



Original research article

Optical soliton perturbation with Radhakrishnan–Kundu–Lakshmanan equation by Lie group analysis

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ABSTRACT

This paper conducts Lie group analysis of the perturbed Radhakrishnan–Kundu–Lakshmanan equation to retrieve optical soliton solutions. This powerful integration scheme retrieves bright and dark optical soliton solutions to the model that is studied with power law nonlinearity. The existence criteria of such solitons are also presented.

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1. Introduction

The study of optical solitons and other nonlinear evolution equations is one of the active areas of research in mathematical physics, applied mathematics and telecommunication sciences. There are several mathematical approaches that can be applied to these studies [1–10]. One of the models that studies dispersive optical solitons is the Radhakrishnan–Kundu–Lakshmanan (RKL) equation. This is the familiar nonlinear Schrödinger's equation that is considered with a couple of perturbations that are due to third order dispersion (3OD) and the self-steepening effect. This model has been studied since many years by a variety of methods that led to the retrieval of bright and dark optical solitons [2,3,10].

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These are the exp-function scheme and the method of undetermined coefficients. This paper will apply one of the most powerful mathematical tools to address the RKL equation that appears with power law nonlinearity and nonlinear dispersion. The self-steepening effect and nonlinear dispersion terms are considered with full nonlinearity. Lie symmetry analysis will retrieve bright and dark singular soliton solutions to this model. The details are discussed in the rest of the paper.

1.1. Governing model

The dimensionless form of the generalized RKL equation with power law nonlinearity, in presence of nonlinear dispersion is given by [2,3,10]

$$iq_t + aq_{xx} + b|q|^{2m}q = i \left[\lambda (|q|^{2m}q)_x - \gamma q_{xxx} + \nu (|q|^{2m})_x q \right], \quad (1)$$

where the first term represents the evolution term, while the second and third terms represent the group velocity dispersion and nonlinear terms respectively. In the right hand side, the terms due to λ , γ , ν are known as self-steepening term, 3OD and nonlinear dispersion, respectively, with constant coefficients. The parameter m indicates the case of power law nonlinearity which is usually the general form of Kerr law nonlinearity. It also needs to be noted that for the stability of the soliton solutions of (1), it is necessary to have $0 < m < 2$. The special case for $\nu = 0$ is called the RKL equation.

2. Lie group analysis

In this section, we will perform Lie symmetry analysis [1,4–9] for generalized RKL Eq. (1). For this, firstly let us assume

$$q(x, t) = u(x, t) + iv(x, t). \quad (2)$$

On substituting (2) in (1) and then, separating real and imaginary parts we get

$$-v_t + au_{xx} + bu(u^2 + v^2)^m = -m(\lambda + \nu)(u^2 + v^2)^{m-1}(2uvu_x + 2v^2v_x) - \lambda(u^2 + v^2)^m v_x + \gamma v_{xxx}, \quad (3)$$

and

$$u_t + av_{xx} + bv(u^2 + v^2)^m = m(\lambda + \nu)(u^2 + v^2)^{m-1}(2u^2u_x + 2uvv_x) + \lambda(u^2 + v^2)^m u_x - \gamma u_{xxx}. \quad (4)$$

Let us consider the Eqs. (3) and (4) are invariant under following transformations:

$$\begin{bmatrix} u \\ v \\ x \\ t \end{bmatrix} \rightarrow \begin{bmatrix} u \\ v \\ x \\ t \end{bmatrix} + \epsilon \begin{bmatrix} \eta(x, t, u) \\ \phi(x, t, u) \\ \xi(x, t, u) \\ \tau(x, t, u) \end{bmatrix}, \quad (5)$$

where ϵ is an infinitesimal parameter. On invoking the invariance criterion, the following relation from the coefficients of the first order of ϵ in (3) and (4) are deduced as:

$$\begin{aligned} -\phi^t + a\eta^{xx} + 2mbu(u^2 + v^2)^{m-1}(u\eta + v\phi) + b\eta(u^2 + v^2)^m &= -\lambda((u^2 + v^2)^m \phi^x) \\ -\lambda(2mv_x(u^2 + v^2)^{m-1}(u\eta + v\phi)) - m(\lambda + \nu)((u^2 + v^2)^{m-1}(2uv\eta^x + 2uu_x\phi + 2u_xv\eta + 2v^2\phi^x + 4vv_x\phi)) & \\ -(\lambda + \nu)2m(m-1)(u^2 + v^2)^{m-2}(u\eta + v\phi)(2uvu_x + 2v^2v_x) + \gamma\phi^{xxx}, & \end{aligned} \quad (6)$$

and

$$\begin{aligned} \eta^t + a\phi^{xx} + 2mbv(u^2 + v^2)^{m-1}(u\eta + v\phi) + b\phi(u^2 + v^2)^m &= \lambda((u^2 + v^2)^m \eta^x) \\ + 2m\lambda u_x(u^2 + v^2)^{m-1}(u\eta + v\phi) + m(\lambda + \nu)(u^2 + v^2)^{m-1}(2u^2\eta^x + 4uu_x\eta + 2uv\phi^x + 2u\phi v_x + 2vv_x\eta) & \\ + 2m(m-1)(\lambda + \nu)(u^2 + v^2)^{m-2}(u\eta + v\phi)(2u^2u_x + 2uvv_x) - \gamma\eta^{xxx}, & \end{aligned} \quad (7)$$

where η^t , ϕ^t , η^x , ϕ^x , η^{xx} , ϕ^{xx} , η^{xxx} , ϕ^{xxx} are extended (prolonged) infinitesimals acting on an enlarged space that includes all derivatives of the dependent variables u_t , v_t , u_x , v_x , u_{xx} , v_{xx} , u_{xxx} , v_{xxx} . The infinitesimals are determined from invariance

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