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International Communications in Heat and Mass Transfer xxx (2016) xxx-xxx



Contents lists available at ScienceDirect

International Communications in Heat and Mass Transfer

journal homepage: www.elsevier.com/locate/ichmt

Numerical investigation of two-phase laminar pulsating nanofluid flow in helical microchannel filled with a porous medium

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ARTICLE INFO

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ABSTRACT

Numerical results of the thermal and hydraulic characteristics of pulsating laminar flow in a three-dimensional 15 helical microchannel heat sink (HMCHS) filled with a porous medium using Al₂O₃-water based nanofluid are 16 presented in this paper. The two-phase mixture model with modified effective thermal conductivity and viscosity 17 equations were used to solve the laminar flow governing equations. The detailed results for thermal and flow 18 fields are reported for the effects of amplitude (1-3) and frequency (5-20 rad/s). The presence of porous 19 media led to better thermal enhancement and significant improvement in heat transfer performance was 20 observed for sinusoidal velocity inlet conditions as compared to steady flow conditions. In addition, the proposed 21 model showed better accuracy in predicting nanofluid flow. 22

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33 1. Introduction

Research interest in the field of electronic thermal management 34has intensified in the past decade in order to address the ever increasing 35heat loads. Microchannel heat sink (MCHS) is a preferred choice of 36thermal management and also widely used since the first introduction 37 by Tuckerman and Pease [1,2] in 1981 and followed by numerous 38 researches [3–10] which have been carried out to further improve its 39 performance. The secondary flow in curved tubes is of particular inter-40 est to researches as it leads to significant heat transfer enhancement. 03 The secondary flow is induced by the centrifugal forces which act on 04 the flowing fluid. This secondary flow is perpendicular to the main 43 44 axial flow. The fluid is directed from the inner wall of the tube to the outer wall of the tube across the center and it is forced back to the 45inner wall again. The Dean number ($De = Re\sqrt{\frac{T}{HR}}$) which depicts the 46 magnitude of the secondary flow for curved tubes was first introduced 47 by Dean, a British applied mathematician and fluid dynamist [11–12]. 48

49Alam and Kim [13] numerically studied the mixing in a curved microchannel with grooves on side-walls. They observed that grooved 5051microchannels lead to better mixing performance compared to smooth microchannels and pressure drop was lesser for grooved microchannels 52than smooth channels. Chu et al. [14] examined the flow characteristics 53in a curved rectangular microchannel with a different aspect ratio and 54curvature ratio. The pressure drop increased on decreasing channel 55width and the effect intensified at larger flow rates. Curved geometries 56had higher pressure drop due to the secondary flow. In another separate 57study, Chu et al. [15] investigated the characteristics of water through 58

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http://dx.doi.org/10.1016/j.icheatmasstransfer.2016.03.016 0735-1933/© 2016 Published by Elsevier Ltd. curved rectangular microchannels for various aspect ratios and curva- 59 ture. The experimental results showed that classical Navier–Stokes 60 equations were still applicable for the incompressible laminar flow for 61 curved microchannels. 62

Analysis of heat and fluid flow using a porous medium has been 63 studied extensively in the past [16-23]. Singh et al. [24] analyzed the 64 sintered porous heat sink used for cooling high-powered compact 65 microprocessors. They found that the system was able to dissipate 66 2.9 MW/m^2 of heat flux with a pressure drop of 34 kPa. Hooman [25] 67 analytically studied cases of fully developed forced convection in a rect- 68 angular channel with and without a (Brinkman) porous medium. The 69 Nusselt number increases on increasing porous medium shape parame-70 ter. Wan et al. [26] developed a novel micro heat sink thermal manage- 71 ment system for high power LEDs. The numerical results showed that 72 the heated surface temperature of the heat sink was low at high heat 73 flux. The average heat transfer coefficient increased on increasing 74 velocity. Hung et al. [27] studied the heat transfer analysis on three 75 dimensional porous microchannel heat sinks of various configurations. 76 They concluded that the presence of a porous medium lead to an 77 additional pressure drop with the sandwich distribution design having 78 better overall thermal performance with the least amount of pressure 79 drop.

