Contents lists available at ScienceDirect



### International Journal of Thermal Sciences

journal homepage: www.elsevier.com/locate/ijts



## Investigation of geometry and dimensionless parameters effects on the flow field and heat transfer of impingement synthetic jets



Mohammad Hatami<sup>a,\*</sup>, Farzad Bazdidi-Tehrani<sup>b</sup>, Ahmad Abouata<sup>b</sup>, Akbar Mohammadi-Ahmar<sup>c</sup>

<sup>a</sup> Mechanical Engineering Department, Ohio University, Stocker Center, Athens, OH 45701, USA

<sup>b</sup> Mechanical Engineering Department, Iran University of Science and Technology, Tehran, Iran

<sup>c</sup> School of Mechanical Engineering, College of Engineering, University of Tehran, Tehran, Iran

#### ARTICLE INFO

Keywords: Synthetic jet Impingement heat transfer Confined Unconfined Turbulent flow

#### ABSTRACT

In order to improve the cooling process in impingement synthetic jets, it is necessary to evaluate the influence of dimensionless parameters and the geometry of the flow field and heat transfer. In this paper, the effects of geometry (confined and unconfined impingement synthetic jets), and jet-to-surface spacings, Reynolds number and dimensionless stroke length on a three-dimensional unsteady impingement synthetic jet are studied. For this purpose, two types of turbulence models, namely the  $v^2 - f$  and  $SST/k - \omega$ , have been employed. The simulation results have been indicated that the  $v^2 - f$  model comparing to the  $SST/k - \omega$  model has a close agreement with the available experimental data. The results show that the flow field and the corresponding heat transfer distribution of the impingement synthetic jet are affected by geometry such that an unconfined impingement synthetic jet is more efficient in a cooling process relative to the confined case. Also, increasing jet-to-surface spacing affects the vortex structure and consequently the heat transfer. The stagnation heat transfer rate reaches to the maximum value at an optimum impingement distance as a result of the appropriate ventilation and the coherence vortex structures. The stagnation heat transfer rate experiences several extrema as a result of the impingement continuous jet.

#### 1. Introduction

An appealing characteristic of impingement synthetic jets is the generation of vortex rings during cycles which has attracted researchers' attention in the field of electronic cooling, separation control, jet vectoring and mixing enhancement [1-4]. Synthetic jets are generated by oscillatory motions of an actuator with a certain frequency in a cavity and injected to a medium through an orifice. In a cooling application, an impingement synthetic jet without needing an extra energy source, can produce an oscillatory flow and has a much better efficiency in comparison with an ordinary impingement jet. The impingement synthetic jet consists of two strokes for the completion of a cycle: blowing and suction strokes. A vortex ring which is generated in the blowing stroke has to have enough energy to escape during the suction stroke. Then it impinges to the heated surface, moves radially along it and potentially can increase the heat transfer efficiency. In each cycle, a vortex ring is generated and leads to the set of vortex rings which moves toward the heated surface. The design and appearance of the impingement synthetic jets can improve the cooling process such that the impingement synthetic jets based on the appearance and design

are categorized to the confined and unconfined impingement synthetic jets. The confined impingement synthetic jets are enclosed by two walls, around the orifice outlet (Fig. 1a). Although the unconfined impingement synthetic jets are similar to the confined ones, they do not have walls around the orifice outlet (Fig. 1b). However, this minor differentiation can affect extremely the heat transfer rate. A schematic of these types of synthetic jets and their components is shown in Fig. 1.

#### 2. Literature review

Steady and oscillation flows have been considered in versatile applications namely, cooling of electronic devices and gas turbine blades, separation control, vectoring and mixing enhancement [1–6]. In the cooling process, employing steady or oscillation flow has been considered by researchers for many years. In the steady impingement jets and the impingement synthetic jets, several key dimensionless parameters affect the flow field and heat transfer rate. In an experimental study, Jambunathan et al. [7] have collated experimental data of impinging turbulent steady jets with Reynolds number in the range of 5000–124000. They found that existing correlations for local heat

https://doi.org/10.1016/j.ijthermalsci.2018.01.011

<sup>\*</sup> Corresponding author. Ohio University, Department of Mechanical Engineering, 418 Stocker Center, Athens, OH 45701, USA. *E-mail address*: mh861316@ohio.edu (M. Hatami).

