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Onset of natural convection of electrolyte on horizontal electrode under non-steady-state mass-transfer conditions

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Abstract

In this work, the problem of convective stability of electrochemical system (a horizontal electrode and a solution containing three types of ions) under the non-steady-state mass-transfer conditions is solved. A set of equations involving the Navier–Stokes equation for incompressible viscous liquid to the Boussinesq approximation, the ion-transfer equations that account for diffusion, migration, and convection, and the electroneutrality condition were used as the mathematical model. The problem is solved within the framework of linear theory of stability. The set of amplitude equations obtained for vertical component of hydrodynamic velocity and concentrations was solved numerically by the shooting method. As a result, the critical Rayleigh number, the wave number, and the critical time of onset of natural convection were determined. It is shown that the supporting electrolyte has a strong effect on the convective stability of the system.

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Keywords: Natural-convective instability; Electrochemical system; Onset time; Non-steady-state mass-transfer; Supporting electrolyte; Migration

1. Introduction

Natural convection of electrolyte, which arises in the electrochemical systems due to the variations in the electrolyte density, has a pronounced effect on the electrode reaction rate [1,2]. In the case of horizontal electrodes, natural convection of electrolyte arises only at the Rayleigh numbers Ra higher than a certain critical value Ra_c . The determination of Ra_c involves a study of stability of stagnant electrolyte under the steady-state mass-transfer conditions. In real systems, stability of stagnant electrolyte is disturbed before the steady-state distribution of concentrations is reached. Therefore, in recent years, non-steady-state convective instability has been studied. In this case, instability arises when the diffusion layers exist near each of electrodes and the distribution of concentrations of ions in the interelectrode space has a complex, non-linear character. This

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significantly complicates the theoretical analysis of the process. In this case, convective instability is characterized by a critical time $t_{\rm c}$, i.e. a period of time from the beginning of experiment (for example, from the instant of imposing a voltage between the electrodes) to the onset of convection. The critical time can be determined by the analysis of stability of interrelated equations of liquid motion and ion-transfer equations. In [3-8] the problem of stability of stagnant fluid under the non-steady-state heat-transfer conditions was studied. The results of these studies can be used directly for the electrochemical systems with a binary electrolyte. In the multi-component electrochemical systems (with three or more types of ions), the distribution of electrolyte density over the diffusion layers depends both on electroactive and supporting electrolyte; this prevents the use of results, which were obtained for heat systems. The natural convection in the multi-component systems with horizontal electrodes was studied experimentally in several works [9-11]. Recently, the effect of stable thermal stratification on the onset of double-diffusive convection in the

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Nomenclature

- а parameter
- parameter, $\overline{Ra_c^{2/3}}$ A
- c_m concentration of ions of *m*th type
- dimensionless concentration of electroactive C_1 ions, $\frac{c_1}{2c_{1_h}}$
- dimensionless concentration of fictitious ions, C_4 $\frac{z_3c_4}{2\alpha c_{1_{\mathbf{b}}}(z_3-z_2)}\frac{\partial\rho}{\partial c_2}$

 \widehat{C}_m modified dimensionless small disturbance of concentration of ions of *m*th type, \tilde{C}_m/Ra

- D_m diffusion coefficient of ions of *m*th type
- diffusion coefficient of supporting electrolyte, D_4 $D_2 D_3 (z_3 - z_2)$ $z_3D_3 - z_2D_2$
- $\frac{D_1}{D_2} \left[\frac{n_2}{n_1} D_1 \right]$ D^* mutual diffusion coefficient, $\frac{D_3(z_3-z_2) + (D_3-D_2)z_1}{z_3D_3 - z_2D_2} D_2 - \frac{n_2 - n_1}{n_1} D_4 \end{bmatrix}$
- dimensionless diffusion coefficient of supporting \overline{D}_{4} electrolyte, $\frac{D_4}{D_1}$
- \overline{D}^* dimensionless mutual diffusion coefficient, $\frac{z_3 D^*}{\alpha D_1(z_3-z_2)} \frac{\partial \rho}{\partial c_2}$
- F Faraday number
- gravitational acceleration g
- Grashof number, $\frac{2gH^3c_{1_b}\alpha}{\rho_bv^2}$ distance between electrodes Gr
- Η
- k_X, k_Y projections of wave vector onto the coordinate axis X and Y, respectively
- wave number, $\sqrt{k_X^2 + k_Y^2}$ modified wave number, $k\sqrt{\tau/Sc_1}$ k
- k
- M_m designation of ions of *m*th type number of electrons exchanged in the electrode n
- reaction, $(z_1n_1 z_2n_2)$ stoichiometric coefficient of ions of *m*th type n_m
- pressure р
- dimensionless pressure, $\frac{H^2}{\rho_{\rm v} v D_{\rm v}} p$ Р
- R gas constant
- Ra_1 Rayleigh number of electroactive ion, $Sc_1 \cdot Gr$
- Rayleigh number of electrolyte with three types Ra of ions, $\frac{\beta}{2}Ra_1$
- Ra modified Rayleigh number of electrolyte with three types of ions, $Ra \cdot (\tau/Sc_1)^{3/2}$

