



## Regular article

# Mechanical softening of thermoelectric semiconductor $Mg_2Si$ from nanotwinning

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## ABSTRACT

Nanotwinning exhibits strengthening effects in many metals, semiconductors, and ceramics. However, we show from *ab-initio* calculations that nanotwins significantly decrease the strength of thermoelectric semiconductor  $Mg_2Si$ . The theoretical shear strength of nanotwinned  $Mg_2Si$  is found to be 0.93 GPa, much lower than that (6.88 GPa) of flawless  $Mg_2Si$ . Stretching the Mg–Si bond under deformation leads to the structural softening and failure of flawless  $Mg_2Si$ . While in nanotwinned  $Mg_2Si$ , the Mg–Si bond at the twin boundary (TB) is expanded to accommodate the structural misfit, weakening the TB rigidity and leading to the low ideal shear strength.

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Solid-state thermoelectric (TE) technology, which can realize the direct conversion from heat to electrical energy, attracts renewable attention recently because it could play an essential role in relieving the global energy crisis [1,2].  $Mg_2Si$  based TE material is a promising high-performance candidate with the  $zT$  value of 1.5 in the mediate temperature range [3–7]. Therefore,  $Mg_2Si$ , which consists of globally abundant and eco-friendly elements, is now being considered for application in the automobile waste heat recovery system [8]. However, under thermal cycling loads, thermo-mechanical stress easily causes the crack initiation of  $Mg_2Si$ , hence leading to the mechanical failure of  $Mg_2Si$  TE devices. Thus, mechanical properties of  $Mg_2Si$  are of vital significance for its engineering applications.

The roles of nanotwins on mechanical properties of metal, semiconductor, and ceramics have been widely examined. In particular, nanoscale twins were found very efficient to strengthening the materials. Lu et al. synthesized stable nanotwinned copper samples with a high twin density, which exhibit a ten times higher tensile strength compared to traditional coarse-grained copper [9]. An. et al. showed from

quantum mechanics (QM) simulations that the ideal shear strength of boron carbide can be enhanced by 11% when inserting nanotwins [10]. Huang. et al. directly synthesized highly twinned diamond samples with an average grain size of ~5 nm, which exhibit unprecedented high hardness (Vickers hardness of ~200 GPa) and high thermal stability [11]. Tian et al. reported that nanotwinned cubic boron nitrides possess an extremely high Vickers hardness (>100 GPa) and a large fracture toughness (>12 MPa m<sup>1/2</sup>) [12]. Our previous *ab-initio* study showed that through introducing nanoscale twin boundaries, the ideal shear strength of TE semiconductor  $Bi_2Te_3$  and  $InSb$  can be enhanced by 215% and 11%, respectively [13,14]. A recent experiment by Jang et al. showed the formation of nanotwins in TE semiconductor  $Mg_2Si$  [15], suggesting that they may play an important role in the mechanical properties. However, this remains unexplored.

To understand how nanotwins influence the mechanical properties of  $Mg_2Si$ , we utilized *ab-initio* calculations at the Perdew-Burke-Ernzerhof (PBE) functional level to examine the theoretical strength and failure mechanism under pure shear and biaxial shear deformations, respectively. We find that the theoretical shear strength of nanotwinned  $Mg_2Si$  is only 0.93 GPa which is much lower than that (6.88 GPa) of flawless  $Mg_2Si$ . This is in contrast with the above examples that the strength of materials is enhanced by nanotwins. This nanotwin-induced softening behavior in  $Mg_2Si$  arises from the expanded and weakened Mg–Si bond in the twin boundary (TB). We also find that

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compression plays an essential role in determining the theoretical strength of flawless  $\text{Mg}_2\text{Si}$  under biaxial shear load. But it has little influence on strength and deformation mode of nanotwinned  $\text{Mg}_2\text{Si}$ .

All *ab-initio* calculations were implemented in the VASP software [16,17], employing the PBE functional and the projector augmented wave (PAW) potentials to account for the core-valence interactions [18,19]. A kinetic energy cutoff of 500 eV was used for geometry optimization. The electronic self-consistent and force convergence criterion are  $1 \times 10^{-6}$  eV and  $1 \times 10^{-2}$  eV/Å, respectively. The Monkhorst-Pack centered  $k$ -points sampling with a fine resolution of  $2\pi \times 1/40 \text{ \AA}^{-1}$  was performed for all calculations. The pure shear and biaxial shear simulation setups are similar with our previous studies on other important TE semiconductors [20–22], which is explained in the Supplementary Material (SM).

