

# Structures and Potential Energy Functions of PdN and PdN<sub>2</sub> Molecules

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**Abstract:** Based on the relativistic effective core potential (RECP/SDD) for Pd atom and AUG-cc-pVTZ basis function for N atom, the structures of PdN and PdN<sub>2</sub> have been optimized using B3LYP method. The results show that the ground state of PdN is  $^4\Sigma^-$  and the ground state of PdN<sub>2</sub> is  $C_{\infty V}$  symmetry and  $^1\Sigma^+$  state. The Murrell-Sorbie potential energy function of PdN molecule has been fitted through the least square fitting, and the potential energy function of PdN<sub>2</sub> is based on many-body expansion theory. The potential energy curves describe the structure character of PdN and PdN<sub>2</sub> ground state molecules correctly, and detail the inner transfer process of Pd. There is a saddle point in  $C_{2V}$  structure ( $R_{NN}=0.11700$  nm,  $R_{PdN}=0.22088$  nm). The energy barrier of inner transfer is 0.5197 eV, which is close to the calculated value 0.4560 eV.

**Key Words:** PdN; PdN<sub>2</sub>; Density functional method; Relativistic effective core potential (RECP); Many-body expansion potential energy function

Using the palladium alloys for hydrogen diffusion membrane method has its unique superiority in pick-up, separation and purification of hydrogen isotopes. But the disadvantage of using Pd alloys membrane is that impurity of gases can poison it through adsorption, dissociation and occupancy course on the membrane surface thus changing its structure. There was a preliminary study of impurity gases poisoning palladium alloys membrane<sup>[1,2]</sup>. In this article, the structural characters of possible resultants such as PdN and PdN<sub>2</sub> of Pd and N<sub>2</sub> are calculated, and the potential energy functions are given. These are very important to understand the molecule reaction dynamics and micro-mechanism occurring during membrane poison of palladium alloys.

## 1 Calculation method

Pd is a heavy transition element. Its relativistic effect is very evident because of *d* orbit interaction. Whereas the character of atom is decided mostly by valence electron, Pd can be dealt using RECP (Relativistic Effective Core Potential). The results have been given using RECP<sup>[3,4]</sup>.

The structures of PdN and PdN<sub>2</sub> are optimized using B3LYP (the hybrid density functional theory (DFT) of

Becke's three parameter exchange function and Lee-Yang-Parr correlation function)<sup>[5]</sup> method, RECP/SDD basis function for Pd and AUG-cc-pVTZ basis function for N. The Murrell-Sorbie potential energy function of PdN molecule has been fitted through the least square fitting, and the potential energy function of PdN<sub>2</sub> is given by many-body expansion theory.

## 2 Result and discussion

### 2.1 Structures of PdN and PdN<sub>2</sub> molecules

The *ab initio* calculation shows that ground state of PdN is  $^4\Sigma^-$ , equilibrium nuclear distance  $R_e=0.18158$  nm, dissociation energy  $D_e=2.1437$  eV. The ground state of PdN<sub>2</sub> is  $C_{\infty V}$  symmetry and  $^1\Sigma^+$  state, equilibrium nuclear distance  $R_{NN}=0.11026$  nm,  $R_{PdN}=0.19537$  nm, dissociation energy  $D_e=10.7047$  eV. Then harmonic vibration frequencies of the ground state PdN<sub>2</sub>( $C_{\infty V}$ ) are given: symmetry stretched vibration frequency  $\nu_1=356.9753$  cm<sup>-1</sup>, bend vibration frequency  $\nu_2=180.9400$  cm<sup>-1</sup> and antisymmetry stretched vibration frequency  $\nu_3=2318.1252$  cm<sup>-1</sup>.

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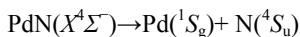
## 2.2 Dissociation limit of PdN and PdN<sub>2</sub> molecules

The ground state of PdN is  $X^4\Sigma^-$ , and its orbit arrangement is shown below:

$\alpha$  electron:  $\sigma\sigma\sigma\pi\pi\sigma\sigma\pi\pi\delta\delta\pi\pi\sigma$

$\beta$  electron:  $\sigma\sigma\sigma\pi\pi\sigma\sigma\pi\pi\delta\delta$

According to the theory of atomic and molecular reaction statics<sup>[6]</sup>, dissociation limit of the ground state PdN( $X^4\Sigma^-$ ) molecule is shown as follows:

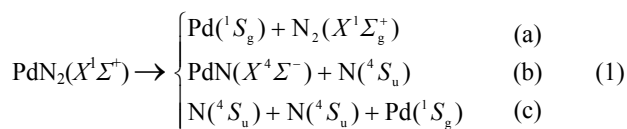


The ground state of PdN<sub>2</sub> is  $X^1\Sigma^+$ ,  $C_{\infty V}$  group and its orbit arrangement is shown below:

$\alpha$  electron:  $\sigma\sigma\sigma\sigma\pi\pi\sigma\sigma\pi\pi\sigma\pi\pi\delta\delta\sigma$

$\beta$  electron:  $\sigma\sigma\sigma\sigma\pi\pi\sigma\sigma\pi\pi\sigma\pi\pi\delta\delta\sigma$

Dissociation channels of the ground state PdN<sub>2</sub>( $X^1\Sigma^+$ ) molecule are shown as follows:



As for channel (a),  $^1S_g$  can be resolved into  $^1A_1$  of  $C_{\infty V}$  group, and  $^1\Sigma_g^+$  into  $^1A_1$  of  $C_{\infty V}$  group. According to the direct product of the  $C_{\infty V}$  group representation:  $^1A_1 \otimes ^1A_1 = ^1A_1$ . This is the ground state of PdN<sub>2</sub> molecule.

