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# Performance of microcrystalline silicon single and double junction solar cells of different degrees of crystallinity

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#### 1. Introduction

Hydrogenated microcrystalline silicon (µc-Si:H) solar cells generally show decreasing open-circuit voltage  $(V_{oc})$  with increasing crystalline volume fraction,  $F_c$  (Fig. 1). In a previous article [1] we had identified three contributing factors to the decrease of  $V_{oc}$  with increasing  $F_c$ , namely, an increase in both mid-gap and tail defects as well as a decrease in the mobility gap. We also explained that in the limiting case of crystalline silicon cells, with an even lower band gap, a higher  $V_{oc}$  is possible due to a sharply reduced gap defect density and a lower effective density of states  $(N_c, N_v)$  at the band edges. In this work, we model fluorinated  $\mu$ c-Si:H solar cells [2], with very high  $F_c$  as well as a significant fraction of large grains  $(F_{lg})$ , which appears to violate the general rule (Fig. 1). A material containing large grains is a dense material with low oxygen content, which can be obtained when using SiF<sub>4</sub> based plasma processes [2]. µc-Si:H with (i)  $F_c = 79\%$ ,  $F_{lg} = 0\%$ , (ii)  $F_c = 93\%$ ,  $F_{lg} = 27\%$  and (iii)  $\mu$ c-Si:F:H  $(F_c = 100\%, F_{lg} = 50\%)$  are used as the second I-layer in hypothetical double junctions and the optimized efficiencies are calculated. Three methods of thickness optimization are discussed.

#### 2. Experiments and modeling

Microcrystalline PIN solar cells having the structure: textured ZnO/P-µc-Si:H/I-µc-Si:H/N-a-Si:H/Aluminum have been deposited

### ABSTRACT

Hydrogenated microcrystalline silicon ( $\mu$ c-Si:H) solar cells show decreasing open-circuit voltage with increasing crystalline volume fraction ( $F_c$ ). Therefore, most research groups focus on  $\mu$ c-Si:H films produced close to the transition with amorphous silicon. Three cells with varying  $F_c$  and large-grain fraction,  $F_{lg}$  have been studied experimentally and by modeling. The large grained fluorinated  $\mu$ c-Si:H cell appears to be an exception to the above rule, and modeling is used to understand why low and high  $F_c$   $\mu$ c-Si:H and  $\mu$ c-Si:F:H are used as the second intrinsic-layer in hypothetical double junction cells. Modeling suggests possible efficiencies of 11.3%, 8.4% and 13%, respectively. Three methods of double junction cell thickness optimization are discussed.

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in a multiplasma monochamber RF-PECVD reactor, employing different ratios of silane to hydrogen flow rates during the intrinsic (I) layer deposition –  $SiH_4$ :  $H_2$  = 4:200, and  $SiH_4$ :  $H_2$  = 6:200 [1]. This results in different  $F_c$  and  $F_{lg}$  (sa and sb in Table 1). The P- $\mu$ c-Si:H and N-a-Si:H layers were deposited under the same conditions for the two cases. A third set (sc) has a fluorinated µc-Si:H I-layer, where hydrogen plasma treatment has been done both before and after deposition of the I-laver, since the latter procedure is shown to improve cell performance [2]. The composition of the films was obtained from the Bruggeman effective medium approximation modeling [3] of the pseudo-dielectric function of the films, deduced from spectroscopic ellipsometry measurements. This approach has been shown to be well adapted to the modeling of  $\mu$ c-Si films [4]. In this approach we used as components for the model the dielectric functions of (i) amorphous silicon; (ii) voids to take into account the porosity of the films; (iii) small grain  $\mu$ c-Si; and (iv) large grain µc-Si material produced by chemical vapor deposition (CVD) at high temperature. (The dielectric function of large grain µc-Si taken corresponds to that of a material produced by CVD at high temperature ( $\sim$ 650 °C) – Ref [5].) The presence of significant fraction of large grains in a material produced by plasma enhanced CVD (PECVD) at 175 °C has been shown to correlate with improved transport properties of the films [6]. In Table 1 we report the values of the total crystalline fraction  $F_c$ (which is the sum of the small grain and large grain fractions) and  $F_{lg}$ , which is the parameter of interest in this study. Indeed, achieving a high value of  $F_{lg}$  indicates that the films obtained by PECVD at 175 °C have a grain size similar to that of polycrystalline layers produced by CVD at  $\sim$ 650 °C [5].

The one-dimensional electrical–optical model Amorphous Semiconductor Device Modeling Program (ASDMP) [7] used in

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**Fig. 1.** The  $V_{oc}$  of silicon solar cells with varying Raman crystallinity.

#### Table 1

Comparison between the measured (Expt) and modeled (Model) solar cell output in PIN devices (sa, sb, sc), having different  $F_c$  and  $F_{lg}$ . da, db and dc are hypothetical double junctions with an a-Si:H top-cell and different  $\mu$ c-Si:H materials in the bottom-cell. The optimized thickness (in  $\mu$ m) of the top and bottom sub-cell I-layers in double junctions are  $d_1/d_2$ =0.23/1.71 in da,  $d_1/d_2$ =0.147/1.01 in db and 0.7/1.2 in dc, showing the importance of an intermediate reflector when using this high  $F_{lg}$  material.