Previous theoretical study showed that (1-1-1)[-11-2] is the least stress slip system for  $\text{Mg}_2\text{Si}$  [23], suggesting that the most plausible slip plane for  $\text{Mg}_2\text{Si}$  is the {1-1-1} plane. Thus, the nanotwinned  $\text{Mg}_2\text{Si}$  was constructed using the {1-1-1} plane as the twin plane, and the nanotwinned model was sheared along twin plane to investigate the role of TB on the mechanical properties. This nanotwinned structure was observed experimentally, suggesting it can be stably formed [15]. We considered the TB plane made of Si atoms, as observed in previous experiments [15]. The TB might be formed by Mg atoms, but this situation will not be considered in this paper since no experimental evidences have shown the existence of this type of TB. For comparison, we investigated the shear deformation of single crystal  $\text{Mg}_2\text{Si}$  along the most plausible slip system (1-1-1)[-11-2] and its opposite slip system (1-1-1)[1-12]. In the initial nanotwinned  $\text{Mg}_2\text{Si}$  (Fig. S2(a)) [15], the pre-stress might exist in the nanotwinned structure. However, structural relaxation allows the local atomic movement at the twin boundary (TB) to release the pre-stress. Thus, the relaxed structure (Fig. S2(b)) is stress-free, which is adopted for the shear simulations. In addition, the neighbor structure at the TB remains ideal structural symmetry, suggesting the initial TB structure can maintain perfect TB interface after relaxation, as shown in Fig. S3 in the SM.

We first investigate the shear responses of nanotwinned  $\text{Mg}_2\text{Si}$  under pure shear deformation, as shown in Fig. 1. We predicted an elastic moduli of 23.8 GPa for nanotwinned  $\text{Mg}_2\text{Si}$  from the slope of stress-strain relations at small strains (Fig. 1(a)). Under pure shear load, the theoretical shear strength of the nanotwinned  $\text{Mg}_2\text{Si}$  is found to be 0.93 GPa at the shear strain of 0.05. For single crystal  $\text{Mg}_2\text{Si}$ , the predicted elastic moduli is 46.4 GPa, which is twice higher than that (23.8 GPa) of its nanotwinned structure. The local atomic rearrangement in the nanotwinned structure leads to the enlarged Mg—Si bond length (2.88 Å) at the TB compared with that (2.75 Å) in the single crystal  $\text{Mg}_2\text{Si}$  (Fig. S2). This gives rise to a more weakened TB rigidity

compared with that of the Mg—Si framework in the single crystal  $\text{Mg}_2\text{Si}$ . Thus, the TB structure is much weaker in resisting the shear deformation compared with flawless  $\text{Mg}_2\text{Si}$ , leading to a much lower elastic modulus of nanotwinned structure than the flawless  $\text{Mg}_2\text{Si}$ . In addition, we calculated the elastic constant matrix of single crystal and nanotwinned  $\text{Mg}_2\text{Si}$ , and then estimated the shear modulus ( $G$ ) using the Voigt-Reuss-Hill method [24]. We found that the shear modulus of nanotwinned structure is 32.4 GPa, which is much lower than that (50.6 GPa) of single crystal  $\text{Mg}_2\text{Si}$ . This is another evidence showing that the elastic modulus decreases with nanotwinning. The ideal shear strength of  $\text{Mg}_2\text{Si}$  along (1-1-1)[-11-2] is 6.88 GPa, which is lower than that (8.63 GPa) along its opposite direction (1-1-1)[1-12]. But it is much higher than that (0.93 GPa) of the nanotwinned model, indicating that nanotwinning significantly softens  $\text{Mg}_2\text{Si}$ . The twin softening effects was also observed in boron rich boron carbide ( $\text{B}_{13}\text{C}_2$ ) [25].

