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Simulation of turbulent natural convection in a porous cylindrical annulus using a macroscopic two-equation model

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Abstract

This work presents numerical computations for laminar and turbulent natural convection within a horizontal cylindrical annulus filled with a fluid saturated porous medium. Computations covered the range $25 < Ra_{\rm m} < 500$ and $3.2 \times 10^{-4} > Da > 3.2 \times 10^{-6}$ and made use of the finite volume method. The inner and outer walls are maintained at constant but different temperatures. The macroscopic k- ϵ turbulence model with wall function is used to handle turbulent flows in porous media. First, the turbulence model is switched off and the laminar branch of the solution is found when increasing the Rayleigh number, $Ra_{\rm m}$. Subsequently, the turbulence model is included and calculations start at high $Ra_{\rm m}$, merging to the laminar branch for a reducing $Ra_{\rm m}$. This convergence of results as $Ra_{\rm m}$ decreases can be seen as an estimate of the so-called laminarization phenomenon. Here, a critical Rayleigh number was not identified and results indicated that when the porosity, Prandtl number, conductivity ratio between the fluid and the solid matrix and $Ra_{\rm m}$ are kept fixed, the lower the Darcy number, the higher is the difference of the average Nusselt number given by the laminar and turbulent models. © 2006 Elsevier Ltd. All rights reserved.

Keywords: Turbulence modeling; Porous media; Natural convection; Cylindrical annuli

1. Introduction

The analyses of natural convection in a horizontal cylindrical annuli filled by a porous material has been subject of a number of studies in recent years. Thermal insulators, cryogenics, thermal storage systems, electronic cooling, inert gas insulation of high-voltage electric cables and the determination of the requirements for aircraft cabin insulation are some examples of applications.

The first basic study on natural convection in clear cylindrical annuli was carried out by [1] and extended by [2]. A very extensive analysis has been made on the concentric annuli by [3]. They have conducted both numerical simulation using the finite elements technique and experimental studies using Mach-Zehnder interferometer. Application of other type of finite differences method with ADI numer-

ical solution has also reported by [4] in solving the laminar horizontal concentric annuli problem formulated in cylindrical polar coordinates. Small eccentric annuli has performed in the work of [5], using an expansion in terms of the double series of eccentricity and Rayleigh number for small values of *Ra*. The work of [6] extended the knowledge on the natural convection heat transfer in the horizontal cylindrical annuli and the numerical analysis has been made using finite difference method in a bipolar coordinate system based on successive over-relaxation iteration.

Flow analysis on clear concentric annulus has been extended to turbulent natural convection using a two-equation turbulence model in the work of [7]. Experimental works yielded new information about the field variable, primarily on the distribution of the mean temperature in the fluid and of Nusselt number around the cylinders circumference in the work of [8]. The turbulence appears in the upper region of the annulus in the plume above the heated inner cylinder. On the other hand, the stable stratification below