electrochemical systems has been studied theoretically by Kim et al. [12]; however, the migration of supporting electrolyte ions has been ignored.

2. Formulation of the problem

Within the framework of the Boussinesq approximation, the theory of dilute electrolytes, and under the assumption of electroneutrality, equations of flow of incompressible viscous liquid and the ion transfer in the electrolyte layer between two horizontal electrodes can be written as follows [1,2,13,14]:

- Sc_1
- Schmidt number, $\frac{v}{D_1}$ amplitude of small perturbation of concentra- $S_1(\varsigma)$ tion of electroactive ions
- amplitude of small perturbation of concentra- $S_4(\varsigma)$ tion of fictitious ions
- $W_Z(\varsigma)$ amplitude of small perturbation of vertical component of hydrodynamic velocity
- time t
- τ dimensionless time, $\frac{v}{H^2}t$
- Т temperature
- hydrodynamic velocity v
- dimensionless hydrodynamic velocity, $\frac{H}{D_1}$ v v
- axial, spanwise and normal coordinate x, y, z
- X, Y, Z dimensionless axial, spanwise and normal coordinate, $\frac{x}{H}$, $\frac{y}{H}$, $\frac{z}{H}$ charge number of ions of *m*th type
- Z_m

Greek symbols

α	densification coefficient of electrolyte with three	
	types of ions. $\frac{\partial \rho}{\partial t} + \frac{z_1 - z_3 n_2 D_1 / n_1 D_2}{2} \frac{\partial \rho}{\partial t}$	
	$O_1 = Z_3 - Z_2 = O_2$	
ß	dimensionless parameter. $\frac{1+\sqrt{D_4}}{2}$	
Р	$1+\sqrt{D_4+D^*}$	
\varDelta_2	two-dimensional Laplacian, $\frac{\partial^2}{\partial X^2} + \frac{\partial^2}{\partial Y^2}$	
3	relative concentration of electroactive ions, $\frac{c_{1_b}}{c_3}$	
ϕ	electric potential	
γ	unity vector directed vertically downwards	
v	kinematic viscosity of electrolyte	
ρ	density of electrolyte	
$\frac{\partial \rho}{\partial a}$	densification coefficient of ions of <i>m</i> th type	
C^{m}	self-similar variable $(S_{C_1}/\tau)^{1/2}Z$	
7		
<u> </u>		
Subscripts		

- b bulk state
- critical conditions с

Overbars

	undisturbed state
\sim	small disturbance
、 、	modified small disturbance

$$\begin{aligned} \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} &= -\frac{1}{\rho_{\rm b}} \nabla p + v \Delta \mathbf{v} + \frac{\gamma g}{\rho_{\rm b}} (\rho - \rho_{\rm b}), \\ \operatorname{div}(\mathbf{v}) &= 0, \\ \frac{\partial c_1}{\partial t} &= D_1 \Delta c_1 + \frac{F z_1 D_1}{RT} (\nabla c_1 \nabla \varphi + c_1 \Delta \varphi) - \mathbf{v} \nabla c_1, \\ \frac{\partial c_2}{\partial t} &= D_2 \Delta c_2 + \frac{F z_2 D_2}{RT} (\nabla c_2 \nabla \varphi + c_2 \Delta \varphi) - \mathbf{v} \nabla c_2, \\ \frac{\partial c_3}{\partial t} &= D_3 \Delta c_3 + \frac{F z_3 D_3}{RT} (\nabla c_3 \nabla \varphi + c_3 \Delta \varphi) - \mathbf{v} \nabla c_3, \\ z_1 c_1 + z_2 c_2 + z_3 c_3 &= 0. \end{aligned}$$
(1)

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