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Investigation of crack tip dislocation emission in aluminum using multiscale molecular dynamics simulation and continuum modeling



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ABSTRACT

This work investigates the dislocation nucleation processes that occur at the tip of a crack in aluminum under a broad range of crystallographic orientations and temperatures. A concurrent multiscale molecular dynamics – continuum simulation framework is employed. The results are then interpreted using a Peierls continuum model that uses finite temperature material properties derived from molecular dynamics simulation. Under ramped loading, partial dislocation nucleation at the crack tip is found to lead to both full dislocation emission and twinning, depending upon the orientation, temperature, and magnitude of the applied load in the simulation. The origins of the dependencies are made apparent by the Peierls continuum model. The continuum model suggests that in many instances dislocation nucleation from the crack tip can be considered to be a strain rate independent process, yet still temperature dependent through the temperature dependence of the stacking fault energies and elastic constants.

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1. Introduction

The ability to accurately model material failure is a key parameter to the analysis and design of advanced engineered systems. In practice, failure models are typically phenomenological in character, based on experimentally measured macroscopic quantities, such as ultimate strength, limit strain, and toughness. While phenomenological failure models have merits, their application to situations outside of the data set to which they were trained can lead to catastrophic errors. As a result, the engineering research community has long pursued the development of more robust failure models that are more closely linked to the underlying physics. Thanks to continually increasing computing resources, improved algorithms, and advanced simulation methodologies (atomistic, phase field, finite element, mesoscale models, etc.), the prospect for further improvement remains strong.

The physics of material failure in engineering alloys begins at the atomic scale, where interatomic decohesion and defect formation, propagation, and interaction occur. At the microstructural scale, the net effect of these mechanisms emerges as

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 $K_{II}^{eff^*}$

resolved shielded stress intensity factor on the

active slip plane

Nomenclature

| | ao | lattice constant of aluminum at 0 K $(a_{1} - 0.405 \text{ nm})$ | K ^{eff*} II crit | resolved shielded stress intensity factor on the active slip for partial dislocation nucleation |
|---|-------------------------|------------------------------------------------------------------|------------------------------|-------------------------------------------------------------------------------------------------|
| | b_p | magnitude of the Burgers vector of a Shockley | Q_{2D} | activation energy per unit length from Peierls |
| | C11 C12 | C_{44} single crystal elastic constants | 0 ₃₀ | total activation energy |
| l | d_{rr} | size of the finite element domain in the | r | distance from the crack tip |
| l | TL | Embedded Statistical Coupling Model | r_p | distance from crack tip to partial dislocation |
| | | (ESCM) model | S | shear stress on the slip plane due to slip |
| l | d _{MD} | size of the molecular dynamics domain in the | Т | system temperature |
| | mb | ESCM model | U | energy of an infinite linear elastic solid with |
| | F | energy factor | | a sharp crack |
| | Н | thickness of the molecular dynamics domain | U_o | energy of an infinite linear elastic solid with |
| l | | in the ESCM model used in the simulations | | a sharp crack with no slip |
| | h | interplanar spacing between the slip planes in | γssf | stable stacking fault energy |
| l | | aluminum | Υusf | unstable stacking fault energy |
| l | k_B | Boltzmann constant | γstf | stable twinning fault energy |
| l | K_I | mode I stress intensity factor | γutf | unstable twinning fault energy |
| l | K_{IC1} | stress intensity factor for leading partial dis- | Δ | relative shear displacement between two |
| l | | location nucleation | \rightarrow | atomic planes |
| l | K_{IC2} | stress intensity factor for complementary par- | δ | slip discontinuity in continuum model |
| l | | tial dislocation nucleation | δ_e | magnitude of edge component of a slip |
| l | K_{IT} | stress intensity factor for trailing partial dis- | | discontinuity |
| l | 1.01 | location nucleation | δ_s | magnitude of screw component of a slip |
| l | K_I^{LSL} | stress intensity factor below which a stable | | discontinuity |
| l | | state for the leading partial dislocation does | ε | applied far-field uniaxial strain |
| l | AT | not exist | θ | slip plane inclination angle with respect to the |
| l | K_{IC1}^{III} | stress intensity factor for leading partial dis- | æ | crack plane |
| l | T AT | location nucleation without thermal activation | Ψ | shear modulus |
| l | K_{IC2}^{III} | stress intensity factor for complementary par- | μ | Sileal illouulus |
| l | | tial dislocation nucleation without thermal | ν | shear stress on the slip plane in the direction |
| l | vAT | activation stress intensity factor for twinning partial | 1 | of $\vec{\delta}$ |
| l | K _{IT} | dialogation purchastion without thermal | (0 | angle between the slip direction and the |
| l | | activation | Ψ | normal to the crack front in the slip plane |
| | KP | most probable stress intensity factor for par | ع | distance from the crack tip |
| | N _I | tial dislocation nucleation | Ψ | generalized stacking fault energy |
| | K ^{eff} | resolved stress intensity factor on the active | ω | expected rate of thermally activated event |
| | * 11 | slin plane | ω_0 | attempt frequency of thermally |
| | | oub brane | | activated event |
| 1 | | | | |

microcracking and crystal plasticity. Ultimately these processes result in macroscopic failure, amenable to traditional continuum plasticity and fracture mechanics. The goal of physics based modeling is to acknowledge and understand this hierarchy. By doing so, it aims to be more predictive, robust, transferable, and potentially provide the insight necessary for creating new materials with improved properties.

Successful physics based failure modeling requires a sound understanding of the controlling mechanisms. A technologically important and illustrative example is the prediction of crack growth in engineering alloys. Crack growth in metals occurs via interatomic decohesion at the crack tip, a mechanism that is in competition with crack tip plasticity in the form of dislocation nucleation and propagation. Additional complexity exists in face-centered-cubic (f.c.c.) alloys, where crack tip dislocation emission can consist of either crystallographically restoring plastic slip in the form of full dislocations or crack tip twinning due to the nucleation of successive partial dislocations (Weertman and Weertman, 1992). While both mechanisms are associated with plastic deformation, they can ultimately lead to drastically different crack tip behavior, influencing the overall material toughness. Accordingly, a key step, necessary for the development of predictive physics based crack growth models in f.c.c. metals, is to develop a comprehensive understanding of the competition between crack tip twinning and full dislocation emission.

Modeling this competition presents a particular challenge in that it is sensitive to strain rate, temperature, and crystallographic orientation (Farkas et al., 2001; Hai and Tadmor, 2003; Yamakov et al., 2006; Warner et al., 2007; Warner and Curtin, 2009; Venables, 1963). This paper reports on the authors' attempts to more comprehensively understand the Download English Version:

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