



Effect of thermal radiation on unsteady mixed convection flow and heat transfer over a porous stretching surface in porous medium

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ABSTRACT

An analysis is performed to investigate the effects of thermal radiation on unsteady boundary layer mixed convection heat transfer problem from a vertical porous stretching surface embedded in porous medium. The fluid is assumed to be viscous and incompressible. Numerical computations are carried out for different values of the parameters involved in this study and the analysis of the results obtained shows that the flow field is influenced appreciably by the unsteadiness parameter, mixed convection parameter, parameter of the porous medium and thermal radiation and suction at wall surface. With increasing values of the unsteadiness parameter, fluid velocity and temperature are found to decrease in both cases of porous and non-porous media. Fluid velocity decreases due to increasing values of the parameter of the porous medium resulting an increase in the temperature field in steady as well as unsteady case.

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1. Introduction

The heat, mass and momentum transfer in the laminar boundary layer flow on a stretching sheet are important from theoretical as well as practical point of view because of their wider applications to polymer technology and metallurgy. The thermal buoyancy force arising due to the heating of stretching surface, under some circumstances, may alter significantly the flow and thermal fields and thereby the heat transfer behaviour in the manufacturing processes. Keeping this fact in mind, Lin et al. [1], Chen [2], Ali and Al-Yousef [3] etc. investigated the flow problems considering the buoyancy force.

Simultaneous heat and mass transfer from different geometries embedded in porous media has many engineering and geophysical applications such as geothermal reservoirs, drying of porous solids, thermal insulation, enhanced oil recovery, packed-bed catalytic reactors, cooling of nuclear reactors, and under ground energy transport. A very significant area of research in radiative heat transfer, at the present time is the numerical simulation of combined radiation and convection/conduction transport processes (Kandasamy et al. [4]). We know that the radiation effect is important under many non-isothermal situations. If the entire system involving the polymer extrusion process is placed in a thermally controlled environment, then radiation could become important. The knowledge of radiation heat transfer in the system can perhaps lead to a desired product with sought characteristic.

All of the above mentioned studies consider the steady-state problem. But, in certain practical problems, the motion of the stretched surface may start impulsively from rest. In these problems, the transient or unsteady aspects become more interesting. Recently, Elbashbeshy and Bazid [5] presented an exact similarity solution for unsteady momentum and heat transfer flow whose motion is caused solely by the linear stretching of a horizontal stretching surface. Since no attempt has been made to analyse the effects of thermal radiation on heat and mass transfer on unsteady boundary layer mixed convection flow over a vertical stretching surface in porous medium in presence of suction, this problem is investigated in this article. The momentum and the thermal boundary layer equations are solved using shooting method and the numerical calculations were carried out for different values of parameters of the problem under consideration for the purpose of illustrating the results graphically. The analysis of the results obtained shows that the flow field is influenced appreciably by the presence of unsteadiness, heat radiation, mixed convection and suction on the wall in presence of porous medium. To reveal the tendency of the solutions, representative results are presented for the velocity, temperature as well as the skin friction and rate of heat transfer. Comparisons with previously published works are performed and excellent agreement between the results is obtained.

2. Equations of motion

We consider the two-dimensional mixed convection boundary-layer flow of an incompressible viscous liquid through porous

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Nomenclature

c	constant
c_p	specific heat at constant pressure
$Da_x = \frac{k_1(1-\alpha t)}{x^2}$	local Darcy number
f	non-dimensional stream function
f', f'', f'''	first order, second order, third order derivatives respectively with respect to η
$Gr_x = \frac{g\beta(T_w - T_\infty)x^3}{\nu^3}$	Grashof number
g	gravity field
k	permeability of the porous medium
k^*	absorption coefficient
$M = \frac{\alpha}{c}$	unsteadiness parameter
$N = \frac{\kappa k^*}{4\sigma T_\infty^3}$	radiation parameter
Pr	Prandtl number
p, q	variables
q_r	radiative heat flux
$Re_x = \frac{u_w x}{\nu}$	local Reynold's number
$S(>0)$	suction parameter
T	temperature of the fluid
T_w	temperature of the wall of the surface

T_∞	free-stream temperature
u, v	components of velocity in the x and y directions
$v_0(>0)$	velocity of suction of the fluid
z	variable

Greek symbols

α	constant
β	volumetric coefficient of thermal expansion
η	similarity variable
κ	coefficient of thermal diffusivity
$\lambda = \frac{Gr_x}{Re_x^2}$	mixed convection parameter
μ	dynamic viscosity
ν	kinematic viscosity
ψ	stream function
ρ	density of the fluid
σ	Stefan-Boltzman constant
θ	non-dimensional temperature
θ', θ''	first order, second order derivatives respectively with respect to η

medium along a permeable vertical wall stretching with velocity $u_w = \frac{cx}{1-\alpha t}$ and with temperature distribution $T_w = T_\infty + 1/2T_0Re_x^{-1}(1-\alpha t)^{-1} = T_\infty + T_0\frac{cx}{2\nu}(1-\alpha t)^{-2}$ (Andersson et al. [6]) where $Re_x = \frac{u_w x}{\nu}$ is the local Reynold's number. The x -axis is directed along the stretching surface and points in the direction of motion. The y -axis is perpendicular to it. u, v are the velocity components in the x - and y -directions. The governing equations under boundary layer and Boussinesq approximations for flow through a porous medium over the stretching surface are, in the usual notation may be written as

