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Theoretical studies on electroluminescent mechanism of a series of thermally activated delayed fluorescence emitters possessing asymmetric-triazine-cored triads



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ABSTRACT

A series of thermally activated delayed fluorescence (TADF) emitters using asymmetric-triazine-cored triad as the electron-accepting unit have been investigated theoretically. Based on two experimentally reported TADF molecules (**TPXZ-as-TAZ** and **oDPXZ-as-TAZ**), two new molecules (**mDPXZ-as-TAZ** and **pDPXZ-as-TAZ**) have been designed to explore the isomeric effect on their TADF properties. The present results reveal that the absorption and emission spectra calculated by the time-dependent density functional theory (TD-DFT) method at the M06-2X level are match well the available experimental findings, and **mDPXZ-as-TAZ** and **pDPXZ-as-TAZ** are found to exhibit the same yellow emission as their analogue **oDPXZ-as-TAZ**. In addition, the rates of reverse intersystem crossing of **mDPXZ-as-TAZ** and **pDPXZ-as-TAZ** estimated by the semiclassical Marcus theory are 2.51 $\times 10^6$ and 4.57 $\times 10^6$ s⁻¹, respectively, one order of magnitude larger than that of **oDPXZ-as-TAZ** can be also considered as potential yellow-light TADF emitters.

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1. Introduction

In recent years, the thermally activated delayed fluorescence (TADF) emitters have attracted a great deal of interest in that these TADF emitters are capable of harvesting both singlet and triplet excitons in organic light-emitting diodes (OLEDs) devices to achieve the internal quantum efficiency (IQE) of 100%, which is desirable to realize high-performance OLEDs for flexible flat-panel displays and solid-state lighting [1–46].

It is known that TADF molecules usually comprise electron donor (D) and acceptor (A) units, and their excited states are observed to exhibit a charge transfer (CT) characteristic from D to A framework. Owing to the strongly electron-deficient character, a symmetric 1,3,5-triazine (*s*-TAZ) group has been extensively adopted as the A moiety of a D-A or D- π -A system that shows significant TADF property [47–54]. A representative example is the report of 9-(4-(4,6-diphenyl-1,3,5-triazin-2-yl)-9'-phenyl-9H,9'H-3,3'-bicarbazole exhibiting a sky-blue emission in OLEDs device with a maximum external quantum efficiency (EQE) of 21.7% [49]. As an isomer of *s*-TAZ, an asymmetric 1,2,4-triazine (*as*-TAZ) unit has recently received attention from researchers. One main motivation is to make it clear whether or not *as*-TAZ can be used as an efficient A unit in the TADF emitters. Based on the speculation, Yang

* Corresponding author. *E-mail address:* xgguo@henu.edu.cn (X. Guo). and co-worker designed two asymmetric-triazine-cored triads, that is, 10,10''.(1,2,4-triazine-3,5,6-triyl)tris(benzene-4,1-diyl))tris(10*H*-phenoxazine) (**TPXZ-as-TAZ**) and 10,10'-((3-phenyl-1,2,4-triazine-5,6-diyl)bis(4,1-phenylene))bis(10*H*-phenoxazine) (**oDPXZ-as-TAZ**) (see Fig. 1), in which the *as*-TAZ core is used as the A unit, phenyl groups as the π -bridges, and phenoxazine (PXZ) groups as the D moieties [55]. It was found that both compounds exhibits remarkable TADF character with short delayed fluorescence lifetimes (1.10 µs for **TPXZ-as-TAZ** and 0.98 µs for **oDPXZ-as-TAZ**) in doped CBP films. More importantly, OLEDs devices incorporating the two compounds as the emitters deliver yellow emission with EQEs of 11.5% and 8.4%, respectively.

To our best knowledge, theoretical investigations on **TPXZ-as-TAZ** and **oDPXZ-as-TAZ** only focus on the structural optimizations for ground state and spectral simulations by employing the density functional theory (DFT) and time-dependent density functional theory (TD-DFT) calculations, respectively [55]. It has been widely accepted that the intersystem crossing (ISC) from the lowest singlet excited S₁ state to the lowest triplet excited T₁ state and reverse ISC (RISC) from T₁ to S₁ are two very important processes for the TADF molecules. Therefore, a comprehensively theoretical studies on ISC and RISC are highly required. In the present calculations, herein, we performed a systematic computational investigation on four D- π -A-type molecules shown in Fig. 1. Two of them are experimentally synthesized molecules (**TPXZ-as-TAZ** and **pDPXZ-as-TAZ**) are newly designed from a



Fig. 1. Chemical structures of four donor-π-acceptor-type molecules investigated in this work. Θ_1 , Θ_2 and Θ_3 denote dihedral angles, and B_1 , B_2 and B_3 represents bond length between the acceptor and phenyl bridges.

theoretical point of view, which are isomers of **oDPXZ-as-TAZ**. A main objective is to elucidate the isomeric effect on their TADF properties. It is expected that the present computational results can serve as a theoretical guide for developing the new efficient TADF-based OLED emitters.

