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¹ Traveling wave solutions to incompressible unsteady 2-D laminar flows with heat ² transfer boundary $\stackrel{\sim}{\sim}$

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ABSTRACT

Analytical solutions play important roles in the understanding of fluid dynamics and heat transfer related prob- 15 lems. Some analytical solutions for incompressible steady/unsteady 2-D problems have been obtained in litera- 16 ture, but only a few of those are found under heat transfer conditions (which brings more complexities into the 17 problem). This paper is focused on the analytical solutions to the basic problem of incompressible unsteady 2-D 18 laminar flows with heat transfer. By using the traveling wave method, fluid dynamic governing equations are de- 19 veloped based on classical Navier-Stokes equations and can be reduced to ordinary differential equations, which 20 provide reliable explanations to the 2-D fluid flows. In this study, a set of analytical solutions to incompressible Q10 unsteady 2-D laminar flows with heat transfer are obtained. The results show that both the velocity field and 22 the temperature field take an exponential function form, or a polynomial function form, when traveling wave 23 kind solution is assumed and compared in such fluid flow systems. In addition to heat transfer problem, the ef- 24 fects of boundary input parameters and their categorization and generalization of field forming or field evolutions 25 are also obtained in this study. The current results are also compared with the results of Cai et al. (R. X. Cai, N. 26 Zhang, International Journal of Heat and Mass Transfer, 2002, 45: 2623-2627) and others using different 27 methods. It is found that the current method can cover the results and will also extend the fluid dynamic 28 model into a much wider parameter ranges (and flow situations). 29

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35 1. Introduction

Analytical solutions have always played important roles in the 36 development of various fluid flow systems, which usually serve as fun-37 damental basis for comparison of fluid dynamic nature and system evo-38 39 lution trends. For example, the analytical solutions of incompressible flow and constant coefficient heat conduction in early days have been 40 the bases of fluid dynamics and heat transfer [1,2]. However, it is diffi-41 cult to derive the analytical solutions of the governing equations with 4243 nonlinear terms, especially the complex governing equations (for fluid systems usually the Navier-Stokes equations are considered) with 44 given initial and boundary conditions. 45

Although with the rapid development of computers and numerical
 methods, much research has focused on formulating efficient numerical
 methods to solve fluid dynamics and coupled heat transfer problems, the
 accuracy of these numerical solutions can only be ascertained by com parison with exact solutions or empirical/half-empirical correlations

http://dx.doi.org/10.1016/j.icheatmasstransfer.2015.05.006 0735-1933/© 2015 Published by Elsevier Ltd. from experiments. Therefore, it is meaningful to find out some analytical 51 solutions not only for the reason that they are found be able to describe 52 the detailed behavior of the concerning system, but also that they can 53 be used as benchmark solutions to check the accuracy, convergence 54 and effectiveness of various numerical methods and solutions, and to im- 55 prove various numerical methods such as their differencing schemes and 56 grid generation skills [3–6]. 57

For fluid dynamic systems, the well-known Navier-Stokes (N-S) 58 equations, first introduced by Navier in 1821, and developed by Stokes 59 in 1845, are the fundamental governing equations. For those two hun- 60 dred years, many groups have tried to solve this problem. The work (Q11 on the exact solutions of the Navier-Stokes equations has also accumu- 62 lated in literature. However, due to the nonlinearity and complexity 63 of Navier-Stokes equations, one can only give the solutions to very 64 limited/simplified cases. Indeed there only exist a small number of 65 exact solutions in literature. In the paper of Wang [7,8], one can find 66 the historical reviews of the trials and solutions up to year 1991. And 67 in the most recent years, super computers have made it possible to nu- 68 merically solve the Navier-Stokes equations and the accuracy of the re- 69 sults can be compared with an exact solution. Thus, the exact solutions 70 are very important as a test to verify numerical or empirical methods for 'Q12 complex flow problems. In the recent twenty or thirty years, major de-72 velopments of the exact solutions of the Navier-Stokes equations can be 73

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B.1	Nomenclature	
B 2	α	thermal diffusivity, [m ² /s]
B.3	v	kinematic viscosity, [m ² /s]
B.4	C_p	specific heat, [J/(kg·K)]
B.5	ρ	density, [kg/m ³]
B.6	g	gravitational acceleration, [m/s ²]
B.7	θ	temperature, [K]
B.8	р	pressure, [Pa]
B.9	u,w	velocity components in <i>x</i> , <i>z</i> directions, [m/s]
B 10	t	time, [s]
B 11	u_0	initial velocity, [m/s]
B 12	p_0	initial pressure, [Pa]
B 13	x	abscissa, [m]
B 14	Ζ	third coordinate, [m]
B 15	ξ	transformation $\xi = ax + bz + ct$, [m]
B.16	a,b	constants, $a^2 + b^2 \neq 0$
B.17	С	wave speed, [m/s]
B.19	$m_{i_i} M_{i_j}$	constants used during the integration of equations
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found in Chandna and Oku-Ukpong [9], Profilo et al. [10], Siddique [11],
Venkatalaxmi et al. [12], Warsi [13], Naeem and Jamil [14], Naeem and
Younus [15] and others [16].

