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Special Issue on CAD/Graphics 2017 As-rigid-as-possible solid simulation with oriented particles

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ABSTRACT

We propose a new as-rigid-as-possible approach to the real-time simulation of physics-based deformable models for interactive applications such as computer games. The key observation is that the efficacy of an embedded oriented particle representation and the stability of a variational implicit formulation of the projective dynamics are complementary to each other. We reformulate the variational implicit formulation to deal with an embedded graph of oriented particles. Our new formulation is extremely stable, and our alternating local/global optimization solver is both easy to implement and computationally efficient. Our method can deal with one-dimensional (cable and rod), two-dimensional (shell), and three-dimensional (solid) models in a uniform manner. Experimental results demonstrate that hundreds of deformable models with an extremely large number of polygons can be simulated robustly in real time using thousands of particles.

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1 1. Introduction

Physics-based simulation is becoming one of the most impor-2 tant techniques for interactive applications such as games. In this 3 paper, we consider physics-based deformation, which has been 4 studied extensively in computer graphics. However, it is still chal-5 6 lenging to robustly produce visually convincing and stable de-7 formation in real time. In real-time applications, the robustness and computational efficiency of deformation simulations are often 8 9 more important than their accuracy.

Recently, several simulation techniques have been developed to 10 compute positions directly instead of integrating velocities or ac-11 celerations to achieve better stability and efficiency. Position-based 12 dynamics (PBD)[1] deals directly with mesh vertices, while mesh-13 less deformation [2] utilizes shape matching to compute the op-14 15 timal rotation and translation for mesh vertices. These positionbased methods are stable regardless of the time step size, and thus 16 17 the deformation can be simulated very efficiently with large time 18 steps.

Two remarkable approaches aim to further improve the PBD method. The first approach employs a simplified structure with a small number of oriented particles to simulate the complex geometry of meshes in a more efficient and robust manner[3]. The second approach called projective dynamics [4] reformulates the implicit time integration of deformation dynamics as energy mini-

http://dx.doi.org/10.1016/j.cag.2017.07.027 0097-8493/© 2017 Elsevier Ltd. All rights reserved. mization in a variational form. In PBD, the material stiffness of a deformable solid is tightly coupled with the convergence of the solver. The formulation of variational implicit integration is similar to PBD, but the projective dynamics method is advantageous because the material stiffness can be specified independently of the solution methods. 30

We found that these two approaches are complementary to 31 each other and they can be combined to take advantage of their 32 benefits. The key challenge is to reformulate the deformation en-33 ergy and momentum potential energy to deal with an embedded 34 graph of oriented particles. Our new formulation yields a com-35 pact formula via clever manipulation of the integral energies and it 36 is extremely stable. Furthermore, our alternating local/global opti-37 mization solver is easy to implement and very efficient for simulat-38 ing complex deformable models in real time. The deformable mod-39 els can be manipulated interactively and the collision between de-40 formable models can be handled efficiently. Our method can deal 41 with one-dimensional (cable and rod), two-dimensional (shell), 42 and three-dimensional (solid) models in a uniform manner. 43

2. Related work

Physics-based simulations of deformable bodies have been researched for decades in computer graphics since the pioneering work by Terzopoulos et al. [5] and many methods have been developed to produce accurate simulations of various types of deformable objects [6]. However, in real-time applications, the robustness and computational efficiency of deformation simulations 50

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51 are often more important than their accuracy, and these require-52 ments are still demanding.

Recently, a number of methods have been proposed for robust 53 54 real-time simulations by evolving the positions of the particles using the initial and predicted positions, before updating the veloci-55 ties based on the positions. Meshless deformation based on shape 56 matching computes the optimal rotation and translation from the 57 initial positions to the predicted positions to deform the mesh ver-58 59 tices [2]. This method is robust regardless of the time step size, and thus it is suitable for real-time applications. However, shape 60 61 matching with a single transformation is restricted to modest de-62 formation. The extension of this method to a lattice applies shape 63 matching to each set of overlapping lattice points with a fast sum-64 mation method to generate large deformation [7]. The fast summation approach has been applied to the simulation of hair with 65 chain shape matching [8], as well as being generalized to an irreg-66 ular structure on a surface mesh with volume preservation [9], and 67 extended to multi-resolution approaches to enhance the conver-68 gence [10,11]. Shape matching with oriented particles is an efficient 69 and robust method for simulating the complex dynamic deforma-70 tion of one-, two-, and three-dimensional deformable bodies even 71 72 with a small number of particles in a single framework [3]. A de-73 formable body is approximated by ellipsoidal particles and shape 74 matching is applied iteratively to each shape matching group comprising a particle and its one-ring neighbors. Our method employs 75 this oriented particle representation so hundreds of deformable 76 models with an extremely large number of polygons can be ap-77 78 proximated using thousands of particles and simulated robustly in real time, as illustrated in Fig. 7. 79

Shape matching can be interpreted as a geometric constraint 80 in PBD [1]. PBD employs an iterative mass-weighted projection 81 82 of constraints in a Gauss-Seidel solver. However, the stiffness is 83 highly dependent on the time step size and the iteration count, 84 which is also problematic in shape matching approaches. Nevertheless, PBD is popular for the real-time simulation of deformable 85 models because of its simplicity and robustness. It has been ex-86 tended to the simulation of fluids [12], rigid bodies [13], elastic 87 88 rods [14,15], and volumetric materials [16], as well as all of them in a unified manner [17]. A comprehensive survey of PBD methods 89 was provided by Bender et al. [18]. Recently, PBD was extended to 90 address the stiffness independently of the time step size and the 91 92 iteration count based on a compliant constraint formulation [19]. However, the extended PBD still employs Gauss-Seidel iterations, 93 94 and thus its convergence is slower than that of projective dynam-95 ics [4,20] where the global solver benefits from the pre-factored 96 system matrix. Moreover, its application to oriented particles has 97 not been addressed previously.

