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# BaTiO<sub>3</sub> as charge-trapping layer for nonvolatile memory applications

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#### 1. Introduction

Conventional floating-gate nonvolatile memories are approaching their scaling limits mainly due to their difficulties in maintaining high gate coupling ratio and suppressing cross talk between neighboring cells. Metal-Oxide-Nitride-Oxide-Silicon (MONOS) type nonvolatile memories have been considered as a promising candidate to replace the floating-gate counterpart because of their discrete charge-storage and coupling-free properties. Si<sub>3</sub>N<sub>4</sub> ( $k \sim 7$ ) was the first dielectric used as charge-trapping layer (CTL). In order to improve the data-storage capability by continual down scaling of the cell size, high-k dielectrics have been widely investigated to replace Si<sub>3</sub>N<sub>4</sub> as CTL [1-5]. Fig. 1 and Table 1 summarize the band-gap  $(E_g)$  and the conduction-band offset with respect to  $SiO_2$  ( $\Delta E_c$ ) of some typical high-k dielectrics used as CTL [6]. For a MONOS-type memory device, it tends to have thinner tunneling oxide for higher program/erase (P/E) speeds and lower operating voltages, whereas a thinner tunneling oxide may deteriorate its retention property. Therefore, there is a trade-off between high P/E speeds and good data retention. Among various high-k dielectrics, BaTiO<sub>3</sub> exhibits some distinguished properties including its large barrier height with respect to SiO<sub>2</sub> (conduction-band offset  $\Delta E_c \sim 3.6 \text{ eV}$  [6] and high dielectric constant ( $k \sim 100$  for the perovskite-type structure) [7], both of which are desirable for the charge-trapping layer of nonvolatile memories to improve

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## ABSTRACT

The charge-trapping (CT) properties of BaTiO<sub>3</sub> are investigated by using an Al/Al<sub>2</sub>O<sub>3</sub>/BaTiO<sub>3</sub>/SiO<sub>2</sub>/Si structure. The memory device with BaTiO<sub>3</sub> as CT layer shows promising performance in terms of large memory window (8.6 V by ±12 V for 1 s), high program speed with low gate voltage (a  $V_{FB}$  shift of 2.9 V at +6 V, 100 µs), negligible  $V_{FB}$  shift after 10<sup>5</sup>-cycle program/erase stressing, and good data retention property (charge loss of 7.9% after 10<sup>4</sup>-s 125 °C baking time), mainly due to the high charge-trapping efficiency of the BaTiO<sub>3</sub> film, as well as the large barrier height between the BaTiO<sub>3</sub> charge-trapping layer and the SiO<sub>2</sub> tunneling layer.

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the P/E speeds and retention properties. However, there have been few reports focusing on BaTiO<sub>3</sub> as CTL for charge-trapping-type nonvolatile memory applications [8]. Therefore, we concentrate on the charge-trapping characteristics of BaTiO<sub>3</sub> based on MONOS capacitors in this work. Experimental results demonstrated that this memory device with BaTiO<sub>3</sub> showed large memory window, high P/E speeds with low operating voltage, as well as good reliability.

### 2. Experiment

MONOS capacitors with Al/Al<sub>2</sub>O<sub>3</sub>/BaTiO<sub>3</sub>/SiO<sub>2</sub>/Si were fabricated on p-type silicon wafers. After a standard RCA cleaning, 2-nm SiO<sub>2</sub> was grown on the wafers by thermal dry oxidation. Then 10-nm BaTiO<sub>3</sub> was deposited on the SiO<sub>2</sub> by sputtering using a BaTiO<sub>3</sub> target in an Ar/O<sub>2</sub> (24 sccm/6 sccm) mixed ambient at a pressure of 3.0 mTorr. Then 15-nm Al<sub>2</sub>O<sub>3</sub> as blocking layer was deposited by means of atomic layer deposition using trimethylaluminum (Al(CH<sub>3</sub>)<sub>3</sub>) and H<sub>2</sub>O as precursors at a substrate temperature of 300 °C. Following that, all the samples went through a post-deposition annealing (PDA) in N<sub>2</sub> ambient at 900 °C for 30 s. Then, Al was evaporated and patterned as gate electrode with a diameter of 100 µm, followed by forming-gas annealing at 300 °C for 20 min. The physical properties of the dielectric films were analyzed by transmission electron microscopy (TEM), X-ray diffraction (XRD) and X-ray photoelectron spectroscopy (XPS). The electrical characteristics of the MONOS capacitors were measured by HP4284A LCR meter and HP4156A semiconductor parameter





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analyzer. Their flat-band voltage ( $V_{\text{FB}}$ ) was extracted from their measured C-V (capacitance–voltage) curve at the capacitance equal to the calculated flat-band capacitance [9].