To examine shear stress changes along different directions during shearing, we extracted the stress tensor of nanotwinned  $\text{Mg}_2\text{Si}$  during pure shear process, as shown in Fig. S4. After relaxation (at 0.0 shear strain), all the shear stress is zero, indicating the nanotwinned  $\text{Mg}_2\text{Si}$  is a stress-free cell. The maximum residual stress is  $\sigma_{yy} = -0.069$  GPa at 0.071 shear strain, which is much less than the value (0.1 GPa) of our convergence criterion on residual stress. Fig. S4 clearly shows that the residual stress along other five directions is negligible.

Under biaxial shear (compression + shear) load, the theoretical strength of nanotwinned  $\text{Mg}_2\text{Si}$  is 0.78 GPa, which is slightly lower than that (0.93 GPa) under pure shear load. This suggests that compression plays a minor role in the mechanical strength of nanotwinned  $\text{Mg}_2\text{Si}$ . However, we find that compression plays an essential role in determining the strength of flawless  $\text{Mg}_2\text{Si}$ , since under biaxial shear load the mechanical strength of  $\text{Mg}_2\text{Si}$  along (1-1-1)[-11-2] is 3.88 GPa, much lower than its ideal shear strength of 6.88 GPa.

Among all the shear-stress – shear-strain relations (Fig. 1), the shear stress monotonically increases with the increasing shear strain. Beyond the maximum shear stress point, the shear stress gradually decreases, suggesting a typical structural softening.

To character this structural softening and understand the deformation mechanism, we examined the atomic configurations and bond changes of single crystal  $\text{Mg}_2\text{Si}$  against shear strain along the (1-1-1)[-11-2] system, as shown in Fig. 2. Under pure shear load, the Mg—Si framework uniformly resists the deformation (Fig. 2(a)) as the shear strain increases to 0.21 corresponding to the ideal strength. The Mg<sub>1</sub>—Si<sub>1</sub> bond is shrunk from 2.75 to 2.65 Å with a shrinking ratio of 3.6%, while the Mg<sub>2</sub>—Si<sub>1</sub> bond is stretched from 2.75 to 2.84 Å with a stretching ratio of 3.3% (Fig. 2(c)). As the shear strain increases to 0.41, the Mg—Si framework was further sheared to accommodate the

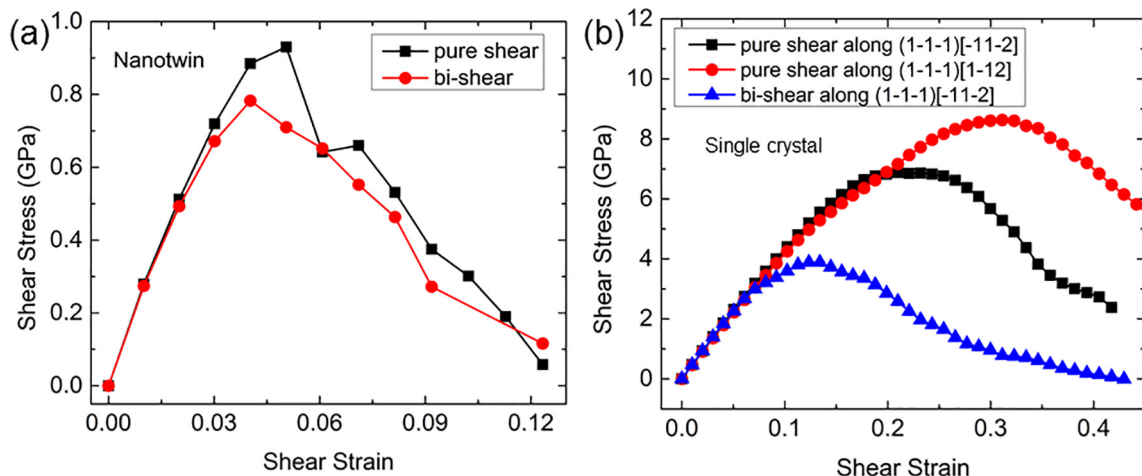


Fig. 1. Computed shear-stress – shear-strain relations of (a) nanotwinned  $\text{Mg}_2\text{Si}$ , as well as a comparison with that of (b) single crystalline  $\text{Mg}_2\text{Si}$  under pure shear and bi-shear loads.

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