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Integrability and generalized center problem of resonant singular point



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ABSTRACT

In this paper, integrability and generalized center condition of resonant singular point for a broad class of complex autonomous polynomial differential system are studied. A new method—integrating factor method of determining integrability of resonant singular point is obtained for any rational resonance ratio. At the same time, the relations of the first integral method and the integrating factor method with the normal form method are obtained. © 2014 Elsevier Ltd. All rights reserved.

1. Introduction

Recently, the classical center problem for real planar polynomial differential system

$$\frac{dx}{dt} = -y + P(x, y), \qquad \frac{dy}{dt} = x + Q(x, y), \tag{1}$$

where P(x, y), Q(x, y) are polynomials of x and y with degree no less than 2, is generalized by [1] to the system

$$\frac{dx}{dt} = px + P(x, y), \qquad \frac{dy}{dt} = -qy + Q(x, y), \tag{2}$$

which has a p:-q resonant singular point at the origin. When P and Q are special polynomials and the resonance ratio is particular, integrability of system (2) has been studied by quite a few people (see, e. g., [1–10] for quadratic polynomials, [11–17] for cubic polynomials, [18,19] for quartic polynomials, and [20] for quintic polynomials) and there is no lack of excellent results such as [1,2,4].

In this paper, based on the previous works, we study further the integrability and generalized center condition of resonant singular point for complex polynomial differential system

$$\begin{cases} \frac{dz}{dT} = pz + \sum_{\alpha+\beta=2}^{\infty} a_{\alpha\beta} z^{\alpha} w^{\beta} = Z, \\ \frac{dw}{dT} = -qw - \sum_{\alpha+\beta=2}^{\infty} b_{\alpha\beta} w^{\alpha} z^{\beta} = -W, \end{cases}$$
(3)

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where z, w, T are independent complex variables, $a_{\alpha\beta}$, $b_{\alpha\beta}$ are complex constants, p, q are positive integers, and (p,q)=1. In Section 2, we introduce some preliminary results. In Section 3, we introduce a new method—integrating factor method of determining integrability of resonant singular point and get another necessary and sufficient condition of determining integrability. What is more, we obtain the relations of the first integral method and the integrating factor method with the normal form method.

2. Some preliminary results

Lemma 2.1 (See [21,4,5] etc.). We can derive successively the following formal series

$$\begin{cases} u = z + \sum_{\alpha+\beta=2}^{\infty} A_{\alpha\beta} z^{\alpha} w^{\beta} \\ v = w + \sum_{\alpha+\beta=2}^{\infty} B_{\alpha\beta} w^{\alpha} z^{\beta}, \end{cases}$$

$$(4)$$

such that system (3) be transformed into its normal form:

$$\begin{cases} \frac{du}{dT} = pu \sum_{k=0}^{\infty} p_k (u^q v^p)^k \\ \frac{dv}{dT} = -qv \sum_{k=0}^{\infty} q_k (u^q v^p)^k, \end{cases}$$
(5)

with $p_0 = q_0 = 1$.

Definition 2.2 ([5]). We call $\mu_k = p_k - q_k$ "generalized singular point quantity of order k" at the origin of system (3). If $\mu_1 = \mu_2 = \cdots = \mu_{k-1} = 0$, $\mu_k \neq 0$, then we call the origin a fine singular point of order k; If for all k, $\mu_k = 0$, then we call the origin a generalized center.

Remark. (i) If system (3) is a real system, then " μ_k " defined in Definition 2.2 is "the saddle point quantity of order k" defined in many other works (see e.g. [21,4] and references there in).

- (ii) $\mu_0 = 0$ is a direct result of $p_0 = q_0 = 1$.
- (iii) If p = q = 1, and $a_{\alpha\beta} = b_{\alpha\beta}$, then under transformation z = x + iy, w = x iy, T = it, $i = \sqrt{-1}$, system (3) becomes system (1). Systems (3) and (1) are concomitant each other. And the singular point quantities of system (3) are actually the focal values of system (1) except for possible constant factor which do not affect the order of singular point or focus.

Theorem 2.3 ([5]). We can derive successively the following formal series

$$F(z, w) = z^q w^p + h.o.t., \tag{6}$$

such that

$$\left. \frac{dF}{dT} \right|_{(3)} = \sum_{k=1}^{\infty} \lambda_k (z^q w^p)^{k+1}. \tag{7}$$

Here and throughout the paper 'h.o.t.' represents 'higher order terms'.

Theorem 2.4 ([5]). The origin of system (3) is a generalized center if and only if there is an analytic first integral.

3. Singular point quantity and integrability

Theorem 3.1. The $\lambda_k's$ appeared in Eq. (7) satisfy the following: if $\mu_1 = \mu_2 = \cdots = \mu_{k-1} = 0$, $\mu_k \neq 0$, then $\lambda_1 = \lambda_2 = \cdots = \lambda_{k-1} = 0$, $\lambda_k \neq 0$, $\lambda_k = pq\mu_k$ for all $k \geq 1$, and vice versa.

Proof. Let $F = u^q v^p = z^q w^p + h.o.t.$, then from Eq. (4), normal form (5) and Eq. (7) we have

$$\frac{dF}{dT}\Big|_{(3)} = \frac{dF}{dT}\Big|_{(5)} = pq \sum_{k=1}^{\infty} (p_k - q_k) (u^q v^p)^{k+1} = pq \sum_{k=1}^{\infty} \mu_k (u^q v^p)^{k+1}
= pq \sum_{k=1}^{\infty} \mu_k [(z^q w^p)^{k+1} + h.o.t] = \sum_{k=1}^{\infty} \lambda_k (z^q w^p)^{k+1}.$$
(8)

From Eq. (8) we can easily get that if $\mu_1 = \mu_2 = \cdots = \mu_{k-1} = 0$, $\mu_k \neq 0$, then by mathematical induction we get $\lambda_1 = \lambda_2 = \cdots = \lambda_{k-1} = 0$, $\lambda_k \neq 0$, $\lambda_k = pq\mu_k$ for $k \geq 1$, and vice versa.

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