77 The major methods used in solving the Navier-Stokes equations are 78variable separation and equation transformation methods; the equation 79transformation method then includes potential function assumption and parameter transformation. The basic process is to change/separate 013 81 the equations to reduce the problem to the class of nonlinear ordinary differential equations for the first step [17]. Indeed, variable separation 82 has been widely used in solving complex equations. Potential functions 014 are also often used in similar methods [16,18]. This method has to 015 assume the independence relations between the variables, making it 85 difficult to justify the fundamentals of the solving process and the re-86 87 sults obtained. Besides the results mentioned in review paper of Wang [7,8], in recent years representative ones can be found by many groups. 88 For example, Al-Mdallal [19] utilized the canonical transformation with 89 complex coefficients; the Navier-Stokes equations were reduced to a 90 91 linear partial differential equations that can be solved by using separa-92tion of variables. According to the type of the canonical transformation 93 constants (real or complex), Al-Mdallal [19] got different types of 94exact solutions. By using variable separation method with addition, 95Cai and Zhang. [20] obtained some exact solutions to incompressible unsteady 2-D laminar flow with heat transfer, neglecting gravity, 96 97 radiation and internal heat source. By this method, Cai discussed a series of problems described by Navier-Stokes equations and found a set of 98 useful solutions to some basic flow situations [21,22]. 99

Variable or equation transformation methods are also generally seen 100 in literatures. For example, Nugroho et al. [17,23] proposed a potential 101 102function and transformed coordinate to alter the three-dimensional 103 incompressible Navier-Stokes equations into simpler forms. Furthermore, a special class of solutions to the three-dimensional incompress-104ible Navier-Stokes equations was obtained by dropping the pressure 105gradient and it was found that constant pressure gradient will produce 106 107 similar solutions to that of a zero pressure gradient. The authors also proposed another potential function and transformed coordinate, 108 which has a nontrivial relation with respect to time, and a general 109 functional form of static pressure was applied. Fang et al. [24] have 110 investigated the steady momentum and heat transfer of a viscous fluid 111 flow over a stretching/shrinking sheet, and have presented new exact 112 solutions for the Navier-Stokes equations. These solutions provide a 113 more general formulation including the linear stretching and shrinking 114 wall problems as well as the asymptotic suction velocity profiles over a 115 116 variety of situations.

As discussed, one case was the transformation of the Navier–Stokes 117 equations to the Schrödinger equation, performed by application of the 118 Riccati equation [25] and to achieve much simpler forms. This has good 119 prospects since the Schrödinger equation is linear and has well defined 120 solutions. The method of Lie group theory was also applied in order to 121 transform the original partial differential equations into ordinary differential systems [26]. The same route was followed by Meleshko [27] and 123 by Thailert [28], in transforming the Navier–Stokes equations to solvable linear systems. 125

The current study is one continued trial of solving channel flow with 126 heat transfer by using the transformation method. The current study is 127 focused on the transforming of partial differential equations into tracta-128 ble ordinary differential equations or some particular partial differential 129 equations. Traveling wave method belongs to this family; due to the ap-**Q16** plication of the transform $\xi = \sum a_i x_i$, partial differential equations can 131 be reduced to tractable ordinary differential equations, where ξ is a var-132 iable of ordinary differential equations and it is called a phase of the 133 wave, $x_1, ..., x_n$ are independent variables of the partial differential 134 equations, and $a_1, ..., a_n$ are arbitrary constants.

In this paper, transformation method and traveling wave solution 136 are used to solve the incompressible unsteady 2-D laminar flow with 137 boundary heat transfer problem. The current study takes the simplified 138 2-D laminar flow conditions, where the instability shear flow with wave 139 transportation happens. Such phenomena indeed are fundamental and 140 critical during the formation of shear flow and the establishment of vis- 141 cous and thermal boundary layers. The basic mathematical model are 142 carefully established and tested with several general cases. It is found 143 that the current model is capable of covering such benchmark cases 144 and can provide new information on the parameter evolutions of 145 Navier-Stokes governed fluid dynamic systems. The following parts of 146 this paper are arranged as follows: In Section 2, the government equa- 147 tions of the incompressible unsteady 2-D laminar flow with heat trans- 148 fer are presented. In Section 3, two kinds of traveling wave solutions are 149 obtained under various conditions. For the first one, the velocity and 150 temperature fields mainly depend on exponential functions, while for 151 the second one, they depend on polynomial functions. In Section 4, Q17 the benchmark case results are analyzed and explained in detail. 018 Further, the physical descriptions of these solutions are explained and 154 some solutions with certain conditions are shown. A short conclusion 155 is given in Section 5. 156

2. Basic model description and governing equations

2.1. Physical model and basic Navier–Stokes equations

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In this paper, the general governing equations of the incompressible 159 unsteady 2-D laminar flow with heat transfer is considered. The 160 governing equations consisted of continuity equation, incompressible 161 fluid Navier–Stokes equation and thermal conservation equation, 162 which are written as follows (neglecting radiation and internal heat 163 source; *z*-direction is opposite to gravitation) [20]: 164

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0 \tag{1}$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + w \frac{\partial u}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + \upsilon \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial z^2} \right)$$
(2)
169

$$\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + w \frac{\partial w}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial z} + v \left(\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial z^2} \right) - g \tag{3}$$

$$\frac{\partial\theta}{\partial t} + u\frac{\partial\theta}{\partial x} + w\frac{\partial\theta}{\partial z} = \alpha \left(\frac{\partial^2\theta}{\partial x^2} + \frac{\partial^2\theta}{\partial z^2}\right) + \frac{\upsilon}{C_p} \left[\left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x}\right)^2 + 2\left(\frac{\partial u}{\partial x}\right)^2 + 2\left(\frac{\partial w}{\partial z}\right)^2 \right]$$
(4)

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