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Long-term H-release of hard and intermediate between hard and soft amorphous carbon evidenced by in situ Raman microscopy under isothermal heating



DIAMOND RELATED MATERIALS

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ABSTRACT

We study the kinetics of the H release from plasma-deposited hydrogenated amorphous carbon films under isothermal heating at 450, 500 and 600 °C for long times up to several days using in situ Raman microscopy. Four Raman parameters are analyzed. They allow the identification of different processes such as the carbon network reorganization and the H release from sp³ or sp² carbon atoms and the corresponding timescales. Carbon reorganization with aromatization and loss of sp³ hybridization occurs first in 100 min at 500 °C. The final organization is similar at all investigated temperatures. Full H release from sp³ carbon occurs on a longer timescale of about 10 h while H release from sp² carbon atoms is only partial, even after several days. All these processes occur more rapidly with higher initial H content, in agreement with what is known about the stability of these types of films. A quantitative analysis of these kinetics studies gives valuable information about the microscopic processes at the origin of the H release through the determination of activation energies.

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1. Introduction

Raman spectroscopy is a powerful tool to rapidly characterize C-based materials that contain sp^2 carbon atoms (C(sp^2)). Interpreting the 1000–1800 cm⁻¹ spectral region gives information on the carbon organization (aromatization). Because the Raman cross section for $C(sp^3)$ is one to two orders of magnitude lower than that of $C(sp^2)$, two bands due to C(sp²) dominate this spectral range even when samples contain only a few percent of C(sp²). They are the well-known G and D bands [1]. They contain information on disorder such as the size of aromatic domains, the $C(sp^2)/C(sp^3)$ ratio for amorphous carbons (a-C), and the H content for hydrogenated carbons (a-C:H) [1–3]. Raman parameters generally used to probe the bonding structure are: the width and the position of the G and D bands, denoted $\Gamma_{G,D}$ and $\sigma_{G,D}$ respectively, the height ratio of these two bands, H_D/H_G , and if applicable, the m/H_G parameter with m being the slope of the photoluminescence background calculated between 800 and 2000 cm^{-1} . In the case of heat-treated a-C:H films for which the H content varied from 2 to '30 at. %, we have recently obtained a linear relation between $log_{10}(m/H_G)$ and the H content $(H/H + C = 25 + 9 \log_{10}(m/H_G))$ [4], similar to what was previously obtained for as-deposited samples for which the H content varied from 20 to 47 at. % [5,6]. We have emphasized that m/H_G is sensitive not only to the H content, but also to defect passivation [7,8], possibly by H migration, and we also found that m/H_G is due only to H bonded to $C(sp^3)$. In addition we have obtained a linear relation between H_D/H_G and the H content (H/H + C = 0.54 - 0.53 H_D/H_G), independent of whether the H is bonded to $C(sp^2)$ or to $C(sp^3)$. H_D/H_G can thus be used to probe the total H content in the range from 2 to 30 at.%. Conversely, we have shown that the Γ_G and the σ_G parameters are sensitive mainly to the structural changes such as the re-organization leading to a larger aromatization and the loss of $C(sp^3)$ induced by heating.

To obtain more information on the different processes involved under heating we investigate here the long-term kinetics of the isothermal hydrogen release from a-C:H films with initial H/H + C ratios of 0.29, 0.32 and 0.37 at temperatures of 450 °C, 500 and 600 °C by means of in situ Raman microscopy. We show how the evolution kinetics provide valuable information on the H release processes and emphasize the role of the initial structure and H-content on these processes.

