



Effect of die exit stress state, Deborah number, uniaxial and planar extensional rheology on the neck-in phenomenon in polymeric flat film production

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ABSTRACT

In this work, effect of the second to first normal stress difference ratio at the die exit, $-N_2/N_1$, uniaxial extensional strain hardening, $\frac{\eta_{E,U,max}}{3\eta_0}$, planar-to-uniaxial extensional viscosity ratio, $\frac{\eta_{E,P}}{\eta_{E,U}}$, and Deborah number, De , has been investigated via viscoelastic isothermal modeling utilizing 1D membrane model and a single-mode modified Leonov model as the constitutive equation. Based on the performed parametric study, it was found that there exists a threshold value for De and $\frac{\eta_{E,U,max}}{3\eta_0}$, above which, the neck-in starts to be strongly dependent on $-N_2/N_1$. It was found that such critical De decreases if $-N_2/N_1$, $\frac{\eta_{E,U,max}}{3\eta_0}$ increases and/or $\frac{\eta_{E,P}}{\eta_{E,U}}$ decreases. Numerical solutions of the utilized model were successfully approximated by a dimensionless analytical equation relating the normalized maximum attainable neck-in with $\frac{\eta_{E,U,max}}{3\eta_0}$, $\frac{\eta_{E,P}}{\eta_{E,U}}$, $-N_2/N_1$ and De . Suggested equation was tested by using literature experimental data considering that $-N_2/N_1$ depends on die exit shear rate, temperature and utilized constitutive model parameters for given polymer melt. It was found that approximate model predictions are in a very good agreement with the corresponding experimental data for low as well as very high Deborah numbers, at which neck-in strongly depends on $-N_2/N_1$. It is believed that the obtained knowledge together with the suggested simple model can be used for optimization of the extrusion die design (influencing flow history and thus die exit stress state), molecular architecture of polymer melts and processing conditions to suppress neck-in phenomenon in production of very thin polymeric flat films.

1. Introduction

The extrusion film casting technology is an industrially important process that has a firm position on the market due to its capability to produce high quality thin polymeric films at high production rates. Those films can be used in different applications such as wrapping materials, barriers reducing permeability for air and vapor, or as a separator membrane for rechargeable batteries in mobile devices and electric vehicles.

In this technology, the polymer melt is extruded through a slit die to form a thick sheet that is subsequently intensively stretched in the axial direction, hauled off and quenched by a rotating drum stabilizing the film dimensions. [1,2] (Fig. 1). Except of an initial swelling, the thickness of the film decreases monotonically due to high Draw ratio (the haul off speed divided by the die exit velocity). In such a case, film width is gradually reduced toward the stretching/cooling roll, which is called neck-in phenomenon.

Aside from that, the interrelated defect of edge-beads promotes the lateral portion of the film to being substantially thicker than its central part (Fig. 2). While the first phenomenon affects a production rate, the second calls for a post-production trimming since solely central portion of the film is uniform in thickness. Understanding the relationship between material parameters and processing conditions including the flow in the die might be the way to effectively control the extent of neck-in and edge-beading as even a small reduction of these defects may bring increased efficiency considering high production rates in this manufacturing process.

Scientific investigation of the extrusion film casting process has been addressed in many works dealing with both a steady and transient approaches to the problem. Initial studies were dedicated to an investigation of the hydrodynamic instability observed during the production of fibers called draw resonance [3–5] and then expanded for films in [6] where the numerical modeling for film casting using the one-dimensional isothermal model of Newtonian fluid was utilized for

