



Dynamic effective thickness in laminated-glass beams and plates



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ABSTRACT

In recent years, several equations have been proposed to calculate deflections and stresses in laminated-glass beams and plates under static loading using the concept of effective thickness, which consists of calculating the thickness of a monolithic element with equivalent bending properties to a laminated element. Recently, an effective thickness for the dynamic behavior of laminated-glass beams has been proposed to enable the modal parameters (natural frequencies, loss factors and mode shapes) to be determined using an equivalent monolithic model. In the present paper, the technique has been extended to the two-dimensional case of rectangular laminated-glass plates and the steps needed to estimate the modal parameters of laminated-glass elements using this methodology are presented. The dynamic effective thickness concept has been validated by experimental tests made on a laminated-glass beam and a laminated-glass plate. The results show that good accuracy is achieved in the natural frequencies and mode shapes but high scatter is encountered in the loss factors.

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1. Introduction

Laminated glass is a sandwich or layered material consisting of two or more plies of monolithic glass with one or more interlayers of a polymeric material. Thus, a laminated-glass element (beam, plate, etc.) is a composite material which combines the properties of the glass with the benefits of a highly elastic polymeric material, i.e. the structural behavior of laminated glass is that of a composite structure. The glass layers may be either normal annealed, heat-strengthened, chemically strengthened or tempered glass. When the cross-section is made of different glasses (e.g. one layer of annealed glass and the other of tempered glass) it is called a hybrid laminated-glass element. However, the treatments affect the ultimate strength but not the Young modulus, and therefore no distinction is made concerning the type of glass if the calculations are made prior to glass breakage. All polymeric interlayers are viscoelastic in nature [1], i.e. their mechanical properties are frequency (or time) and temperature dependent. Polyvinyl butyral (PVB) is the most commonly used interlayer material and is marketed in thicknesses of 0.38 mm or a multiple of this value (0.76 mm, 1.12 mm, 1.52 mm). However, the new ionoplastic interlayers improve the mechanical properties of laminated glass and maintains a significant advantage (higher stiffness and strength) over the PVB for a large range of temperatures [1]. This

interlayer material is now in flat sheet form, in thicknesses of 0.89, 1.52, 2.27, and 3.05 mm, and as rolled sheeting, at 0.89 mm thickness. The simplest laminated-glass configuration consists of three layers: two monolithic glass plies and a polymeric core (see Fig. 1).

The response of laminated-glass elements varies between two borderlines [2]: (1) The layered limit corresponding to the case when the beam consists of free-sliding glass plies and (2) the monolithic limit, when the Euler–Bernoulli assumptions hold (plane sections remain plane) for the entire section of the laminated-glass element (the response of the composite beam approaches that of a homogeneous glass beam with an equal cross-section) [3,4]. As the tensile modulus of the PVB is far less in comparison with that corresponding to glass, significant transverse shear appears in the viscoelastic layer [1,8,10].

In the analytical and numerical models, glass mechanical behavior is usually modeled as linear-elastic prior to glass breakage, whereas the polymeric interlayer is characterized as linear-viscoelastic. Laminated glass is easy to assemble in a finite-element model but many small 3D elements are needed to mesh accurately, which, on the other hand, are very high time consuming. In the last few years, some papers have been published on the calculation of laminated-glass elements, examining the concept of effective thickness [1,3,4,8]. The method consists of calculating the thickness of a monolithic element with bending properties equivalent to those of the laminated one. The effective thickness can then be used in analytical equations and simplified finite-element models instead of the laminated-glass element [9].