2. Experimental

Plasma-deposited hydrogenated amorphous carbon films, a-C:H, with H/H + C \sim 0.29, 0.32 and 0.37 (DC self bias of - 300, -200 and



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-100 V, respectively), and a thickness of e ~ 0.20 μm (see ref. [9] for details on the deposition method) were studied together with a thicker film with H/H + C = 0.32 (applied substrate bias: -200 V) and a thickness of e ~ 1 μm (see ref. [5] for details on the deposition method). The 0.29 and 0.32 films are hard films while the 0.37 film is intermediate between a hard and a soft film [9,10]. Note that the thick film has a larger C(sp³) content than the three thin films: this is indicated by its G band position at 1540 cm⁻¹ instead of 1522 cm⁻¹ for the thin films, their G band widths being similar (see Fig. 9 in [2]). The samples with the as-deposited films were cut into several pieces, each piece being heated under argon flow in a cell at 1.5 bar with a linear ramp of 150 °C min⁻¹ to reach the working temperature. Time zero was taken at the end of the ramp. Temperatures 450, 500 and 600 °C were studied.

Raman spectra were obtained using a Horiba-Jobin-Yvon HR LabRAM apparatus (laser wavelength: $\lambda_L = 514.5$ nm, 50X objective, numerical aperture of 0.5 leading to a laser focus diameter of 2.5 µm, resolution ~ 1 cm⁻¹). The laser power was chosen to have a good signal/noise ratio and/or to prevent damages and three values were tested. 0.01, 0.1 and 1 mW, corresponding to power densities from 2×10^2 to 2×10^4 J cm $^{-2}$ s $^{-1}$. Unless otherwise stated, spectra were recorded in situ at the working temperature. The cell was a commercial LINKAM TS1500, argon was Alphagaz 2 (99.9995% purity, Air Liquid Company). The Raman parameters analyzed were the G band wavenumber, σ_{G} , the G band full-width at half-maximum, Γ_G , the relative height of the G and D bands, H_D/H_C , and the m/H_C parameter. Heights were measured on the raw data after the linear baseline calculated between 800 and 2000 \mbox{cm}^{-1} (slope m) was subtracted. In that case, $\mbox{H}_{\mbox{D}}$ was therefore measured at its apparent maximum, except when the D band maximum was not sufficiently well defined (i.e. for the as-deposited sample). In the latter case H_D was taken at 1370 cm⁻¹. All these parameters are presented in Fig. 1.

3. Results and discussion

Fig. 1 displays Raman spectra of the a-C:H film ($e = 0.2 \mu m$, H/H + C = 0.32), as deposited at room temperature and after 4, 45 and 500 min at 500 °C. The Raman spectrum of the as-deposited film is composed of the underlying silicon wafer signature at 520 cm⁻¹, a G band at 1522 cm⁻¹, and a broad shoulder at ~1200 – 1400 cm⁻¹ containing the D band. When the sample is heated, the silicon signature diminishes because the a-C:H becomes more absorbent [11,12]. The photoluminescence background significantly increases, due to the defect passivation (and the quenching of non-radiative relaxation



Fig. 1. Raman spectra of the a-C:H film ($e = 0.2 \ \mu m$ and H/H + C = 0.32) at room temperature, as deposited (AD) and at 500 °C after 4, 45 and 500 min at 500 °C.

processes [4,6]) before it starts decreasing as hydrogen is released. The G band narrows and slightly blueshifts, as previously observed in [4,13]. The D band also narrows and its height increases. According to previous studies [2,13,14], these evolutions reveal the organization of the material, i.e. a loss of $C(sp^3)$ and increasing aromatization. When the heating time increases, the Raman spectra keep changing: the G and D bands become more clearly distinguishable, H_D/H_G increases and m/H_G decreases. A long timescale process thus occurs, which we investigate in more detail below.