Received 9 September 2017; Received in revised form 7 January 2018; Accepted 8 January 2018 1290-0729/ © 2018 Elsevier Masson SAS. All rights reserved.

fFrequency of diaphragm (Hz)(m/s)HJet-to-surface distance (mm) $\delta(r, t)$ Deformation of diaphragm (mm)hCavity height (nun) $\lambda$ Thermal conductivity (W/m. K)h_cHeat transfer coefficient (W/m². K) $\Delta$ Peak to peak displacement at the center of the diaphLOrifice height (mm) $\beta$ Local surface temperature (K)NuNusselt number (= $h_c$ . D/ $\lambda$ ) $\theta_{\infty}$ Ambient temperature (K)q''convective heat flux (W/m²) $v$ Kinematic viscosity (m²/s)rRadial distance from the center of diaphragm (mm) $\varphi$ Phase angle of oscillation (degree)ReReynolds number (= $f_c$ . D/ $U_0$ ) $\pi$ Pi number	Nomenclature		Т	Time period of cycle ( <sup>s</sup> )
DOrifice diameter $\binom{nm}{r}$ $U_0$ Reference velocity $\binom{m/s}{r}$ $D_T$ Uniform temperature wall diameter $\binom{mm}{r}$ $u(r, t)$ Instantaneous velocity at the diaphragm inlet bour $\binom{m/s}{r}$ $f$ Frequency of diaphragm $(Hz)$ $u(r, t)$ Instantaneous velocity at the diaphragm inlet bour $\binom{m/s}{r}$ $H$ Jet-to-surface distance $\binom{nm}{r}$ $\delta(r, t)$ Deformation of diaphragm $\binom{mm}{r}$ $h$ Cavity height $\binom{mm}{r}$ $\lambda$ Thermal conductivity $(W/m, K)$ $h_c$ Heat transfer coefficient $(W/m^2, K)$ $\Delta$ Peak to peak displacement at the center of the diaphragm $L$ Orifice height $\binom{mm}{r}$ $\theta$ Local surface temperature $(K)$ $L_0$ Stroke length of synthetic jet $\binom{mm}{r}$ $\theta_{\infty}$ Ambient temperature $(K)$ $Nu$ Nusselt number $(=h_c, D/\lambda)$ $\theta_{\infty}$ Ambient temperature $(K)$ $q''$ convective heat flux $(W/m^2)$ $v$ Kinematic viscosity $(m^2/s)$ $r$ Radial distance from the center of diaphragm $\binom{mm}{r}$ $\varphi$ Phase angle of oscillation (degree) $Re$ Reynolds number $(=f_0, D/v_0)$ $\pi$ Pi number $Sr$ Strouhal number $(=f, D/U_0)$ $\pi$ Pi number			t	Time ( <sup>s</sup> )
$D_T$ Uniform temperature wall diameter (mm) $u(r, t)$ Instantaneous velocity at the diaphragm inlet bour (m/s) $f$ Frequency of diaphragm (Hz) $u(r, t)$ Instantaneous velocity at the diaphragm inlet bour (m/s) $H$ Jet-to-surface distance (num) $\delta(r, t)$ Deformation of diaphragm (mm) $h$ Cavity height (num) $\lambda$ Thermal conductivity ( $W/m$ . $K$ ) $h_c$ Heat transfer coefficient ( $W/m^2$ . $K$ ) $\Delta$ Peak to peak displacement at the center of the diaphragm $L$ Orifice height (num) $\partial$ Local surface temperature ( $K$ ) $L_0$ Stroke length of synthetic jet (mm) $\theta$ Local surface temperature ( $K$ ) $Nu$ Nusselt number ( $=h_c$ . $D/\lambda$ ) $\theta_{\infty}$ Ambient temperature ( $K$ ) $q''$ convective heat flux ( $W/m^2$ ) $v$ Kinematic viscosity ( $m^2/s$ ) $r$ Radial distance from the center of diaphragm (num) $\varphi$ Phase angle of oscillation (degree) $Re$ Reynolds number ( $=f_c$ . $D/U_0$ ) $\pi$ Pi number	$D_c$	Cavity diameter ( <sup>mm</sup> )	U(t)	Instantaneous velocity (m/s)