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Nomenclature

| | | | |
|-----------------|--|--------------------------|--|
| D | flexural stiffness in plates | T | temperature |
| E | Young modulus | T_0 | reference temperature |
| E_{eff} | effective Young modulus | Y | $\frac{H_2^2 E_1 H_1 E_3 H_3}{E I_T (E_1 H_1 + E_3 H_3)}$ |
| E_1 | Young's modulus of glass layer 1 | Y_1 | $\frac{12 H_0^2 H_1 H_3}{(H_1^3 + H_3^3)(H_1 + H_3)}$ |
| E_3 | Young's modulus of glass layer 3 | | |
| $E I^*$ | complex flexural stiffness in beams | <i>Lowercase letters</i> | |
| $E_2^*(\omega)$ | complex tensile modulus for the polymeric interlayer | a_T | shift factor |
| $E'(\omega)$ | real component of the tensile complex modulus (storage) | b | width of a glass beam |
| $E''(\omega)$ | imaginary component of the tensile complex modulus (loss) | e_i | modulus coefficient in Prony's series viscoelastic model |
| $E I_T$ | $E_1 I_1 + E_3 I_3$ | $g(x)$ | shape function (Galuppi and Royer Carfagni model) |
| $E_2(t)$ | viscoelastic relaxation tensile modulus for polymeric interlayer | g_M^* | shear parameter M&Ms model |
| E_∞^2 | equilibrium tensile modulus for the polymeric interlayer | g_R^* | shear parameter RKU model |
| $G_2(t)$ | viscoelastic relaxation shear modulus for the polymeric interlayer | i | imaginary unit |
| $G_2^*(\omega)$ | complex shear modulus for the polymeric interlayer | k_l | wavenumber |
| G_0 | glassy shear modulus | k_M^* | complex wave number |
| H_1 | thickness of glass layer 1 in laminated glass | \bar{m} | mass per unit area |
| H_2 | thickness of polymeric layer in laminated glass | t | time |
| H_3 | thickness of glass layer 3 in laminated glass | w | deflection |
| H_0 | $H_2 + \left(\frac{H_1 + H_3}{2}\right)$ | | |
| I | second moment of area | <i>Greek symbols</i> | |
| I_1 | $\frac{H_1^3}{12}$ | Ω^* | non-dimensional complex frequency |
| I_3 | $\frac{H_3^3}{12}$ | β | buckling ratio for a beam |
| I_T | $I_1 + I_3 = \frac{H_1^3 + H_3^3}{12}$ | ζ | modal damping ratio |
| $K(t)$ | viscoelastic bulk modulus | η | loss factor |
| L | length of a glass beam | η_2 | loss factor of the polymeric interlayer of laminated glass |
| | | ν_i | Poisson ratio of the i -th glass layer |
| | | ρ_i | mass density of the i -th glass layer |
| | | τ_i | time coefficient in Prony's series viscoelastic model |
| | | ω | frequency |

The aim of the present paper is to propose a simplified method to estimate the modal parameters of rectangular laminated-glass plates while avoiding the use of finite-element models or complicated analytical models. The method is based on the dynamic effective thickness proposed in a previous paper [9] for laminated-glass beams, which is here extended to the two-dimensional case of rectangular laminated-glass plates. An alternative to the effective thickness is the concept of effective Young modulus, which can be used interchangeably for laminated-glass elements with the same accuracy. This technique can be applied to three-layered laminated-glass plates with glass showing a linear elastic behavior and the polymeric core showing viscoelastic behavior. Thus the glass layers can be made of different types of glass (annealed, tempered, heat-strengthened, etc.), and the traditional cores (PVB, ionoplastic, etc.) can be considered in this model. In this paper, the modal parameters (natural frequencies, loss factors, and mode shapes) of a 1400 × 1000 × 16 mm laminated-glass plate pin-supported at the four corners, and of a beam 1 m long and 12 mm thick, both the beam and the plate with annealed glass plies and PVB core, were estimated using the effective thickness concept. For the validation of the model, operational modal tests

were performed on the beam and the plate, and the modal parameters identified from the experimental responses were compared with those predicted using the effective thickness concept.

2. State of the art

2.1. Viscoelastic behavior

The mechanical properties of a linear-viscoelastic material are frequency (or time) and temperature dependent [11]. In the frequency domain, the complex tensile modulus, $E_2^*(\omega)$, at temperature T is given by:

$$E_2^*(\omega, T) = E_2'(\omega, T) + i \cdot E_2''(\omega, T) = E_2'(\omega, T)(1 + i \cdot \eta_2(\omega, T)) \quad (1)$$

where superscript ‘*’ indicates complex, ω represents the frequency, i is the imaginary unit, $E_2'(\omega, T)$ and $E_2''(\omega, T)$ are the storage and the loss tensile moduli, respectively, and

$$\eta_2(\omega) = \frac{E_2''(\omega, T)}{E_2'(\omega, T)} \quad (2)$$

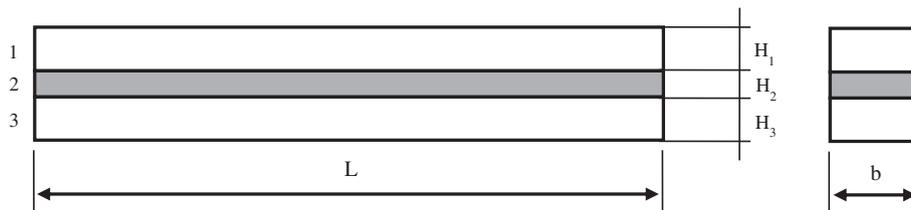


Fig. 1. Laminated glass.

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