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Efficiency of the lumped parameter concept and the role of liquid properties in modelling micro droplet evaporation

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HIGHLIGHTS

- Evaporation study of alkane droplets <100 µm in increasing thermal load is done.
- Effective conductivity model variants are developed to predict droplet behaviour.
- Infinite conductivity model proves high efficiency in a number of conditions.
- Except for density, variations of liquid properties are unimportant in this work.
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ABSTRACT

This paper presents a study of the micro-sized droplets of dodecane, decane and nonane undergoing evaporation under increasing thermal load through a laminar flow reactor. The objective is to use the effective conductivity concept for the liquid phase to explore the droplet internal state during its lifetime and to assess the efficiency of the infinite conductivity (lumped parameter) concept by comparing these concepts with each other and against the experimental data. Another objective of this work is to assess the sensitivity of the process on the liquid phase properties by developing the constant properties effective conductivity model to more enhanced versions and comparing them to the fully numerical one. Extension to higher Reynolds number and higher temperature conditions is also made. While the enhanced models give almost identical predictions to the full numerical version for all conditions examined, the lump model also found to yields close predictions. Examination of the liquid properties shows that relative to the other properties, accounting for variation in the liquid density is most important. © 2015 Published by Elsevier Ltd.

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1. Introduction 49

50 This paper examines the ability of a variety of approaches to model evaporation of micro-sized droplets (<100 µm). This is 51 prompted by moves toward reduced fuel droplet size as a means 52 to improve efficiency and reduce pollutant formation in many 53 combustion systems. One of the problems in undertaking such an 54 investigation is the limited experimental data available (e.g. 55 [1–5]), with even fewer studies providing high temporal and spa-56 tial resolution [6,7]). The difficulty in obtaining experimental data 57 for the length and time scales necessary under the high tempera-58 59 ture and pressure conditions found in combustion systems is the 60 main reason behind the absence of a large body of data. Of the 61 studies available, we have selected one previously published by 62 the authors for the limited condition of droplets undergoing evaporation in an environment of increasing thermal heating load,

In this series of studies, fuel droplet evaporation is investigated 70 for droplets flowing through a laminar flow quartz glass reactor, 71 described in [6,7]. The reactor's increasing convective heating load 72 is created by the mixing of the cooler carrier gas used to transport 73 droplets into the reactor with the flow produced by the reactor's 74 flat flame burner. The advantage of this approach over other 75 methods is that it provides precise control of the droplet conditions 76 (initial size, temperature and velocity), and once created, the droplet contacts only gases, rather than filaments or similar droplet suspension devices. The setup is characterised by a prolonged preheating time that the droplet experiences as the carrier gas heats

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such as might be found in a partially premixed compression ignition engine [7]. The advantage of this data set is that it contains almost the entire evaporation history of droplets <100 µm in diameter. The disadvantages are the long preheating period and low Reynolds number, which are not typical of many combustion systems.

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Nomenclature

1D, 2D	one or two-dimensional	T _{dmax}	maximum attained (steady state) droplet temperature
B_H, B_Y	heat and mass transfer (Spalding) numbers		(K)
C_{pl}, C_{pf}, C_{pf}	pg specific heat of the fuel liquid, fuel vapour and reactor	T_{ave}	average droplet temperature (K)
	gases (J/kg K)	T_o, \overline{T}	droplet initial temperature (K) and non-dimensional
C_F	droplet skin friction coefficient		temperature
D, D_o	droplet diameter, initial droplet diameter (µm)	$T_{\rm ref}$	reference temperature for estimating gas properties (K)
F_D, F_B	drag force, buoyancy (N)	T_H	high temperature condition
F_H, F_Y	heat and mass transfer number correction factors	t	time (s or ms)
g	gravitational acceleration (m ² /s)	u, u _{rel}	droplet velocity, relative velocity between droplet and
Ĥ	fuel latent heat of vaporisation (J/kg)		gas (m/s)
k _{fg}	binary diffusion coefficient (m ² /s)	u _s	maximum velocity at the droplet surface (m/s)
K_{v}, K_{vs}	evaporation rate, steady evaporation rate (mm^2/s)	x	axial coordinate (m)
Le	lewis number	Y_s	vapour mass fraction at the droplet surface
m, ṁ _d	droplet mass (kg), droplet mass flow rate (kg/s)	α_l	liquid thermal diffusivity (mm ² /s)
Nu	Nusselt number	$\rho_{\rm l}, \rho_{\rm g}$	droplet and gas densities (kg/m ³)
<i>p</i> , <i>p</i> _v	pressure, fuel vapour pressure (Pa)	δ_M, δ_T	thicknesses of the gas films for mass and heat transfers
Pel	Peclet number in the liquid phase	λ_{g}, λ_{l}	thermal conductivities of gas and liquid fuel (W/m K)
Pr, Pr _l	Prandtl numbers in the gas and liquid phases	λ_{eff}	effective liquid conductivity (W/m K
q	total heat flux to the droplet (W/m^2)	Vg	kinematic viscosity of gas (m^2/s)
$Q_{i} Q_{i}$	total and sensible heat transfer rates to the droplet (W)	μ_{g}, μ_{l}	dynamic viscosities of gas and liquid $(N s/m^2)$
r_s, r_o, \bar{r}_s	droplet radius (µm), initial radius (µm) and non-	ζ	correction factor for liquid reference temperature
	dimensional radius	η	non-dimensional radial coordinate within the droplet
Re, Re _l	droplet and liquid phase Reynolds numbers	Ť	non-dimensional time
Sh	Sherwood number	χ	liquid phase effective heat transport factor
T_d, T_s, T_g	droplet temperature, droplet surface temperature, gas		
0	temperature (K)		