Fig. 2 displays the isothermal evolution with time up to 800 min of $\sigma_{\rm G}$, $\Gamma_{\rm G}$, m/H_G and H_D/H_G at 500 °C at the three laser powers, 0.01, 0.1 and 1 mW for the a-C:H film ($e = 0.2 \mu m$, H/H + C = 0.32). No significant effect of the laser power is observed, indicating that no additional heating due to the laser irradiation is induced even at 1 mW, contrary to what was measured in [15] for thinner (15 nm) films. Fig. 2a and b displays a similar isothermal evolution for Γ_G and σ_G , reaching plateaus in ~ 100 min. $\Gamma_{\rm G}$ is at 180 cm⁻¹ for the as-deposited film and the plateau at 500 °C is at 85 cm $^{-1}$. σ_{G} is at 1522 cm $^{-1}$ for the as-deposited film and the plateau at 500 °C is at 1580 cm⁻¹. These evolutions suggest that at 500 °C, the film gets rapidly more organized (in typically 100 min) and that this organization then does not significantly evolve. Conversely, Fig. 2c and d show that for m/H_c and H_D/H_c this rapid evolution is followed by a slow and roughly linear change (decrease for m/H_{C} and increase for H_{D}/H_{C}). H_{D}/H_{C} and m/H_{C} are both correlated with the H content [4]. We thus interpret these evolutions as due to the isothermal H release, with a first phase occurring rapidly together with structural changes (within about 100 min), followed by a second phase occurring on a longer timescale.

To test the spatial homogeneity of heating-induced changes, we map the film at room temperature after heating at 500 °C for 800 min, in $a \sim 35 \times 35 \ \mu m^2$ area, much larger than the 5 μm^2 area of the laser spot, with the 0.01 mW laser power. Histograms obtained by analyzing the maps of the four parameters, σ_G , Γ_G , m/H_G and H_D/H_G are displayed in Fig. 3. Their widths are small ($\sigma_G = 1601 \pm 2 \text{ cm}^{-1}$, $\Gamma_G = 76 \pm$ 4 cm $^{-1}$, m/H_G = 0.25 \pm 0.09 μm and H_D/H_G = 0.77 \pm 0.02), indicating that the sample is uniform, both in structure and in composition. The values obtained for Γ_G , m/H_G and H_D/H_G at the end of the 800 min heating at 500 °C fall in this range, confirming that the spot irradiated by the laser is not more heated than the other parts of the sample. Conversely, $\sigma_{\rm C}$ is at 1580 cm⁻¹ at the end of the heating process, significantly outside the histogram values, and this is due to a reversible temperature effect due to multi-phonon interactions, similarly to what was found for graphite or graphene [16]. We have checked that the effect was reversible by varying alternately the temperature between room temperature and 500 °C.

Fig. 4 is a plot similar to that of Fig. 2 to compare the isothermal evolutions of the three thin films ($e = 0.2 \mu m$ and H/H + C = 0.29, 0.32and 0.37) at 500 °C and that of the thick film ($e = 1 \ \mu m$ and H/H + C = 0.32) at 450 °C. The laser power was 1 mW, except for the first 400 min at H/H + C = 0.29 for which it was only 0.1 mW (as the spot was located at another place on the sample, the small jump in the data is due to the small width of the structure and H-content distribution, see Fig. 3). For all the Raman parameters at 500 °C and for the first 100 min the faster the evolution the higher the initial H content. This is in agreement with what is known about the stability and the H desorption of this type of hydrogenated films [17-21], the higher the H content, the less stable is the film and the more rapid is the H desorption. For longer times, the plateaus reached by σ_G and Γ_G are almost independent of the initial H content, indicating that the structure of the film at 500 °C is independent of the initial H content. m/H_G goes to zero for the three films at 500 °C while H_D/H_G keeps slowly increasing along three distinct lines. As H_D/H_G depends on to the total H content and $log_{10}(m/H_G)$ only on H bonded to $C(sp^3)$, this indicates that at this stage of the kinetics all remaining H is bonded to C(sp²). Furthermore, the remaining H content is higher for lower initial H content. The thick film evolution of $\sigma_{\rm C}$ and $\Gamma_{\rm C}$ and m/H_{C} at 450 °C is intermediate between the thin films having the H

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