For channel (b),  $^4\Sigma^-$  can be resolved into  $^4A_1$  of  $C_{\infty V}$  group, and  $^4S_u$  into  $^4A_1$  of  $C_{\infty V}$  group. According to the direct product of the  $C_{\infty V}$  group representation:  $^4A_1 \otimes ^4A_1 = (^1A_1 \otimes ^3A_1 \otimes ^5A_1 \otimes ^7A_1)$ . It also includes component of  $^1A_1$ .

For channel (c), according to dissociation channel of the ground state PdN( $X^4\Sigma^-$ ) and the microscopic process transitive principle, it is possible.

## 2.3 The potential energy function and spectral data of PdN molecule

After the structure of PdN has been optimized using B3LYP method, the single point energy (SP) of the ground state PdN( $X^4\Sigma^-$ ) is calculated, the bond length is maintained from 0.10 to 0.58 nm and step length is 0.01 nm. Then the form of Murrell-Sorbie (M-S) potential energy function<sup>[7]</sup> is fitted through the least square fitting.

$$V = -D_e(1 + a_1\rho + a_2\rho^2 + a_3\rho^3)\exp(-a_1\rho) \quad (2)$$

Where,  $\rho = R - R_e$ ,  $R$  and  $R_e$  are the nuclear distance and equilibrium nuclear distance, respectively.  $D_e$ ,  $a_1$ ,  $a_2$  and  $a_3$  are fitting parameters. The fitting results are shown in Table 1.

As can be seen in Fig.1, the fitting curve (solid line)

Table 1 M-S potential energy functions for PdN and N<sub>2</sub> molecules

State	$D_e/\text{eV}$	$R_e/\text{nm}$	$a_1/\text{nm}^{-1}$	$a_2/\text{nm}^{-2}$	$a_3/\text{nm}^{-3}$
PdN( $X^4\Sigma^-$ )	2.1437	0.18158	35.413	186.38	705.01
N <sub>2</sub> ( $X^1\Sigma^+$ )	9.905	0.10977	53.96	732.8	4985

shows good agreement with *ab initio* calculation values of PdN( $X^4\Sigma^-$ ). By using the Murrell-Sorbie potential energy function (2), harmonic force constants and spectral data of PdN( $X^4\Sigma^-$ ) are calculated according to the formulas<sup>[7]</sup> about parameters  $D_e$ ,  $a_1$ ,  $a_2$  and  $a_3$  with harmonic force constants and spectral data. The results are shown in Table 2.

## 2.4 The many-body expansion potential energy function of PdN<sub>2</sub> molecule

The reference structure of  $C_{2V}$  is chosen in this article (PdN<sub>2</sub>:  $R_1^0 = R_{NN}$ ,  $R_2^0 = R_3^0 = (R_{NN} + R_{PdN} + R_{PdN})/2$  are reference coordinates) to study the potential energy function of PdN<sub>2</sub> molecule. The internal coordinates are written as  $\rho_i = R_i - R_i^0$  ( $i=1,2,3$ ), then the corresponding symmetric internal coordinates are given as follows:

$$\begin{bmatrix} S_1 \\ S_2 \\ S_3 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1/\sqrt{2} & 1/\sqrt{2} \\ 0 & 1/\sqrt{2} & -1/\sqrt{2} \end{bmatrix} \begin{bmatrix} \rho_1 \\ \rho_2 \\ \rho_3 \end{bmatrix}$$

The many-body expansion potential energy function of PdN<sub>2</sub> molecule is as follows:

$$V(R_1, R_2, R_3) = V_{NN}^{(2)}(R_1) + V_{PdN}^{(2)}(R_2) + V_{PdN}^{(2)}(R_3) + V_{NNPd}^{(3)}(R_1, R_2, R_3) \quad (3)$$

Where,  $V_{NN}^{(2)}(R_1)$ ,  $V_{PdN}^{(2)}(R_2)$  and  $V_{PdN}^{(2)}(R_3)$  are two-body terms, and use form of Murrell-Sorbie potential energy function (2), which parameters are shown in Table1, the pa-

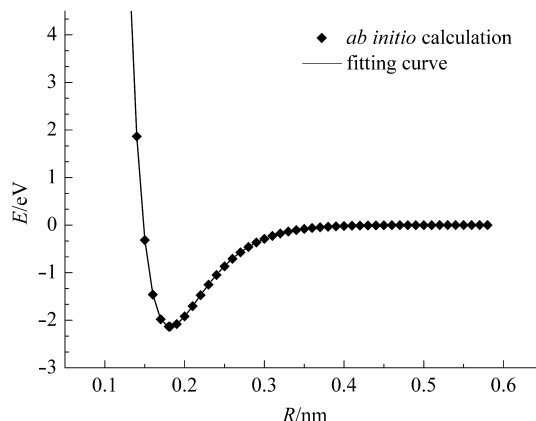


Fig.1 The potential curve of PdN( $X^4\Sigma^-$ )

Table 2 Harmonic force constants and spectral data of PdN( $X^4\Sigma^-$ )

State	$10^{-2}f_2/\text{aJ}\cdot\text{nm}^{-2}$	$10^{-3}f_3/\text{aJ}\cdot\text{nm}^{-3}$	$10^{-4}f_4/\text{aJ}\cdot\text{nm}^{-4}$	$\omega_e/\text{cm}^{-1}$	$\omega_e x_e/\text{cm}^{-1}$	$B_e/\text{cm}^{-1}$	$\alpha_e/\text{cm}^{-1}$
$X^4\Sigma^-$	3.0270	-18.3581	86.2961	$6.4427 \times 10^3$	558.2997	0.4131	$5.6743 \times 10^{-3}$

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