Case	$F_c$ (%)	$F_{lg}\left(\%\right)$		$J_{SC}$ (mA/cm <sup>2</sup> )	$V_{OC}\left(V\right)$	FF	η (%)
sa	79	0	Expt	16.76	0.54	0.687	6.22
			Model	16.76	0.54	0.685	6.20
sb	93	27	Expt	19.94	0.46	0.631	5.79
			Model	19.85	0.45	0.640	5.72
SC	100	50	Expt	23.2	0.523	0.680	8.30
			Model	22.9	0.526	0.704	8.50
da				10.86	1.45	0.72	11.34
db				9.47	1.27	0.70	8.42
dc				13.92	1.37	0.69	13.16

this study solves Poisson's equation and the two carrier continuity equations under steady state conditions for the given device structure and yields the resulting *J*–*V* characteristics and the quantum efficiency. The electrical part of the modeling program is described in Ref. [8]. The gap state model used in our calculations consists of the tail states and Gaussian distribution functions to simulate the deep dangling bond states. The junction between the P–I–N cells of the stacked structure is modeled by Palit and Chatterjee [9], and Hou et al. [10] a heavily defective "recombination" layer (RL) with a reduced mobility band gap. The potential barriers for carriers moving towards this junction are reduced by band gap grading. Fig. 2 represents the blow-up of the area surrounding this inter-cell contact region and our approach to simulating the junction region.

The generation term in the continuity equations has been calculated using a semi-empirical model, that has been integrated into the modeling program [7]. Both specular interference effects and diffused reflectances and transmittances due to interface roughness are taken into account. The complex refractive indices for each layer of the structure are required as input and have been measured in-house by SE.

#### 3. Results and discussion

The experimental solar cell output parameters for "low  $F_c$ ", "high  $F_c$ "  $\mu$ c-Si:H and fluorinated  $\mu$ c-Si:H are compared to the modeling results in Table 1. From these simulations, typical parameters that characterize the intrinsic  $\mu$ c-Si:H layers of these



**Fig. 2.** Energy band diagram under short-circuit conditions and AM1.5 bias light around the inter-cell recombination layer (RL – shaded black in the figure), indicating our approach for simulating this contact region. The conduction band edge ( $E_c$ ) and the valence band edge ( $E_v$ ) are indicated.

cells have been extracted and are given in Table 2. Modeling indicates that high  $F_c$  µc-Si:H cell (sb) has a lower band gap, higher carrier mobilities, and both higher mid-gap defect density and broader band tails as compared to the cell having low  $F_c$  µc-Si:H (sa) [1]. In order to simulate the improved cell performance (Table 1) of the  $\mu$ c-Si:F:H cell, we had to take higher carrier mobilities, lower mid-gap defects and lower effective DOS at the band edges (similar to crystalline silicon). Indeed, the use of  $SiF_4$ as a gas precursor allows to reduce the concentration of oxygen in the films and leads to a material with a higher fraction of large grains. This in turn results in better transport properties. In particular we have shown that the electron mobility, as measured by time resolved microwave conductivity, increases with the fraction of large grains [6]. The lower defect density in  $\mu$ c-Si:F:H is confirmed by Fourier transform photocurrent spectroscopy (FTPS) measurements (not yet published); and a decrease of the band gap of  $\mu$ c-Si:H with increasing  $F_c$  (Table 2) has been observed by Hamma and Roca i Cabarrocas [11]. The vastly improved spectral response at long wavelengths confirms the very high crystalline volume fraction of µc-Si:F:H [2]. Since the highly crystallized  $\mu$ c-Si:F:H cell violates the general trend of lower V<sub>oc</sub> with increasing  $F_{c}$ , we checked the experimental results on various sets of solar cells all having very high  $F_c$  and  $F_{lg}$  and found the results reproducible. These cell results have been confirmed independently by measuring at the Jülich research center.

The parameters of  $\mu$ c-Si:H (low and high  $F_c$ ) and of  $\mu$ c-Si:F:H are then used as input to calculate the efficiency of hypothetical double junctions (da, db, and dc, respectively), with a standard a-Si:H top-cell. After suitable thickness optimization, the cell outputs are given in Table 1. The results indicate that the double junction cell dc, with fluorinated  $\mu$ c-Si:H as bottom sub-cell has the potential to give a higher efficiency ( $\eta \sim 13\%$ ), compared to the low  $F_c$  cell da ( $\eta \sim 11.3\%$ ) and the high  $F_c$  cell db( $\eta \sim 8.4\%$ ), but requires an exceptionally thick top-cell. This is due to the fact that current capacity of the bottom cell with  $\mu$ c-Si:F:H is very high (22.9 mA/cm<sup>2</sup> in the single junction, sc); so to match the current between the two sub-cells, the first cell has to be very thick. This suggests that we need an intermediate reflector to selectively reflect the short wavelength light into the top cell, so that higher current can be achieved with a thinner top sub-cell.

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