fFrequency of diaphragm (Hz)(m/s)HJet-to-surface distance (num) $\delta(r, t)$ Deformation of diaphragm (num)hCavity height (num) $\lambda$ Thermal conductivity (W/m. K)h_cHeat transfer coefficient (W/m <sup>2</sup> . K) $\Delta$ Peak to peak displacement at the center of the diaphLOrifice height (num) $\theta$ Local surface temperature (K)NuNusselt number (= $h_c$ . $D/\lambda$ ) $\theta_{\infty}$ Ambient temperature (K)q''convective heat flux (W/m <sup>2</sup> ) $v$ Kinematic viscosity (m <sup>2</sup> /s)rRadial distance from the center of diaphragm (num) $\varphi$ Phase angle of oscillation (degree)ReReynolds number (= $f_c$ . $D/U_0$ ) $\pi$ Pi number	D	Orifice diameter ( <sup>mm</sup> )	$U_0$	Reference velocity ( <i>m</i> / <i>s</i> )
HJet-to-surface distance (mm) $\delta(r, t)$ Deformation of diaphragm (mm)hCavity height (nm) $\lambda$ Thermal conductivity (W/m. K)h_cHeat transfer coefficient (W/m². K) $\Delta$ Peak to peak displacement at the center of the diaphLOrifice height (mm) $\Delta$ Local surface temperature (K)L_0Stroke length of synthetic jet (mm) $\theta$ Local surface temperature (K)NuNusselt number (= $h_c$ . $D/\lambda$ ) $\theta_{\infty}$ Ambient temperature (K)q"convective heat flux (W/m²) $v$ Kinematic viscosity (m²/s)rRadial distance from the center of diaphragm (mm) $\varphi$ Phase angle of oscillation (degree)ReReynolds number (= $f_0$ . $D/U_0$ ) $\pi$ Pi number	$D_T$	Uniform temperature wall diameter (mm)	u(r, t)	Instantaneous velocity at the diaphragm inlet boundary
hCavity height (mm) $\lambda$ Thermal conductivity (W/m. K)h_cHeat transfer coefficient (W/m². K) $\Delta$ Peak to peak displacement at the center of the diaphLOrifice height (mm) $\Delta$ Peak to peak displacement at the center of the diaphL_0Stroke length of synthetic jet (mm) $\theta$ Local surface temperature (K)NuNusselt number $(=h_c. D/\lambda)$ $\theta_{\infty}$ Ambient temperature (K)q"convective heat flux (W/m²) $v$ Kinematic viscosity (m²/s)rRadial distance from the center of diaphragm (mm) $\varphi$ Phase angle of oscillation (degree)ReReynolds number $(=f_0. D/v_0)$ $\pi$ Pi numberSrStrouhal number $(=f. D/U_0)$ $\pi$ Pi number	f	Frequency of diaphragm (Hz)		(m/s)
$h_c$ Heat transfer coefficient $(W/m^2, K)$ $\Delta$ Peak to peak displacement at the center of the diaph $(^{num})$ $L$ Orifice height $(^{mm})$ $\Delta$ Peak to peak displacement at the center of the diaph $(^{num})$ $L_0$ Stroke length of synthetic jet $(^{num})$ $\theta$ Local surface temperature $(K)$ $Nu$ Nusselt number $(=h_c. D/\lambda)$ $\theta_{\infty}$ Ambient temperature $(K)$ $q''$ convective heat flux $(W/m^2)$ $v$ Kinematic viscosity $(m^2/s)$ $r$ Radial distance from the center of diaphragm $(^{num})$ $\varphi$ Phase angle of oscillation (degree) $Re$ Reynolds number $(=U_0. D/v)$ $\pi$ Pi number $Sr$ Strouhal number $(=f. D/U_0)$ $\pi$	H	Jet-to-surface distance (mm)	$\delta(r, t)$	Deformation of diaphragm ( <sup>mm</sup> )
