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Optimal convergence rates for the strong solutions to the compressible Navier–Stokes equations with potential force

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1. Introduction

This paper is concerned with the Cauchy problem of the full compressible Navier–Stokes equations affected by the external potential force in \mathbb{R}^3 :

$$\begin{cases} \rho_t + \nabla \cdot (\rho u) = 0, \\ \rho[u_t + (u \cdot \nabla)u] + \nabla P(\rho, \theta) = \mu \Delta u + (\mu + \mu') \nabla (\nabla \cdot u) + \rho F, \\ \rho c_V[\theta_t + (u \cdot \nabla)\theta] + \theta P_{\theta}(\rho, \theta) \nabla \cdot u = \kappa \Delta \theta + \Psi[u], \end{cases}$$
(1.1)

and the initial data

$$(\rho, u, \theta)(0, x) = (\rho_0, u_0, \theta_0)(x) \to (\rho_\infty, 0, \theta_\infty), \quad \text{as } |x| \to \infty.$$

$$(1.2)$$

Here the unknown functions $\rho > 0$, $u = (u_1, u_2, u_3)$, and θ denote the density, the velocity and the temperature; $x = (x_1, x_2, x_3) \in \mathbb{R}^3$ is the space variable, t > 0 is the time variable; $P = P(\rho, \theta)$, μ , μ' , $\kappa > 0$, and c_V are the pressure, the first and second viscosity coefficients, the coefficient of heat conduction,

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In this paper, we consider the effect of external force on the large-time behavior of solutions to the Cauchy problem for the three-dimensional full compressible Navier–Stokes equations. We construct the global unique solution near the stationary profile to the system for the small $H^2(\mathbb{R}^3)$ initial data. Moreover, the optimal L^p-L^2 $(1 \le p \le 2)$ time decay rates of the solution to the system are established via a low frequency and high frequency decomposition.

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and the specific heat at constant volume, respectively. In addition, F = F(x) is an external force and $\Psi = \Psi[u]$ is the dissipation function:

$$\Psi[u] = \frac{\mu}{2} \sum_{i,j=1}^{3} (\partial_i u_j + \partial_j u_i)^2 + \mu' \sum_{j=1}^{3} (\partial_j u_j)^2.$$
(1.3)

Throughout this paper, we assume that the above physical parameters satisfy $\mu > 0$ and $2\mu + 3\mu' \ge 0$ which deduce $\mu + \mu' > 0$. ρ_{∞} and θ_{∞} are positive constants, and $P(\rho, \theta)$ is smooth in a neighborhood of $(\rho_{\infty}, \theta_{\infty})$ with $P_{\rho}(\rho_{\infty}, \theta_{\infty}) > 0$ and $P_{\theta}(\rho_{\infty}, \theta_{\infty}) > 0$.

In this work, we only consider the potential force, that is, $F = -\nabla \Phi(x)$. Under aforementioned assumptions, the existence of the stationary solution to the problem (1.1) and (1.2) has been established in [1]. The solution (ρ_*, u_*, θ_*) in a neighborhood of $(\rho_\infty, 0, \theta_\infty)$ is given by

$$\int_{\rho_{\infty}}^{\rho_*(x)} \frac{P_{\rho}(\eta, \theta_{\infty})}{\eta} \mathrm{d}\eta + \Phi(x) = 0, \qquad u_*(x) = 0, \qquad \theta_*(x) = \theta_{\infty}, \tag{1.4}$$

and satisfies

$$\|\rho_* - \rho_\infty\|_{H^k(\mathbb{R}^3)} \le C \|\Phi\|_{H^k(\mathbb{R}^3)}, \quad 0 \le k \le 4,$$
(1.5)

$$\sum_{k=1}^{4} \|(1+|x|)\nabla^{k}(\rho_{*}-\rho_{\infty})\|_{L^{2}(\mathbb{R}^{3})} \leq C \sum_{k=1}^{4} \|(1+|x|)\nabla^{k}\Phi\|_{L^{2}(\mathbb{R}^{3})}.$$
(1.6)

We will construct the global unique solution to (1.1) near the steady state $(\rho_*, 0, \theta_{\infty})$ when the initial perturbation belongs to the Sobolev space $H^2(\mathbb{R}^3)$. Our main results are stated as the following theorem.

Theorem 1.1. Let $(\rho_0 - \rho_\infty, u_0, \theta_0 - \theta_\infty) \in H^2(\mathbb{R}^3)$, there exists some small constant $\varepsilon > 0$ such that if

$$\|(\rho_0 - \rho_\infty, u_0, \theta_0 - \theta_\infty)\|_{H^2(\mathbb{R}^3)} + \|\Phi\|_{H^4(\mathbb{R}^3)} + \sum_{k=1}^4 \|(1+|x|)\nabla^k \Phi\|_{L^2(\mathbb{R}^3)} \le \varepsilon,$$
(1.7)

then the initial value problem (1.1) and (1.2) admits a unique solution (ρ , u, θ) globally in time which satisfies

$$\begin{aligned} \rho - \rho_* &\in C^0([0,\infty); H^2(\mathbb{R}^3)) \cap C^1([0,\infty); H^1(\mathbb{R}^3)), \\ u, \theta - \theta_\infty &\in C^0([0,\infty); H^2(\mathbb{R}^3)) \cap C^1([0,\infty); L^2(\mathbb{R}^3)) \end{aligned}$$

Moreover, if the initial data $(\rho_0 - \rho_\infty, u_0, \theta_0 - \theta_\infty)$ is bounded in $L^p(\mathbb{R}^3)$ for any given $1 \le p \le 2$, the solution (ρ, u, θ) enjoys the following decay-in-time estimates:

$$\|\nabla(\rho - \rho_*, u, \theta - \theta_\infty)\|_{H^1(\mathbb{R}^3)} \le C(1+t)^{-\frac{3}{2}\left(\frac{1}{p} - \frac{1}{2}\right) - \frac{1}{2}} \quad for \ all \ t \ge 0, \tag{1.8}$$

$$\|(\rho - \rho_*, u, \theta - \theta_\infty)\|_{L^q(\mathbb{R}^3)} \le C(1+t)^{-\frac{3}{2}\left(\frac{1}{p} - \frac{1}{q}\right)} \quad \text{for all } t \ge 0, \ 2 \le q \le 6, \tag{1.9}$$

$$\|\partial_t (\rho - \rho_*, u, \theta - \theta_\infty)\|_{L^2(\mathbb{R}^3)} \le C(1+t)^{-\frac{3}{2}\left(\frac{1}{p} - \frac{1}{2}\right) - \frac{1}{2}} \quad for \ all \ t \ge 0,$$
(1.10)

for some positive constant C.

Remark 1.1. In Theorem 1.1, using the Sobolev imbedding inequalities in Lemma 2.1, (1.7) together with (1.5) and (1.6) yields

$$\|\rho_* - \rho_\infty\|_{H^4(\mathbb{R}^3)} + \sum_{k=1}^3 \|(1+|x|)\nabla^k(\rho_* - \rho_\infty)\|_{L^2(\mathbb{R}^3) \cap L^3(\mathbb{R}^3)} \le C\varepsilon.$$
(1.11)

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