh Nam^{a,b}, Pham Ngoc Dinh Alain^{a,*}, Dang Duc Trong^b, Pham Hoang Quan^b^a Mathematics Department, MAPMO UMR 6628, BP 67-59, 45067 Orleans cedex, France^b Department of Mathematics, HoChiMinh City National University, Viet Nam

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ABSTRACT

The only source of ice formation in the Arctic during summer is a layer of ice called false-bottoms between an under-ice melt pond and the underlying ocean. Of interest is to give a mathematical model in order to determine the simultaneous growth and ablation of false-bottoms, which is governed by both of heat fluxes and salt fluxes. In one dimension, this problem may be considered mathematically as a two-phase Stefan problem with two free boundaries. Our main result is to prove the existence and uniqueness of the solution from the initial condition.

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1. Introduction

There are some different reservoirs of fresh water in the Arctic during summer (see, e.g., Eicken et al. [2]). First, melt water collects in surface melt pond (melting under the sun) which is the most important reservoir. Second, this melt water can percolate into the ice matrix to form an *under-ice melt pond* (see [5] for more detail). At the interface between this fresh water and the underlying salt water, double-diffusive convection of heat and salt leads to the formation of a layer of ice called *false-bottoms* (see Fig. 1 below). Very early, Nansen [8] in 1897 noted that this is the only source of forming new ice in the Arctic during the summer. This phenomenon has been considered for a long time by many authors (see, e.g., [1,2,5–7,9–11]). However, it has been considered in geophysical view-point based on practical experiments rather than rigorously mathematical formulations.

One of the most interesting ones is the simultaneous growth and ablation of false-bottoms, which is governed by both of heat fluxes and salt fluxes. The ablation of the sea-ice interface is caused by dissolution rather than by melting. Note that salt water has the double properties: it does not freeze even for temperature less than 0 °C, and it dissolves ice when it is in contact with ice. McPhee et al. [7] emphasized that properly describing heat and salt flux at the ice-ocean interface is essential for understanding and modeling the false-bottoms, and in particular without the double diffusion at this interface false bottoms would be so short-lived. The growth of the upper interface between a false bottom and an under-ice melt pond is governed by the purely thermodynamic condition at the interface.

Recently, in 2003, Notz et al. [9] gave a model simulating successfully the simultaneous growth and ablation of false-bottoms. They formulated mathematically the problem by a system of partial differential equations and solved them numerically by using a numerical routine in Mathematica. Although their numerical result fits quite well to early experimental data from Martin and Kauffman [6], a rigorous proof of the existence and uniqueness of the solution for the system of equations is still unavailable. Our aim in this paper is to give a such a mathematical proof. More precisely, we shall represent the problem explicitly by a system of partial differential equations associated with free boundary conditions similar to [9], and then show that the system has a unique solution from given initial conditions.

* Corresponding author.

E-mail addresses: alain.pham@univ-orleans.fr, alain.pham@math.cnrs.fr (P.N. Dinh Alain).

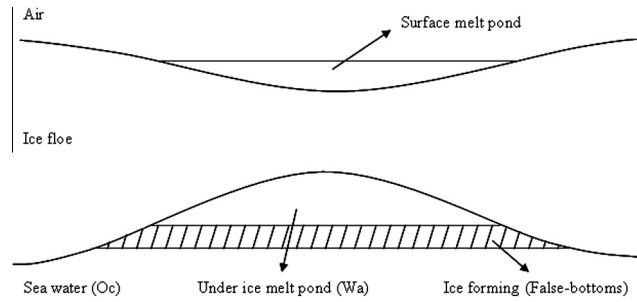


Fig. 1. Ice formation in the Arctic during summer.

Now we consider a one-dimensional model describing the simultaneous growth and ablation of the ice of false-bottoms. Here we have three environments: the ocean (Oc), the ice of false-bottoms (Fb) and the fresh water (Wa). Denote by $T(x, t), S(x, t)$ the temperature and the salinity, and denote by $h_0(t), h_u(t)$ the free boundaries at the interfaces ice-ocean (Fb–Oc) and ice-water (Fb–Wa), respectively (see Fig. 2 below).

At the interface Fb–Oc, we apply the first principle of thermodynamics, (i.e. variation of energy = variation of heat flux through the interface)

$$\Delta U = \Delta \Phi.$$

The net amount of heat transferred through the interface Fb–Oc in a section s is equal to

$$\Delta U = -dh_0 \rho_i L_f s,$$

where ρ_i is the density of the ice and L_f is the latent heat of fusion. On the other hand, the difference of the heat fluxes through a section s in the ice and the ocean during a time dt is

$$\Delta \Phi = (-\lambda_i T_x(h_0(t)+, t) + \lambda_o T_x(h_0(t)-, t))s dt,$$

where λ_i, λ_o are thermal conductivities of the ice and the ocean. Here the notations $h_0(t)+$ and $h_0(t)-$ stand for the right limit and the left limit at $x = h_0(t)$. Thus the law of conservation of energy mentioned above, i.e. $\Delta U = \Delta \Phi$, leads to the Stefan condition for the heat balance at the interface

$$h'_0(t) = \tilde{\lambda}_i T_x(h_0(t)+, t) - \tilde{\lambda}_o T_x(h_0(t)-, t), \tag{1.1}$$

where

$$\tilde{\lambda}_i = \frac{\lambda_i}{\rho_i L_f} > 0, \quad \tilde{\lambda}_o = \frac{\lambda_o}{\rho_i L_f} > 0.$$

For simplicity, we can neglect the salt of the ice of false-bottoms. The water near the interface Fb–Oc is a mixture of melt water, which melts from the ice of false-bottoms, and sea water. This water freshens at the rate $S_0(t)h'_0(t)$, while salt diffuses into this water at the rate $-DS_x(h_0(t)-, t)$, where $S_0(t) = S(h_0(t)-, t)$ is the salinity of the ocean at the interface and $D > 0$ is the molecular diffusivity of salt in sea water. The balance of salt at this interface leads to the conservation condition

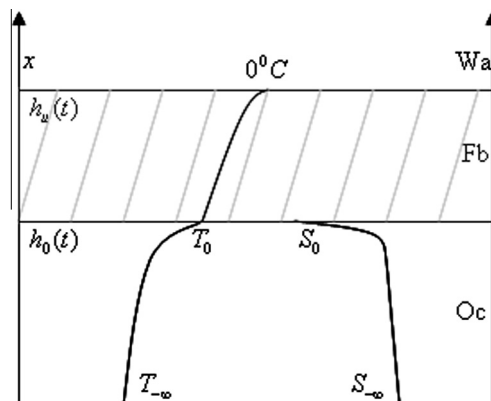


Fig. 2. One-dimensional model.

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