Recently, a number of fast and parallel techniques have been 98 proposed by formulating implicit Euler integration as energy min-99 imization [21,22] and by employing an alternating local/global op-100 timization solver. Liu et al. [20] proposed a local/global solver for 101 102 mass-spring systems, where the local solver deals with nonlin-103 ear terms for direction and the global solver handles stretching. This idea was generalized to other constraints in projective dynam-104 105 ics [4]. A Chebyshev semi-iterative approach was also proposed, which combines the results obtained from previous iterations to 106 107 achieve better convergence in projective dynamics and PBD [23]. Narain et al. [24] employed the alternating direction method of 108 multipliers (ADMM) for implicit time integration and showed that 109 projective dynamics is a special case of ADMM. Their method 110 also allows nonlinear constitutive models and hard constraints. In 111 addition, Liu et al. [25] interpreted projective dynamics as quasi-112 Newton optimization and applied the L-BFGS method to accelerate 113 convergence. In these techniques, the stiffness is largely indepen-114 dent of the iteration count and the solution becomes more ac-115 116 curate as the number of iterations increases. Our method is also

based on implicit Euler integration formulated as energy minimization, and thus it differs from the original oriented particles approach where the stiffness is highly dependent on the time step size and the iteration count.

In as-rigid-as-possible (ARAP) surface modeling [26], a block co-121 ordinate descent method is employed to iteratively minimize the 122 shape deformation energy in alternating local/global optimization 123 steps. This method introduces a local rotation at each vertex to de-124 fine an ARAP deformation energy by using the squared distances 125 between the locally rotated positions of its neighbors and the ac-126 tual deformed positions. The auxiliary rotations and the positions 127 of the vertices should be optimized. An optimal rotation at a vertex 128 is computed in parallel with shape matching of its neighboring ver-129 tices, which requires the polar decomposition of a shape matching 130 matrix. The optimal positions are computed efficiently by solving 131 a linear system, which depends only on the initial mesh, so it can 132 be pre-factored with a sparse Cholesky decomposition. 133

Embedded deformation for shape manipulation [27] introduces 134 a deformation graph of nodes corresponding to rigid transforma-135 tions that deform nearby space, and then defines an ARAP defor-136 mation energy over the deformation graph. In contrast to ARAP 137 surface modeling, this approach can deal with a wide range of 138 shape representations, such as meshes, polygon soups, mesh an-139 imations, and animated particle systems. However, the node has 140 no volume, and thus it cannot deal with the deformation of 141 a one-dimensional structure robustly, unlike the oriented parti-142 cles approach and our proposed method. In addition, the itera-143 tive Newton-Gauss method employed for nonlinear optimization 144 requires more time than the iterative local/global optimization 145 method when only seeking a plausible solution in a small number 146 of iterations. 147

3. Method

In this paper, we introduce a deformation graph \mathcal{G} comprising 149 oriented particles, which approximate a flexible body and deform 150 nearby spaces robustly even with a small number of particles. The 151 positions and orientations of the oriented particles are computed 152 stably by solving an energy minimization problem formulated as 153 implicit Euler integration in a variational form. The total energy 154 $E(\mathcal{G})$ comprises the momentum potential energy $E^k(\mathcal{G})$, ARAP de-155 formation energy $E^{e}(\mathcal{G})$, and direct manipulation constraint energy 156 $E^{c}(\mathcal{G})$: 157

$$E(\mathcal{G}) = E^{k}(\mathcal{G}) + E^{e}(\mathcal{G}) + E^{c}(\mathcal{G}).$$
⁽¹⁾

We develop an iterative local/global optimization solver to seek a plausible solution in an efficient and robust manner. The final deformed vertices are obtained by linear blend skinning of the rigid transformations stored in the oriented particles, which can be implemented with GPU skinning.

Our main contribution lies in developing the ARAP deforma-163 tion energy over a deformation graph comprising oriented parti-164 cles. Thus, we first explain the deformation graph and the cor-165 responding space deformation in Section 3.1, before deriving the 166 ARAP deformation energy in Section 3.2. We explain the alternat-167 ing local/global optimization solver before considering other en-168 ergies to make the explanation clearer. The momentum potential 169 energy and direct manipulation constraint energy are described in 170 Sections 3.3 and 3.4, respectively. 171

3.1. Deformation graph

Suppose that the *j*th ellipsoidal particle \mathcal{E}_j is transformed from 173 the rest position $\bar{\mathbf{x}}_j$ and orientation $\bar{\mathbf{E}}_j$ to the current position \mathbf{x}_j 174 and orientation $\mathbf{E}_j = [\mathbf{e}_j^a | \mathbf{e}_j^b | \mathbf{e}_j^c]$, where \mathbf{e}_j^a , \mathbf{e}_j^b , \mathbf{e}_j^c are the current 175 unit axes of \mathcal{E}_j . We denote the set of particles directly connected to 176

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