Zhao and Lu [28] carried out a two-dimensional analytical and nu- 81 merical study to analyze the forced convective heat transfer characteris- 82 tic in a rectangular MCHS. They observed an increase of the overall 83 Nusselt number due to a reduction in the effective thermal conductivity 84 ratio. Moraveji et al. [29] performed a numerical study on the cooling 85 performance and pressure drop in a mini-channel heat sink using 86 nanofluids. They found that the heat transfer coefficient increased on in- 87 creasing nanoparticle concentration and Reynolds number. Fani et al. 88

Please cite this article as: S. Sivasankaran, K. Narrein, Numerical investigation of two-phase laminar pulsating nanofluid flow in helical microchannel filled with a porous medium, Int. Commun. Heat Mass Transf. (2016), http://dx.doi.org/10.1016/j.icheatmasstransfer.2016.03.016

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S. Sivasankaran, K. Narrein / International Communications in Heat and Mass Transfer xxx (2016) xxx-xxx

T1.1	Nomenclature		
T1.2	Α	dimensionless amplitude	
T1.3	C_p	heat capacity (J/kg K)	
T1.4	D	channel depth (m)	
T1.5	D_h	hydraulic diameter $(2W_{ch}D / W_{ch} + D) (m)$	
T1.6	De	Dean number	
T1.7	HR	helix radius	
T1.8	h	heat transfer coefficient (W/m ² K)	
T1.9	Kp	permeability (m ²)	
T1.10	λ	thermal conductivity (W/m K)	
T1.11	L	length of channel (m)	
T1.12 Q2	Nu	Nusselt number ($Nu = hD_h/k$)	
T1.13	Р	pressure (Pa)	
T1.14	q	heat transfer rate (W)	
T1.15	r	inner tube radius (m)	
T1.16	R_c	curvature ratio	
T1.17	Re	Reynolds number ($Re = \rho VD_h / \mu$)	
T1.18	Т	temperature (K)	
T1.19	t	time (s)	
T1.20	V	velocity (m/s)	
T1.22	W_{ch}	channel width	
T1.23	Greek syn	Greek symbols	
T1.24	α	channel aspect ratio (D / W _{ch})	
T1.25	к	Boltzmann constant	
T1.26	3	porosity (%)	
T1.27	ρ	density (kg/m ³)	
T1.28	μ	viscosity (kg m/s)	
T1.29	Ø	volume fraction of nanoparticle	
T1.30	f	frequency (rad/s)	
T1.32	Subscript		
T1.33	ch	channel	
T1.34	f	fluid	

[30] examined the effect of nanoparticle size on thermal performance of
nanofluid through a trapezoidal MCHS. The pressure drop increased on
increasing concentration. Heat transfer decreased on increasing particle
size. Emran and Islam [31] numerically studied the flow dynamics and
heat transfer characteristics in a MCHS. They observed that the highest
temperature was at the bottom of the heat sink immediately below
the channel outlet and the lowest at the channel inlet.

Lotfi et al. [32] studied the forced convective heat transfer of 96 nanofluid flow through a horizontal circular tube using the two-phase 97 mixture model and two-phase Eulerian model. They concluded that 98 the mixture model provided better results as compared to all the 99 100 other models. Kalteh et al. [33] carried out both a numerical and an ex-101 perimental study on the laminar convective heat transfer of Al₂O₃–H₂O nanofluid flow through a wide rectangular MCHS. The two-phase 102Eulerian-Eulerian model was used. They observed that the two-phase 103numerical results were in good agreement with experimental results 104 as compared to single-phase results. Targui and Kahalerras [34] ob-105served heat transfer enhancement due to the addition of an oscillating 106 component to the mean flow structure. The peak performance was ob-107 tained for the case of pulsating hot fluid. Nandi and Chattopadhyay [35] 108 numerically studied the heat transfer characteristics of water through a 109two-dimensional wavy microchannel subject to pulsating inlet fluid 110 condition. They deduced that the pulsating inlet fluid condition led to 111 improved heat transfer performance with pressure drop within 112 acceptable limits. The research was further continued by Akdag et al. 113 114 [36] using nanofluid. They observed an increase in heat transfer performance on increasing nanoparticle volume fraction and amplitude 115 of pulsation. 116

