



# The fractional Schrödinger equation with Hardy-type potentials and sign-changing nonlinearities



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## ABSTRACT

We look for solutions to a fractional Schrödinger equation of the following form

$$(-\Delta)^{\alpha/2}u + \left(V(x) - \frac{\mu}{|x|^\alpha}\right)u = f(x, u) - K(x)|u|^{q-2}u \text{ on } \mathbb{R}^N \setminus \{0\},$$

where  $V$  is bounded and close-to-periodic potential and  $-\frac{\mu}{|x|^\alpha}$  is a Hardy-type potential. We assume that  $V$  is positive and  $f$  has the subcritical growth but not higher than  $|u|^{q-2}u$ . If  $\mu$  is positive and small enough we find a ground state solution, i.e. a critical point of the energy being minimizer on the Nehari manifold. If  $\mu$  is negative we show that there is no ground state solutions. We are also interested in an asymptotic behaviour of solutions as  $\mu \rightarrow 0^+$  and  $K \rightarrow 0$ .

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## 1. Introduction

We consider the following nonlinear, fractional Schrödinger equation with external, Hardy-type potential

$$(-\Delta)^{\alpha/2}u + \left(V(x) - \frac{\mu}{|x|^\alpha}\right)u = f(x, u) - K(x)|u|^{q-2}u \text{ on } \mathbb{R}^N \setminus \{0\} \quad (1.1)$$

where  $\alpha \in (0, 2)$ ,  $\mu \in \mathbb{R}$  and  $N > \alpha$ , with  $u \in H^{\alpha/2}(\mathbb{R}^N)$ . The fractional Schrödinger equation arises in many models from mathematical physics, e.g. nonlinear optics, quantum mechanics, nuclear physics (see e.g. [16,30,33,39,40,46,52,56,58,59] and references therein). We focus on the external potential of the form  $V(x) - \frac{\mu}{|x|^\alpha}$ , where  $V \in L^\infty(\mathbb{R}^N)$  is close-to-periodic potential and  $-\frac{\mu}{|x|^\alpha}$  is Hardy-type potential. Note that

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the Hardy-type potential does not belong to the Kato's class, hence it is not a lower order perturbation of the operator  $-\Delta + V(x)$  (see [45]).

The fractional Laplacian can be defined via Fourier multiplier  $|\xi|^\alpha$ , i.e. the operator  $(-\Delta)^{\alpha/2}$ , for a function  $\psi : \mathbb{R}^N \rightarrow \mathbb{R}$ , is given by the Fourier transform by the formula

$$\mathcal{F}\left((-\Delta)^{\alpha/2}\psi\right)(\xi) := |\xi|^\alpha \hat{\psi}(\xi),$$

where

$$\mathcal{F}\psi(\xi) := \hat{\psi}(\xi) := \int_{\mathbb{R}^N} e^{-i\xi \cdot x} \psi(x) dx$$

denotes the usual Fourier transform. When  $\psi : \mathbb{R}^N \rightarrow \mathbb{R}$  is rapidly decaying smooth function, it can be defined by the principal value of the singular integral

$$(-\Delta)^{\alpha/2}\psi(x) = c_{N,\alpha} P.V. \int_{\mathbb{R}^N} \frac{\psi(x) - \psi(y)}{|x - y|^{N+\alpha}} dy, \quad (1.2)$$

where

$$c_{N,\alpha} := \frac{2^\alpha \Gamma\left(\frac{N+\alpha}{2}\right)}{2\pi^{N/2} |\Gamma(-\alpha/2)|} > 0.$$

Here,  $\Gamma$  denotes the Gamma function, i.e. a function defined for complex numbers  $z$  with  $\operatorname{Re}(z) > 0$  by the formula

$$\Gamma(z) := \int_0^\infty x^{z-1} e^{-x} dx$$

and extended to a meromorphic function on the set  $\mathbb{C} \setminus \{0, -1, -2, \dots\}$ . Both definitions of the fractional Laplacian are equivalent, i.e. on  $L^2(\mathbb{R}^N)$  they give operators with common domain and they coincide on this domain (see [35]). It is known that  $(-\Delta)^{\alpha/2}$  reduces to  $-\Delta$  as  $\alpha \rightarrow 2^-$  (see [17]). In this paper we identify  $(-\Delta)^{\alpha/2}$  with the classical Laplace operator  $-\Delta$  for  $\alpha = 2$ . In what follows we will use the following characterization of the fractional Sobolev space, for  $0 < \alpha < 2$ :

$$H^{\alpha/2}(\mathbb{R}^N) := \left\{ u \in L^2(\mathbb{R}^N) : \iint_{\mathbb{R}^N \times \mathbb{R}^N} \frac{|u(x) - u(y)|^2}{|x - y|^{N+\alpha}} dx dy + \int_{\mathbb{R}^N} |u(x)|^2 dx < \infty \right\}$$

with the associated scalar product:

$$H^{\alpha/2}(\mathbb{R}^N) \times H^{\alpha/2}(\mathbb{R}^N) \ni (u, v) \mapsto \iint_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+\alpha}} dx dy + \int_{\mathbb{R}^N} u(x)v(x) dx \in \mathbb{R}.$$

See e.g. [10,17] for more background about the fractional Laplace operator and fractional Sobolev spaces.

Eq. (1.1) describes the behaviour of the so-called standing wave solutions  $\Phi(x, t) = u(x)e^{-i\omega t}$  of the following time-dependent fractional Schrödinger equation

$$i \frac{\partial \Phi}{\partial t} = (-\Delta)^{\alpha/2} \Phi + \left( V(x) - \frac{\mu}{|x|^\alpha} + \omega \right) \Phi - g(x, |\Phi|).$$

Such an equation was introduced by Laskin by expanding the Feynman path integral from the Brownian-like to the Lévy-like quantum mechanical paths (see [36,37]). The time-dependent equation is also intensively studied (see e.g. [28,38]).

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