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81 up to the much higher reactor temperature. Accounting for the heating up effect is thus important in modelling the present dro-82 83 plet behaviour. In addition to this, and common to many studies, 84 droplet Reynolds number is less than unity; the droplet velocity 85 is low and close to that of the reactor flow. Dodecane, decane and nonane droplets were investigated with their sizes monitored 86 87 over almost the entire trajectory as they pass through the reactor, thus creating entire evaporation histories from ensembles of indi-88 89 vidual droplets. The need to track the droplets until the end of their 90 lifetimes poses a difficulty in practice, as does the need for repeat-91 able behaviour under complex thermal loading conditions. 92 Through the coupling of magnified digital inline holography and 93 digital particle image correlation, simultaneous measurement of 94 droplet size and velocity yields the evaporation rate time history of individual droplets in the reactor. A detailed description of the 95 experimental technique can be found in [6]. 96

There are a number of options to model droplets the choice of 97 98 which depends on the strength of the convective heat transfer. 99 Options range from the simple lump parameter approach as represented by the infinite conductivity method, to more complex mod-100 els which include internal circulation thus enabling 2D 101 descriptions of temperature and fluid properties. In this paper, a 102 range of approaches with differing complexity will be used to 103 model the data published in [7]. The temperature profile chosen 104 105 for this paper is the high gas temperature condition T_H as described in [7] which has a peak of 680 K. While this temperature resembles 106 107 the peak temperature in a naturally aspirated compression ignition 108 engine, the use of atmospheric pressure with isolated droplets and the very low gas and droplet speeds used, render the conditions 109 different to those found in an engine [8]. 110

Despite its relative simplicity, the infinite conductivity concept [9–11] which treats the liquid phase as a lump model of unique parameters unvarying in space but varying with time, has been shown to be efficient in predicting the near steady droplet evaporation rates in the weak convective conditions studied here,

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although with some small difference in evaporation times [7]. One of the aims of this work is to enhance the treatment of the liquid phase to explore the internal state of the present droplets in relation to their behaviour to explain the efficiency of the lump model.

In general, for convective conditions 2D modelling approaches enable a more detailed description of the droplet internal transport which is generated by the slip effect between the droplet and the surrounding gas. The internal process can be visualised by showing the isothermal streamlines which, for strong convective case, can take the form typical of Hill vortices [12,13]. However, for weak convection it is more efficient to approximate the process by the effective conductivity approach introduced in [9], which treats the liquid phase as a diffusion-limited process with an enhanced coefficient representative of the improved heat transport. Furthermore, while essential for multicomponent fuels, such a detailed 2D description is not essential when modelling for a single component liquid. Another consideration is the possibility of internal transport being caused by the thermo-capillary (Marangoni) effect. This effect is due to an uneven distribution of temperature dependant droplet surface tension. This effect only becomes significant for conditions of high relative speed and/or where variation of the surrounding environment is large (see e.g. [14-16]), which is not the case in this work. For this work, the implementation is simplified to a 1D approach as represented by the choice of the effective conductivity concept.

The simplest form of effective conductivity model used in this 142 study makes use of average liquid properties estimated from a ref-143 erence temperature that is kept constant during the whole process. 144 The aim is then to improve this model in planned steps to assess 145 the importance of liquid phase properties. These models are com-146 pared to the more complete model which numerically solves the 147 liquid heat transfer equation with variable properties. Assessments 148 are finally made for some extended, more realistic conditions 149 which are difficult to study experimentally. 150

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