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List of symbols

A	aspect ratio (1)	$u(X)$	(mm·s ⁻¹)
A_1, A_2	fitting parameters of analytical model (1)	u_0	chill roll speed (mm·s ⁻¹)
b	dissipation term (s ⁻¹)	\bar{u}	axial velocity component at the die exit (mm·s ⁻¹)
\bar{b}	dimensionless dissipation term (1)		dimensionless axial velocity component of the film at any x location (1)
\underline{c}	recoverable Finger tensor (1)	v	transverse velocity component of the film at any x location (mm·s ⁻¹)
\underline{c}^{-1}	inverse recoverable Finger tensor (1)	W	elastic potential (Pa)
$\underline{\underline{c}}_0$	jaumann (corotational) time derivative of the recoverable Finger strain tensor (s ⁻¹)	w	thickness velocity component of the film at any x location (mm·s ⁻¹)
c_{xx}	normal component of the recoverable Finger tensor in axial x -direction (1)	X	take-up length (stretching distance, air gap) (mm)
c_{yy}	normal component of the recoverable Finger tensor in transverse y -direction (1)	x	position in axial x -direction (mm)
c_{zz}	normal component of the recoverable Finger tensor in thickness z -direction (1)	\bar{x}	dimensionless position in axial x -direction (1)
\underline{D}	deformation rate tensor (s ⁻¹)	y_i	observed value (1)
De	Deborah number (1)	\hat{y}_i	predicted value (1)
DR	draw ratio (1)	x, y, z	spatial coordinates in axial, transverse and thickness direction, respectively (1)
\underline{e}_p	irreversible rate of strain tensor (s ⁻¹)	Z_x, Z_y, Z_z	substitution variables (1)
E	dimensionless take-up force (1)	$\left(\frac{du}{dx}\right)_M$	mean value of extensional strain rate in the air gap (s ⁻¹)
E_a	flow activation energy (J·mol ⁻¹)	$\frac{dc_{xx}}{d\bar{x}}, \frac{dc_{yy}}{d\bar{x}}, \frac{dc_{zz}}{d\bar{x}}$	derivative of Finger tensor components with respect to dimensionless \bar{x} position (1)
e	half-thickness of the film at any x location (mm)	$\frac{d\bar{u}}{d\bar{x}}, \frac{d\bar{v}}{d\bar{x}}, \frac{d\bar{w}}{d\bar{x}}$	derivative of dimensionless axial, transverse and thickness velocity component with respect to dimensionless \bar{x} position (1)
e_0	die half-gap (half-thickness of the film at the die exit) (mm)		
\bar{e}	dimensionless half-thickness of the film at any x location (1)	Greek symbols	
F	take-up force (stretching force) (N)	α	Arrhenius law parameter (K)
f	rate of deformation in transverse y -direction (s ⁻¹)	α_1, α_2	fitting parameters of analytical model (1)
G	linear Hookean elastic modulus (Pa)	β	non-linear Leonov model parameter (1)
g	rate of deformation in thickness z -direction (s ⁻¹)	$\dot{\gamma}_{COR}$	corrected shear rate by Rabinowitsch correction for the rectangle channel (s ⁻¹)
i	index i , noting the spatial direction (1)	$\underline{\delta}$	unit tensor (Kronecker delta) (1)
$I_{1, c}$	first invariant of recoverable Finger tensor (1)	δ	δ shift function (1)
$I_{2, c}$	second invariant of recoverable Finger tensor (1)	η_0	Newtonian viscosity (Pa·s)
j	index j (1)	$\eta_{E, P}$	steady planar extensional viscosity (Pa·s)
L	half-width of the film at any x location (mm)	$\eta_{E, P, \max}$	maximal steady planar extensional viscosity (Pa·s)
L_0	half-width of the die (half-width of the film at the die exit) (mm)	$\eta_{E, U}$	steady uniaxial extensional viscosity (Pa·s)
\bar{L}	dimensionless half-width of the film at any x location (1)	$\eta_{E, U, \max}$	maximal steady uniaxial extensional viscosity (Pa·s)
MFI	melt flow index (g/10 min)	θ	fitting parameters of analytical model (1)
MFR	mass flow rate (kg·h ⁻¹)	λ	relaxation time (s)
M_n	number average molar mass (g·mol ⁻¹)	ν	non-linear Leonov model parameter (1)
M_w	mass average molar mass (g·mol ⁻¹)	ξ	non-linear Leonov model parameter (1)
NI	maximum attainable neck-in (mm)	ρ	polymer density (g·cm ⁻³)
NI^*	normalized maximum attainable neck-in (1)	$\underline{\tau}$	extra stress tensor (Pa)
N_1	first normal stress difference (Pa)	τ_{xx}	normal stress in axial x -direction (Pa)
N_2	second normal stress difference (Pa)	τ_{yy}	normal stress in transverse y -direction (Pa)
n	non-linear Leonov model parameter (1)	τ_{zz}	normal stress in thickness z -direction (Pa)
n_0	non-Newtonian index (1)	$\bar{\tau}_{xx}$	dimensionless normal stress in axial x -direction (1)
n_s	number of sample points (1)	$\bar{\tau}_{yy}$	dimensionless normal stress in transverse y -direction (1)
Q	volumetric flow rate (m ³ ·s ⁻¹)	$\bar{\tau}_{zz}$	dimensionless normal stress in thickness z -direction (1)
R	gas constant (J·K ⁻¹ ·mol ⁻¹)	φ_1, φ_2	fitting parameters of analytical model (1)
T	melt temperature at the die (°C)	$\psi_1, \psi_2, \psi_3, \psi_4$	fitting parameters of analytical model (1)
T_0	reference temperature in the Arrhenius law (°C)		
u	axial velocity component of the film at any x location		

the first time. Other authors followed and employed different constitutive equations for power-law [7], and viscoelastic fluids using modified convected-Maxwell [8] and modified Giesekus model [9,10]. Due to the assumed kinematics for the free surface flow at the drawing zone, the model could not capture edge-bead defect and contraction in film width.

First efforts to overcome this limitation and to accommodate ability to predict neck-in were made for a Newtonian fluid at isothermal [11]

and non-isothermal conditions [12–14]. Lately, improved isothermal two-dimensional membrane model having the capability to capture the development of edge-beads under the stationary conditions was released; isothermal Newtonian [15] and viscoelastic Maxwell and Giesekus fluid [16], and models that take thermal effects into account for Newtonian [17] and viscoelastic Larson fluid [18]. In the meantime, simplified one-dimensional membrane approach based on a supplementary kinematic hypothesis, that originally brought for a float glass

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