LOrifice height (mm)(mm) $L_0$ Stroke length of synthetic jet (mm) $\theta$ Local surface temperature (K) $Nu$ Nusselt number $(=h_c. D/\lambda)$ $\theta_{\infty}$ Ambient temperature (K) $q''$ convective heat flux $(W/m^2)$ $v$ Kinematic viscosity $(m^2/s)$ $r$ Radial distance from the center of diaphragm (mm) $\varphi$ Phase angle of oscillation (degree) $Re$ Reynolds number $(=U_0. D/v)$ $\pi$ Pi number $Sr$ Strouhal number $(=f. D/U_0)$ $\pi$ $Pi$ number	h	Cavity height ( <sup>mm</sup> )	λ	Thermal conductivity $(W/m. K)$
$L_0$ Stroke length of synthetic jet ( $^{mm}$ ) $\theta$ Local surface temperature (K) $Nu$ Nusselt number ( $=h_c$ . $D/\lambda$ ) $\theta_{\infty}$ Ambient temperature (K) $q''$ convective heat flux ( $W/m^2$ ) $v$ Kinematic viscosity ( $m^2/s$ ) $r$ Radial distance from the center of diaphragm ( $^{mm}$ ) $\varphi$ Phase angle of oscillation (degree) $Re$ Reynolds number ( $=U_0$ . $D/v$ ) $\pi$ Pi number $Sr$ Strouhal number ( $=f$ . $D/U_0$ ) $\pi$ Pi number	$h_c$	Heat transfer coefficient $(W/m^2, K)$	Δ	Peak to peak displacement at the center of the diaphragm
NuNusselt number $(=h_c. D/\lambda)$ $\theta_{\infty}$ Ambient temperature $(K)$ $q''$ convective heat flux $(W/m^2)$ $v$ Kinematic viscosity $(m^2/s)$ $r$ Radial distance from the center of diaphragm $(^{num})$ $\varphi$ Phase angle of oscillation (degree) $Re$ Reynolds number $(=U_0. D/v)$ $\pi$ Pi number $Sr$ Strouhal number $(=f. D/U_0)$ $\pi$ Pi number	L	Orifice height ( <sup>mm</sup> )		( <sup>mm</sup> )
$q''$ convective heat flux $(W/m^2)$ $v$ Kinematic viscosity $(m^2/s)$ $r$ Radial distance from the center of diaphragm $\binom{nm}{}$ $\varphi$ Phase angle of oscillation (degree) $Re$ Reynolds number $(=U_0. D/v)$ $\pi$ Pi number $Sr$ Strouhal number $(=f. D/U_0)$ T $\pi$	$L_0$	Stroke length of synthetic jet ( <sup>mm</sup> )	θ	Local surface temperature (K)
rRadial distance from the center of diaphragm ( $^{mm}$ ) $\varphi$ Phase angle of oscillation (degree)ReReynolds number (= $U_0$ . $D/v$ ) $\pi$ Pi numberSrStrouhal number (= $f$ . $D/U_0$ ) $\pi$ Pi number	Nu	Nusselt number $(=h_c. D/\lambda)$	$ heta_{\infty}$	Ambient temperature (K)
ReReynolds number (= $U_0$ . $D/v$ ) $\pi$ Pi numberSrStrouhal number (= $f$ . $D/U_0$ )	q''	convective heat flux $(W/m^2)$	υ	Kinematic viscosity $(m^2/s)$
Sr Strouhal number $(=f. D/U_0)$	r	Radial distance from the center of diaphragm (mm)	arphi	Phase angle of oscillation (degree)
	Re	Reynolds number $(=U_0. D/v)$	π	Pi number
St Stokes number $(=\sqrt{f/D^2/n})$	Sr	Strouhal number $(=f. D/U_0)$		
$(-\sqrt{1}, 2, 7)$	St	Stokes number $(=\sqrt{f. D^2/v})$		