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Nomenclature			
$c_{ m F}$ c 's	Forchheimer coefficient non-dimensional turbulence model constants	$Re_{\mathfrak{p}}$	$Re_{\rm p} = \frac{ \mathbf{u}_{\rm D} d}{v_{\rm f}}$, Reynolds number based on the pore diameter
$c_{\rm p}$	specific heat	T	temperature
d	pore diameter	u	microscopic velocity
D	$\mathbf{D} = [\nabla \mathbf{u} + (\nabla \mathbf{u})^{\mathrm{T}}]/2$, deformation rate tensor	\mathbf{u}_{D}	Darcy or superficial velocity (volume average of u)
Da	Darcy number, $Da = \frac{K}{r^2}$		
D_{p}	particle diameter	Greek symbols	
g	gravity acceleration vector	α	thermal diffusivity
G^{i}	generation rate of $\langle k \rangle^{i}$ due to the action of the	β	thermal expansion coefficient
	porous matrix	ΔV	representative elementary volume
$G^{ m i}_{eta}$	generation rate of $\langle k \rangle^{i}$ due to the buoyant effects	$\Delta V_{ m f}$	fluid volume inside ΔV
h	heat transfer coefficient	3	$\varepsilon = \mu \nabla \mathbf{u}' : (\nabla \mathbf{u}')^{\mathrm{T}} / \rho$, dissipation rate of k
I	unit tensor	μ	dynamic viscosity
K	$K = \frac{D_{\rm p}^2 \phi^3}{144(1-\phi)^2}$, permeability	μ_{t}	microscopic turbulent viscosity
k	$k = \overline{\mathbf{u}' \cdot \mathbf{u}'}/2$, turbulent kinetic energy per mass	μ_{t_ϕ}	macroscopic turbulent viscosity
, r	$k = \mathbf{u} \cdot \mathbf{u}/2$, turbulent kinetic energy per mass unit	v	kinematic viscosity
$k_{ m f}$	fluid thermal conductivity	ho	density
$k_{\rm s}$	solid thermal conductivity	σ 's	non-dimensional constants
$\mathbf{K}_{\mathrm{disp}}$	conductivity tensor due to the dispersion	ϕ	$\phi = \Delta V_{\rm f}/\Delta V$, porosity
$\mathbf{K}_{\mathrm{disp,t}}$	conductivity tensor due to the turbulent disper-	θ	angle
uisp,t	sion	<i>c</i> .	7 7
\mathbf{K}_{t}	conductivity tensor due to the turbulent heat	•	l characters
(flux	$_{-}^{arphi}$	general scalar variable
$\mathbf{K}_{\mathrm{tor}}$	conductivity tensor due to the tortuosity	$ar{\phi}$	time average
Nu	Nusselt number	$arphi' \ \langle arphi angle^{ m i}$	time fluctuation
Pe	Peclet number	$\langle \varphi \rangle$	intrinsic average volume average
Pe_{D}	modified Peclet number, $Pe_D = Pe(1 - \phi)^{1/2}$	$\langle \varphi \rangle^{\mathrm{v}}$ $\langle \varphi \rangle^{\mathrm{v}}$	spatial deviation
P^{i}	production rate of $\langle k \rangle^{i}$ due to gradients of $\bar{\mathbf{u}}_{\mathrm{D}}$	arphi	absolute value (Abs)
Pr	v_f/α_f , Prandtl number	$\mathbf{\phi}$	general vector variable
r	radius	$arphi_{ m eff}$	effective value, $\varphi_{\rm eff} = \varphi \varphi_{\rm f} + (1 - \varphi) \varphi_{\rm s}$
R	r_0/r_i	$\varphi_{ m s,f}$	solid/fluid
$Ra_{\rm f}$	$Ra_{\rm f} = \frac{g\beta r_{\rm i}^3 \Delta T}{v_{\rm f} \alpha_{\rm f}}$, fluid Rayleigh number	$\varphi_{i,o}$	inner/outer
$Ra_{\rm m}$	$Ra_{ m m}=Ra_{ m f}\cdot Da=rac{geta_{\phi}r_{ m i}\Delta TK}{v_{ m f}lpha_{ m eff}},$ Darcy–Rayleigh	$\varphi_{\mathrm{H,C}}$	hot/cold
	number		macroscopic value
Ra_{cr}	critical Rayleigh number	$()^{\mathrm{T}}$	transpose

the heated inner cylinder suppresses any movement and turbulence that might be convected downward. Direct numerical simulation (DNS) was performed by [9]. However, because of the extreme requirements for computing resources, these computations are limited to low Ra_f numbers. The work of [10] use a Large Eddy Simulation (LES) for a case with a high Ra_f number indicated that this method may serve in the future as an alternative to the (DNS), but there are many critics about the use of such model in 2-D cases due the 3-D characteristic of the turbulence. The work of [11], reports on the modeling and computational study of the natural convection in concentric and eccentric annuli by means of several variants of the Algebraic Stress Models (ASM), which is based on the expression for turbulent heat flux $\overline{\theta u_i}$ obtained by trunca-

tion of the second-moment transport equation for this correlation. Various levels of closure included the Low-Re number form of the k- ε model, but also a version in which the differential transport equations are solved for the temperature variance θ^2 and its decay rate ε_θ . Ascending flows of liquid metal [12] and anisotropic effects in flows over rod-bundles, [13] have also been investigated via ASMs.

The natural convection in cylindrical annular geometry filled with porous material also have been studied by distinct numerical approaches, such as the finite-difference method reported by [14 and 15]. Finite element method is found in the work of [16] and the Galerkin spectral method in the work of [17–20]. The work of [21] have shown that the Fourier–Chebyshev method gives better accuracy than

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