To the authors' best knowledge, no analyses were made in the past 117 to study the combined effect of secondary flow, porous medium, 118 nanofluid and pulsating flow on the heat transfer and fluid flow charac- 119 teristics; hence, this has motivated the present study. The present study 120 examines the three-dimensional laminar flow in HMCHS filled with a 121 porous medium and saturated with Al_2O_3 -water based nanofluid. The 122 results for the effects of amplitude (1–3) and frequency (5–20 rad/s) 123 are reported. 124

2. Mathematical model	12

2.1. Governing equations

The helical microchannel heat sink considered in this study is 127 shown schematically in Fig. 1. Heat is transferred from the outer wall 128 of the HMCHS to the nanofluid which flows through the channel. Sever- 129 al assumptions are made on the operating conditions of the HMCHS as 130 follows: (i) the MCHS operates under steady-state conditions; (ii) the 131 properties of the HMCHS material are temperature independent; 132 (iii) the external heat transfer effects are ignored; (iv) the outer walls 133 of the MCHS are adiabatic. 134

The two-phase mixture model governing equations for mass, 135 momentum and energy are [37]: 136 Continuity equation: 137

nanaty equation.

$$\frac{\partial}{\partial t}(\rho_m) + \nabla \cdot \left(\rho_m \vec{v}_m\right) = \mathbf{0} \tag{1}$$

$$\rho_m = \sum_{k=1}^n \emptyset_k \rho_k. \tag{2}$$

Momentum equation:

$$\frac{\partial}{\partial t} \left(\rho_m \overrightarrow{\mathbf{v}}_m \right) + \nabla \cdot \left(\rho_m \overrightarrow{\mathbf{v}}_m \overrightarrow{\mathbf{v}}_m \right) = \nabla \rho_m + \nabla \cdot \left[\mu_m \left(\nabla \overrightarrow{\mathbf{v}}_m + \overrightarrow{\mathbf{v}}_m^T \right) \right] + \rho_m \overrightarrow{g} + \nabla \cdot \left(\sum_{k=1}^n \partial_k \rho_k \overrightarrow{\mathbf{v}}_{dr,k} \overrightarrow{\mathbf{v}}_{dr,k} \right) \right]$$
(3)

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where *n* is the number of phases, \emptyset is the volume fraction and μ_m is the viscosity of the mixture: 144

$$\mu_m = \sum_{k=1}^n \mathcal{O}_k \mu_k \tag{4}$$

and $\overrightarrow{v}_{dr,k}$ is the drift velocity for secondary phase k defined as:

$$\overrightarrow{v}_{dr,k} = \overrightarrow{v}_k - \overrightarrow{v}_m. \tag{5}$$

Energy equation:

$$(\Theta_k \rho_k h_k) + \nabla \cdot \left(\sum_{k=1}^n \Theta_k \overrightarrow{v}_k (\rho_k h_k + p) \right) = \nabla \cdot \left(\lambda_{\text{eff}} \nabla T \right). \tag{6}$$

The representative viscosity for the nanoparticle is estimated as:

$$\mu_{np} = \frac{\frac{\mu_m}{\mathcal{O}_{bf}\mu_{bf}}}{\mathcal{O}_{np}} \tag{7}$$

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 $\frac{\partial}{\partial t} \sum_{k=1}^{n}$

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