

# Oxide interfacial charge engineering towards normally-off AlN/GaN MOSHEMT



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## ABSTRACT

In this paper a detail insight into the role of oxide/barrier interfacial charges ( $N_{ox}$ ) for shifting the threshold voltage ( $V_T$ ) of AlN/GaN metal oxide semiconductor high electron mobility transistors (MOSHEMTs) is gained. A model is developed for  $V_T$  considering all possible charges arise at different interfaces. To validate the model the proposed device is simulated by considering different insulators and  $N_{ox}$  into account. It is very fascinating to observe that  $V_T$  is highly sensitive towards change in  $N_{ox}$  at higher oxide dimensions, whereas at lower dimensions  $N_{ox}$  has very negligible effect. Normally-off operation can be achieved by increasing or decreasing  $N_{ox}$  in MOSHEMT with  $Al_2O_3$  or  $HfO_2$  as gate dielectric respectively.

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## 1. Introduction

Inherent high gate leakage current in High Electron Mobility Transistors (HEMTs) not only degrades the device performance in terms of gate voltage swing, power added efficiency and noise [1,2] but also responsible for drain current collapse [3]. In case of MOSHEMT an oxide layer is grown on barrier layer before gate metallization, which not only reduces the gate leakage current but also avoids current collapse and improves on-state drain current performance by nearly 40% [4]. The insulator layer can be  $SiO_2$  [5],  $HfO_2$  [6],  $Al_2O_3$  [7],  $Ga_2O_3$  [8],  $Si_3N_4$  [9],  $MgO$  [10],  $Sc_2O_3$  [10] or  $TiO_2$  [11]. Recently AlN/GaN MOSHEMTs have been demonstrating promising performance in high power applications because of high bandgap and polarization property of AlN [12]. It's very substantial that the AlN barrier grown over GaN buffer should be very thin i.e. below critical thickness in order to obtain normally-off operation [13]. This results in lower sheet charge density ( $n_s$ ), and hence higher sheet resistance ( $R_s$ ) [14]. Meanwhile, there is an opportunity of achieving positive threshold voltage with lower  $n_s$  because at zero gate bias ( $V_{gs}$ ) the channel is depleted and a positive  $V_{gs}$  is required to induce the channel. A positive  $V_T$  is obtained by considering sub-critical barrier thickness i.e. only 1.5 nm [15]. For power electronics applications normally-off operation is required to simplify the design of the driving circuit by eliminating the requirement of negative voltage supply and to reduce the switching [16] and off-state standby power loss [17].

In MOSHEMTs, oxide charges formed because of defects generated during oxidation play a vital role in tuning  $V_T$  depending on amount and polarity. The amount of  $N_{ox}$  depends on the deposition technique adopted during oxidation process [18]. There are many literatures evident regarding the impact of  $V_{gs}$  dependent oxide/semiconductor interface trap charges on  $V_T$  [18,19]. However study related to the effect of variation in interfacial and bulk oxide charges for different dielectric materials on shift in  $V_T$  forms the motivation behind the present work. The remaining of the paper is described as follows. Section 2 describes the device structure, simulation setup and model development. In Section 3 results are discussed. Finally in Section 4 the paper is concluded.

## 2. Model development for threshold voltage

The structure considered for simulation is similar to that of fabricated in Ref. [20] except AlGaN barrier layer which is replaced by AlN layer as shown in Fig. 1. The gate bias dependent 2-dimensional electron gas (2DEG) at AlN/GaN interface forms the channel between source and drain. The dimensions of different layers of the device are enlisted in Table 1. The gate width is considered to be by default 1  $\mu m$  in Silvaco TCAD [21]. In this study three oxide layers such as  $SiO_2$ ,  $Al_2O_3$ , and  $HfO_2$  with different dielectric constants and charges have been considered. Similarly the thickness of the oxides also varies from 2 nm to 20 nm during the study.

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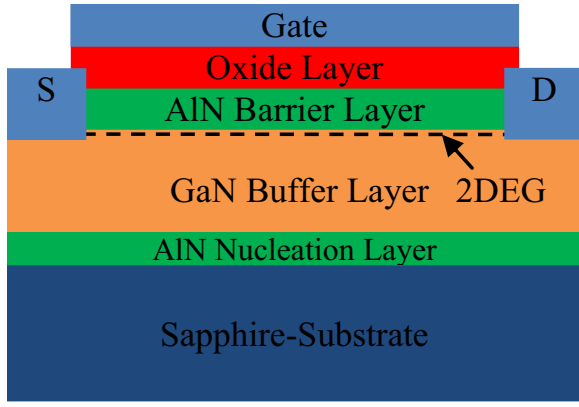


Fig. 1. AIN/GaN MOSHEMT structure.

Table 1

Parameters used for the Simulated Device.

Parameters	Dimensions
Gate Contact Length	6 $\mu\text{m}$
Source to Drain Distance (Effective Gate Length)	5.5 $\mu\text{m}$
AlN Barrier Thickness	2 nm
GaN Buffer Thickness	3 $\mu\text{m}$
AlN Nucleation layer thickness	5 nm
Sapphire Thickness	350 $\mu\text{m}$

### 2.1. MOSHEMT structure and dimensions

See Fig. 1 and Table 1.