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \quad (1)$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} + g\beta(T - T_\infty) - \frac{\nu}{k}u, \quad (2)$$

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{\kappa}{\rho c_p} \frac{\partial^2 T}{\partial y^2} - \frac{1}{\rho c_p} \frac{\partial q_r}{\partial y} \quad (3)$$

along with the boundary conditions

$$u = u_w = \frac{cx}{1-\alpha t}, \quad v = v_w = -\frac{v_0}{(1-\alpha t)^{\frac{1}{2}}}, \quad T = T_w \quad \text{at } y = 0, \quad (4)$$

$$u \rightarrow 0, T \rightarrow T_\infty \quad \text{as } y \rightarrow \infty. \quad (5)$$

Here $k[=k_1(1-\alpha t)]$ is the permeability of the porous medium, k_1 is the initial permeability, μ is the coefficient of fluid viscosity, ρ is the fluid density, $\nu = \mu/\rho$ is the kinematic viscosity, β is the volumetric coefficient of thermal expansion, g is the gravity field, T is the temperature, κ is the coefficient of thermal conductivity of the fluid, $v_0(>0)$ is the velocity of suction of the fluid, $c(>0)$ and $\alpha(>0)$ are constants with dimension $(\text{time})^{-1}$, T_w is the uniform wall temperature, T_∞ is the free-stream temperature, c_p is the specific heat at constant pressure and q_r is the radiative heat flux. The viscous dissipative term in the energy equation is neglected here.

Using Rosseland approximation, we get $q_r = -\frac{4\sigma}{3k^*} \frac{\partial T^4}{\partial y}$ where σ is the Stefan-Boltzman constant, k^* is the absorption coefficient. We assume that the temperature difference within the flow is such that T^4 may be expanded in a Taylor's series. Expanding T^4 about T_∞ and neglecting higher orders we get, $T^4 = 4T_\infty^3 T - 3T_\infty^4$.

Now Eq. (3) becomes

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \left(\frac{\kappa}{\rho c_p} + \frac{16\sigma T_\infty^3}{3\rho c_p k^*} \right) \frac{\partial^2 T}{\partial y^2}. \quad (6)$$

2.1. Method of solution

We now introduce the following relations for u, v and θ as

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x} \quad \text{and} \quad \theta = \frac{T - T_\infty}{T_w - T_\infty} \quad (7)$$

where ψ is the stream function.

Using the relation (7) in the boundary layer Eq. (2) and in the energy Eq. (6) we get the following equations

$$\frac{\partial^2 \psi}{\partial t \partial y} + \frac{\partial \psi}{\partial y} \frac{\partial^2 \psi}{\partial x \partial y} - \frac{\partial \psi}{\partial x} \frac{\partial^2 \psi}{\partial y^2} = \nu \frac{\partial^3 \psi}{\partial y^3} + g\beta(T_w - T_\infty)\theta - \frac{\nu}{k} \frac{\partial \psi}{\partial y} \quad (8)$$

and

$$\frac{\partial T}{\partial t} + \frac{\partial \psi}{\partial y} \frac{\partial T}{\partial x} - \frac{\partial \psi}{\partial x} \frac{\partial T}{\partial y} = \left(\frac{\kappa}{\rho c_p} + \frac{16\sigma T_\infty^3}{3\rho c_p k^*} \right) \frac{\partial^2 T}{\partial y^2}. \quad (9)$$

We now introduce the similarity variable η and the dimensionless variables f and θ as follows:

$$\eta = \left(\frac{c}{\nu(1-\alpha t)} \right)^{\frac{1}{2}} y, \quad \psi = \left(\frac{\nu c}{(1-\alpha t)} \right)^{\frac{1}{2}} x f(\eta), \quad T = T_\infty + T_0 \frac{cx}{2\nu} (1-\alpha t)^{-2} \theta(\eta). \quad (10)$$

In view of the relations (10), the Eqs. (8) and (9) become

$$M \left(\frac{\eta}{2} f'' + f' \right) + f'^2 - f f'' = f''' + \lambda \theta - \frac{1}{D} f', \quad (11)$$

$$\frac{M}{2} (\eta \theta' + 4\theta) + f' \theta - f \theta' = \frac{1}{Pr} \left(1 + \frac{4}{3N} \right) \theta'', \quad (12)$$

where $\lambda = \frac{g\beta T_0}{2\nu c} = \frac{Gr_x}{Re_x^2}$ is the mixed convection parameter, $Gr_x = \frac{g\beta(T_w - T_\infty)x^3}{\nu^3}$ is the Grashof number, $D = Da_x Re_x = \frac{k_1 c}{\nu}$, $Da_x = \frac{k}{x^2} = \frac{k_1(1-\alpha t)}{x^2}$ is the local Darcy number, $M = \frac{\alpha}{c}$ is the unsteadiness parameter, $N = \frac{\kappa k^*}{4\sigma T_\infty^3}$ is the radiation parameter.

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