#### 3. Results and discussion

The inset of Fig. 2a shows the cross-sectional TEM image of the MONOS capacitor with BaTiO<sub>3</sub> as CTL, where the thickness of Al<sub>2</sub>O<sub>3</sub>/BaTiO<sub>3</sub>/SiO<sub>2</sub> is determined to be 15.6 nm/10.6 nm/2.0 nm respectively. It is also observed that the BaTiO<sub>3</sub> film displays an amorphous phase, which can be further confirmed by the XRD spectrum with no peaks shown in Fig. 2a. It is worth mentioning that the charge-trapping film in an amorphous state is favorable for the reliability of the memory device because charge loss via grain boundaries and defects at the CTL/SiO<sub>2</sub> interface induced by lattice mismatch can be avoided [3]. These defects along the grain boundaries or at the CTL/SiO<sub>2</sub> interface can enhance charge leakage and degrade the dielectric quality, and so are considered as reliability degraders, instead of effective traps for charge storage. Fig. 2b exhibits the Ti  $2p_{3/2}$  spectrum as well as the curve-fitting lines. Each curve-fitting line is assumed to follow the general shape of Lorentzian–Gaussian function. The Ti  $2p_{3/2}$  spectrum can be decomposed into two components, corresponding to Ti-O bonding in the BaTiO<sub>3</sub> film ( $\sim$ 458.4 eV for Ti 2 $p_{3/2}$ ) [10] and Ti silicate at the BaTiO<sub>3</sub>/SiO<sub>2</sub> interface ( $\sim$ 459.0 eV for Ti 2p<sub>3/2</sub>) [11], respectively, where the Ti-silicate component displays a weak peak, indicating only a small fraction of BaTiO<sub>3</sub> transformed into silicate. This is consistent with the observation of the TEM image with no obvious interlayer at the CTL/SiO<sub>2</sub> interface. This non-stoichiometric Ti-silicate interlayer normally has a much smaller barrier height with respect to BaTiO<sub>3</sub> ( $E_g \sim 3.3 \text{ eV}$ ) as well as more defects compared with that of thermal SiO<sub>2</sub> tunneling layer ( $E_g \sim 9.0 \text{ eV}$ ) [6,10]. Consequently, an abrupt interface is desirable for good data retention properties because the charge-loss process related to the interlayer (e.g. defect-assisted tunneling) can be suppressed [12].

Fig. 3a depicts the 1-MHz *C*–*V* hysteresis characteristics of the memory device with BaTiO<sub>3</sub> as CTL. Sweep starts from inversion region to accumulation region and back to inversion region again, corresponding to electron trapping and de-trapping respectively. As the sweeping voltage increases from ±8 V to ±12 V, the hysteresis window, defined as the difference of  $V_{FB}$  corresponding to the backward and forward *C*–*V* curves, increases from 3.0 V to 7.6 V. The large window indicates a high trap density in the BaTiO<sub>3</sub> film. Moreover, the  $V_{FB}$  of the *C*–*V* loop is more positive than the neutral  $V_{FB}$  ( $V_{FB}$  of the fresh device). Therefore, electrons trapped in the CTL during the backward sweeping are not completely removed after



**Fig. 1.** Energy band-gap ( $E_g$ ) and conduction-band offset with respect to SiO<sub>2</sub> ( $\Delta E_c$ ) of typical high-*k* dielectrics used as charge-trapping layer.

Table 1
Comparison of $E_g$ and $\Delta E_c$ for typical high- <i>k</i> dielectrics used as charge-trapping layer.

310	2 BaliO <sub>3</sub>	1a <sub>2</sub> O <sub>5</sub>	$ZrO_2$	HfO <sub>2</sub>	Y <sub>2</sub> O <sub>3</sub>	S13N4	
$ \begin{array}{cc} E_g \left( eV \right) & 9.0 \\ \Delta E_c \left( eV \right) & 0 \end{array} $	3.3 3.6	5.8 3.2	5.8 2.1	6.0 2.0	6.0 1.2	5.3 1.1	

the forward sweeping, and thus a larger forward sweeping stress  $(-16 \text{ V} \rightarrow +12 \text{ V})$  is required to bring the C-V curve back to the neutral position as shown in Fig. 3b. These phenomena should be associated with deep traps in the BaTiO<sub>3</sub> film [13] because electrons located in deep traps are more difficult to escape than those in shallow traps. As further increasing the forward sweeping stress  $(-18 \text{ V} \rightarrow +12 \text{ V})$ , the V<sub>FB</sub> slightly shifts towards the negative direction with respect to the neutral state, indicating only a small amount of hole traps in the BaTiO<sub>3</sub> film. Different from SiO<sub>2</sub> with covalent bonds, the ionic bonding in BaTiO<sub>3</sub> suggests its poor ability to remove a defect once created. It has been reported that oxygen can escape from BaTiO<sub>3</sub> during high-temperature annealing in a N<sub>2</sub> protective ambient, thus leading to oxygen vacancies [14]. Moreover, the Ti<sup>4+</sup> sites in titanate can also act as acceptor levels for electron trapping [15]. Both of these reasons can contribute to high electron-trap density in the BaTiO<sub>3</sub> film. Considering the Ti 2p XPS spectrum with the fully oxidized Ti<sup>4+</sup> state shown in Fig. 2b, it is believed that few oxygen vacancies are present in the BaTiO<sub>3</sub> film [16].



Fig. 2. (a) XRD spectrum of the BaTiO<sub>3</sub>/SiO<sub>2</sub> (10 nm/2 nm) film on Si substrate. The XRD sample received the same thermal cycles as the MONOS capacitor. The inset shows the cross-sectional TEM image of the MONOS capacitor with Al/Al<sub>2</sub>O<sub>3</sub>/BaTiO<sub>3</sub>/SiO<sub>2</sub>/Si. (b) Ti 2p<sub>3/2</sub> XPS spectrum of the BaTiO<sub>3</sub>/SiO<sub>2</sub> stack on Si substrate.

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