### 2.2. Simulation setup

Drift-diffusion numerical model including carrier continuity equation and Poisson equation is used for the simulation of carrier transport in the channel at 300 K. Newton method is used for model calculations. Mesh value of the structure is carefully chosen by the Deck Build Editor for the accurate simulation of the important regions of the device as well as to accelerate the computational efficiency. A fine meshing is done in Deck build Editor around the gate foot and channel region such that there are 60 nodes within 2 nm thick barrier, i.e. distance between two nodes is around 0.6 Å. Simulations have been established using the physical models such as, FLDMOB which accounts for the electric field-dependent mobility, CONMOB, which accounts for the concentration dependent mobility, and CVT-a stand-alone model which incorporates all the effects required for simulating the carrier mobility, Fowler-Nordheim model for gate leakage current, SRH recombination and generation and Auger recombination model. Traps have been modeled by including Shockley-Read-Hall recombination terms in the continuity equation. The results are analysed in Tony plot and for output data extraction the Tony plot export feature is used.

During simulation the fixed oxide charges ( $V_{gs}$  independent) at oxide/semiconductor interfaces are considered as follows.  $N_{ox}$  is varied from  $-4.5 \times 10^{12} \text{ cm}^{-2}$  to  $-9.2 \times 10^{12} \text{ cm}^{-2}$  for  $\text{Al}_2\text{O}_3$ ,  $4.9 \times 10^{12} \text{ cm}^{-2}$  to  $9.8 \times 10^{12} \text{ cm}^{-2}$  for  $\text{HfO}_2$  [18] and  $-1.6 \times 10^{12} \text{ cm}^{-2}$  to  $-3.2 \times 10^{12} \text{ cm}^{-2}$  for  $\text{SiO}_2$  [22]. As the amount of oxide charges is dependent on deposition technique, for simulation purpose and to maintain uniformity among the oxides, in one case this value is referred from literature and in another case it is doubled. The strain dependent polarization charges of  $6.8 \times 10^{13} \text{ cm}^{-2}$  and  $5 \times 10^{13} \text{ cm}^{-2}$  are calculated and considered across oxide/AlN and AlN/GaN interface respectively [23].

### 2.3. Model Development for $V_T$

The schematic representation of conduction energy band across the metal/oxide/barrier/GaN heterojunction along with the charges formed at different interfaces is shown in Fig. 2.

The charge distribution across the junctions is also shown at the bottom of the figure, where  $\sigma_1$  represents the negative charge due to both spontaneous and piezoelectric polarization at oxide/AlN junction,  $n_s$  represents 2DEG density at AlN/GaN interface. The net charge at oxide/AlN junction consists of mainly four charge components such as positive surface donor trap charges ( $qN_d$ ), negative net polarization charges ( $\sigma_1$ ), negative oxide/AlN interfacial charge and fermi level dependent interface charge (due to interface density of state,  $d_{it}$ ) [24]. Among the above mentioned charges,  $N_d$  is always positive because of empty donor traps present at the surface of barrier and these trap states are the source of electrons inside the 2DEG which is still under debate [25].  $N_{ox}$  is negative for  $\text{SiO}_2$  and  $\text{Al}_2\text{O}_3$  but positive for  $\text{HfO}_2$  [18]. The polarity of  $N_{it}$  in fact depends on the position of fermi energy dependent, so it is determined by gate bias. As most of these charges concentrate at the oxide/barrier interface the effect of bulk oxide charges can be neglected. So the  $V_T$  expression for this structure can be written as,

$$V_T = \phi_b - \Delta E_c - \phi_F - \frac{t_{ox}}{\epsilon_{ox}}(\sigma_1 + \sigma_2) - \frac{t_b}{\epsilon_b} \sigma_2 - \frac{t_{ox}}{\epsilon_{ox}} q(N_d + N_{it}(E_F) + N_{ox}) \quad (1)$$

where  $\phi_b$  is the metal to semiconductor barrier height at zero gate bias,  $\Delta E_c$  is the conduction band offset at oxide/AlN interface,  $\phi_F$  is the GaN bulk fermi potential,  $\sigma_2$  is the polarization charge (equivalent of spontaneous and piezoelectric polarization charge of AlN and spontaneous polarization charge of GaN) at AlN/GaN interface,  $t$  is the thickness,  $\epsilon$  is the permittivity, and the subscripts  $ox$  and  $b$  refer to oxide and barrier respectively,  $N_d$  is the surface donor trap charge,  $N_{it}(E_F)$  is the oxide/AlN interface trap charge which is a function of  $E_F$  as  $N_{it}(E_F) = D_{it} \cdot (E_F - E_{NL})$ , where  $E_{NL}$  is the charge neutrality level present at the center of oxide band gap as shown in Fig. 2. It can be assumed as constant for the same range of  $V_{gs}$  when  $V_T$  is to be studied with respect to variation in  $N_{ox}$ . As

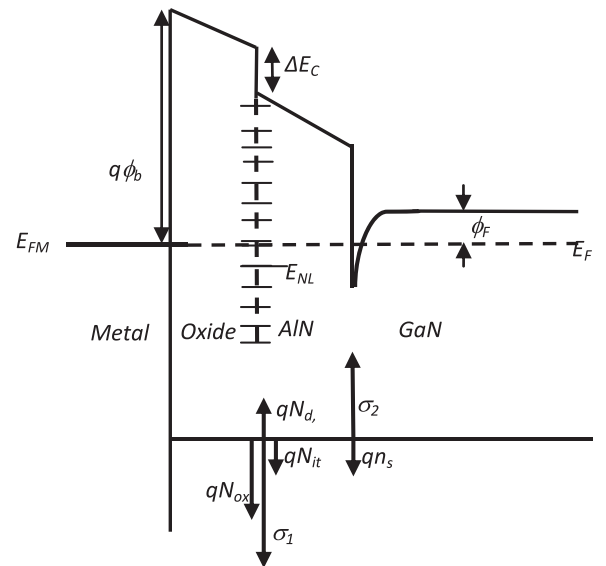


Fig. 2. Conduction band profile along with various charge distribution (bottom) in AIN/GaN MOSHEMT. Upward arrow shows positive charges and downward arrow shows negative charges.

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