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## Effect of material hysteresis in magnetic refrigeration cycles

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#### Abstract

In this paper we make use of the concepts of out-of-equilibrium thermodynamics applied to systems with hysteresis in order to model the magnetic material employed as a working substance in magnetic refrigeration cycles. The approach developed leads to a detailed description of heat fluxes as well as the entropy production connected with hysteresis. As an example we discuss Carnot and AMR refrigeration cycles where hysteresis effects are included.

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Keywords: Magnetic refrigeration; Preisach model; Thermodynamic cycle

# Effet de l'hystérésis du matériau dans les cycles frigorifiques magnétiques

Mots clés : Froid magnétique ; Modèle de Preisach ; Cycle thermodynamique

#### 1. Introduction

The magnetocaloric effect has been studied for a long time for its potential use in refrigeration cycles. Recent progresses in material science, notably with the development of new materials exhibiting magnetocaloric effect (see the reviews of [1,2]), have generated a renewed interest in the development of magnetic refrigeration systems working around room temperature. Materials of interest include rare-earth-based ferromagnets, like Gd, Y–Fe and Gd–Fe,

this new family of materials has been accompanied by a growing interest in their physical properties including, temperature changes and entropy changes achievable by magnetic field. In this context, the fact that the working material displays hysteresis [3] calls for a better physical understanding of the role of

irreversibility in the magnetization process.

Within the limits of equilibrium thermodynamics all the machinery of classical Gibbs thermostatics may be successfully used to describe the materials and to derive the performance of thermodynamic cycles. For example, the temperature change or entropy change due to a magnetic field

with a Curie point around room temperature, and alloys like Gd-Si-Ge [1], which display the so-called giant mag-

netocaloric effect connected to the presence of a first-order

The technological outlook promised by the introduction of

phase transformation induced by magnetic field.

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S	specific entropy	$z_{\rm c}$	normalized local coercive field of the Preisach
T	absolute temperature		unit
M	magnetization	$Q_{ m u}$	$h_{\rm u}$ scaling factor
$\mu_0$	permeability of vacuum	$Q_{ m c}$	$h_{\rm c}$ scaling factor
Н	magnetic field	$\tilde{p}(z_{ m c},z_{ m u})$	normalized Preisach distribution
$\delta s$	variation of the specific entropy	$ ilde{b}(z_{ m c})$	normalized state line
$\delta_{\rm e} s$	variation of the exchanged specific entropy	$ ilde{H}$	normalized magnetic field
$\delta_{i}s$	variation of the specific entropy production	$g_{\mathbf{M}}$	specific Gibbs free energy of the magneti
$M_{\rm s}$	saturation magnetization		subsystem
$\Delta m$	magnetic moment of the Preisach unit	$f_{M}$	specific free energy of the magnetic
$h_{ m u}$	local interaction field of the Preisach unit		subsystem
$h_{\rm c}$	local coercive field of the Preisach unit	$s_{\mathbf{M}}$	specific entropy of the magnetic subsystem
$p(h_{\rm c}, h_{\rm u})$	Preisach distribution	$s_{ m B}$	specific entropy of the non-magnetic subsyster
$b(h_{\rm c})$	state line	$s_{\rm R}$	specific entropy of the regenerative fluid
$z_{\mathbf{u}}$	normalized local interaction field of the	$v_{\rm B}$	volume of the magnetic material
	Preisach unit	$v_{\rm R}$	volume of the regenerative fluid

H may be derived from the graphical representation of the specific entropy s(T, H) constitutive relation in the s-T plane [4]. The performance of different kinds of cycles (Carnot, Ericsson, Bryton, AMR) can be directly derived from these diagrams.

However, if the material displays hysteresis, the equilibrium s-T diagram is no longer sufficient to describe the material behavior. Internal entropy production takes place during the magnetization process and a specific out-of-equilibrium approach to the material behavior is needed.

The aim of our study is to develop a modelling approach permitting one to represent the out-of-equilibrium properties of the magnetic material. The constitutive relation is replaced by the entropy balance equation  $\delta s = \delta_i s + \delta_e s$ , in which the measurable quantity, i.e. the heat flux  $T\delta_e s$ , is proportional to the exchanged entropy  $\delta_e s$  and the internal entropy production term  $\delta_i s$  is present.

The model we describe in this paper is far from being a general theory of hysteretic systems. However, it has the merit of working out in detail a case for which both the system entropy s and the entropy production  $\delta_i s$  can be expressed in closed analytical forms. Furthermore, the equations of state describing the magnetization as a function of the field and temperature have been shown to be an excellent tool for the description of ferromagnetic materials behavior [5,6]. In addition, the formalism of our approach bears strong similarities with a well known approach to out-of-equilibrium systems which is known as thermodynamics of materials with memory [7].

Irreversible thermodynamics is the appropriate context to study dissipation phenomena associated with hysteresis. In recent papers (see Ref. [8] for a review) the thermodynamics of hysteresis was studied for a system made of the superposition of non-interacting two-level systems (TLS) using an internal variables thermodynamic approach. When the magnetic field *H* and the temperature *T* are externally controlled

quantities, explicit expressions for the magnetization M, the entropy s and the entropy production  $\delta_i s$  were derived. One of the merits of this approach is that the isothermal constitutive relation M(H;T) is equivalent to the very well known Preisach model of hysteresis [8] whose applications to magnetic materials have been extensively studied.

A common way to describe systems characterized by long-living non-equilibrium states consists in the introduction of a proper set of additional variables, vanishing when the system reaches equilibrium, which are assumed to still enable the use of state functions, like Gibbs free energy, when the system is not at equilibrium. One of the main properties of these non-equilibrium variables is that they appear in the constitutive relations, but not in the work expression of the first law. For this reason they are called internal variables [7]. In our case internal variables account for dissipative processes taking place along hysteresis curves. This approach is extensively discussed in Ref. [8].

In Section 2 we make use of the internal variable approach in order to formulate the thermodynamics of hysteresis for a system consisting of the superposition of two-level subsystems (Preisach model) [8]. The approach gives the explicit expressions for the magnetization, the entropy and the entropy production for a system with hysteresis [8]. The resulting model is an appropriate description of the magnetic part (M) of a ferromagnetic body.

In Section 3 we discuss how this approach can be employed to model the thermodynamic behavior of a ferromagnetic body in order to predict the exchanged heat along a thermodynamic transformation [9,10].

In Section 4 we employ the same model to compute thermodynamic transformations in refrigeration cycles. Due to the conceptual difficulties associated with constitutive laws with hysteresis, the literature dealing with the design and optimization of thermodynamic cycles for magnetic

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