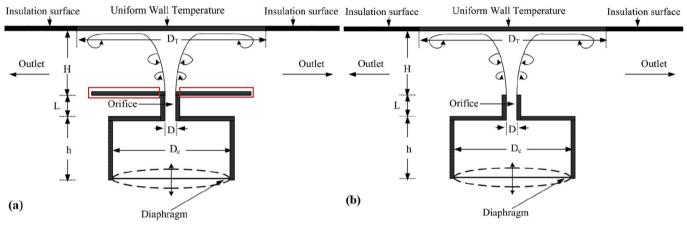


Fig. 1. A schematic of (a) confined (b) unconfined impingement synthetic jets.

transfer coefficient express Nusselt number as a function of Reynolds number raised to a constant exponent. However, the available empirical data suggest that this exponent should be a function of nozzle-to-plate spacing and of the radial displacement from the stagnation point. It is also suggested that the Nusselt number is independent of nozzle-toplate spacing up to a value of 12 nozzle diameters at radii greater than six nozzle diameters from the stagnation point. Hoogendoorn [8], in an experimental study on heat transfer of the impinging steady jet, has shown the effect of turbulence on the stagnation zone. It has been reported that when a low turbulent jet is examined a Nusselt number value slightly higher than the stagnation point is observed. Colucci and Viskanta [9] have studied the effect of hyperbolic nozzle geometry on the local heat transfer coefficient for the confined impinging steady jets. Their experiments have been performed at the small nozzle to plate distances (between 0.25 and 6) and Reynolds numbers in the range of 10000-50000 on a uniformly heated impingement surface. It has been concluded that the local heat transfer coefficient for confined jets is more sensitive to Reynolds number and nozzle to plate spacing than that for unconfined jets. Xie et al. [10] have evaluated the application of ribs in gas turbine blade cooling process. For fixed mainstream and cooling flow Reynolds numbers, three ribs including continuous rib, centrally truncated rib and laterally truncated rib at two different blowing ratios (i.e., 0.5 and 1.0) are considered. They have reported that at lower blowing ratio, the result ensuing for each rib is identical but significantly outperforms the case without ribs. At higher blowing ratio, the heat transfer rate for the cases with ribs is lower than the case without ribs. Increasing efficiency with a synthetic jet firstly has been reported by Campbell et al. [1]. It has been shown through

experimental studies that synthetic jets can improve the cooling process of a laptop processor. With an optimum combination of design parameters, they have shown that the synthetic jet can decrease the processor temperature rise by 22% in comparison with the uncontrolled case (i.e. without a synthetic jet). Nevins and Ball [11] have examined an oscillation impingement jet in a limited range of frequencies. The experiments have been repeated for several forms of waves: sinusoidal, quadratic, and triangular, however, in all cases, a significant impact has not been observed in the heat transfer rate. In the recent years, utilization of the synthetic jet for the improvement of heat transfer has grown significantly. The cooling process on a constant heat flux wall has been investigated by using a synthetic jet and the experimental results have been compared with a continuous jet [12]. In this study, the influence of several parameters namely, Reynolds number, orifice to surface distances and high and low frequencies have been examined. The results have shown that for small distances at high oscillation frequencies, the heat transfer coefficient is higher than low oscillation frequencies, whereas low frequencies have a better performance for lager distances. The performance of synthetic jets at small jet-to-surface spacings on a constant temperature surface has been addressed by McGuinn et al. [13]. The experimental observations have shown that the mean heat transfer distribution has a secondary peak for low jet-tosurface spacings due to high turbulent flow in the wall jet boundary layer. In another experimental study, Vukasinovic and Glezer [14] have proved that maximum heat transfer is not in the stagnation point; rather it is in a distance about half of the orifice radius. It has also been concluded that the heat transfer coefficient increases as a result of reducing jet-to-surface spacings and increasing frequency. The influence Download English Version:

# https://daneshyari.com/en/article/7060750

Download Persian Version:

https://daneshyari.com/article/7060750

Daneshyari.com