2. Computational Details

Geometry optimizations of molecules as shown Fig. 1 were performed to determine minimum energy structures in the gas phase for the ground S₀ state by DFT, and for the S₁ and T₁ excited states by TD-DFT using the hybrid meta-generalized gradient approximation (meta-GGA) exchange correlation functional M06-2X [56]. This functional provides a reasonable description of the ground-state and excited-state geometries [33]. Here, Pople's 6-31G(d) basis set was used in both cases. In addition, their absorption and emission spectra were also estimated by employing the TD-M06-2X/6-31G(d) calculations. Considering that the absorption spectra of **TPXZ-as-TAZ** and **oDPXZ-as-TAZ** were measured in dichloromethane solution, the polarizable continuum model (PCM) using CH₂Cl₂ (ε = 8.93) as solvent was adopted [57,58].

Based upon the optimized structures, the adiabatic excitation energy method was applied to evaluate the singlet-triplet energy differences between the S₁ and T₁ states. Fig. 2 plots the potential energy surfaces (PESs) of the S₀, S₁ and T₁ electronic states. $E_{00}(S_1)$ and $E_{00}(T_1)$ represent the adiabatic excitation energies between the energy minima of S₀ and S₁/T₁ along PESs.

 $E_{00}(S_1) = E_{S_1}(S_1 - \text{geometry}) - E_{S_0}(S_0 - \text{geometry})$ (1)

$$E_{00}(T_1) = E_{T_1}(T_1 - \text{geometry}) - E_{S_0}(S_0 - \text{geometry})$$
(2)

Therefore, the adiabatic excitation energy difference (ΔE_{ST}) can be defined as:

$$\Delta E_{ST} = E_{00}(S_1) - E_{00}(T_1) \tag{3}$$

The (R)ISC rates can be calculated by using semiclassical Marcus theory expression [59–61]:

$$k = \frac{V^2}{\hbar} \sqrt{\frac{\pi}{k_{\rm B} T \lambda}} \exp\left[-\frac{(\Delta E_{\rm ST} + \lambda)^2}{4k_{\rm B} T \lambda}\right] \tag{4}$$

where $k_{\rm B}$ is Boltzmann constant and *T* is temperature, which is set to 300 K. *V* denotes the spin-orbital coupling (SOC) matrix element between the S₁ and T₁ states, which could be treated as a perturbation on the basis of the scalar relativistic (SR) orbitals after the SCF and TD-DFT calculations [62], together with an Slater-type all-electron DZP basis set [63]. Additionally, λ is the reorganization energy induced by the electron or the energy transfer. For the S₁ and T₁ states, λ could be formulated as (see Fig. 1):

$$\lambda_{\rm S} = E_{\rm S_1}({\rm T}_1 - \text{geometry}) - E_{\rm S_1}({\rm S}_1 - \text{geometry}) \tag{5}$$

$$\lambda_{\rm T} = E_{\rm T_1}(S_1 - \text{geometry}) - E_{\rm T_1}(T_1 - \text{geometry}) \tag{6}$$

where $E_{S_1}(T_1/S_1 - \text{geometry})$ is the S_1 energy at the T_1/S_1 minimum geometry, and $E_{T_1}(T_1/S_1 - \text{geometry})$ is the T_1 energy at the T_1/S_1 minimum geometry. It should be mentioned that in calculation of the k_{ISC} rate, the λ_T value is used and $\Delta E_{ST} = E_{00}(T_1) - E_{00}(S_1)$.

All the DFT and TD-DFT computations were carried out with the Gaussian 09 package [64]. The SOC calculations were performed using the ADF2016 package [65].



Fig. 2. Potential energy surfaces of the ground S₀ state, the lowest singlet excited S₁ state, and the lowest triplet excited T₁ state.

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