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# CO<sub>2</sub> trapping in amorphous H<sub>2</sub>O ice: Relevance to polar mesospheric cloud particles



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#### ARTICLE INFO

# ABSTRACT

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Keywords: Amorphous ice Carbon dioxide trapping Temperature programmed desorption Polar mesospheric clouds Polar mesospheric clouds form in the summer high latitude mesopause region and are primarily comprised of H<sub>2</sub>O ice, forming at temperatures below 150 K. Average summertime temperatures in the polar mesosphere (78°N) are approximately 125 K and can be driven lower than 100 K by gravity waves. Under these extreme temperature conditions and given the relative mesospheric concentrations of CO<sub>2</sub> and H<sub>2</sub>O (~360 ppmv and ~10 ppmv, respectively) it has been hypothesised that CO<sub>2</sub> molecules could become trapped within amorphous mesospheric ice particles, possibly making a significant contribution to the total condensed volume. Studies of CO<sub>2</sub> trapping in co-deposited gas mixtures of increasing CO<sub>2</sub>:H<sub>2</sub>O ratio (deposited at 98 K) were analysed *via* temperature programmed desorption. CO<sub>2</sub> trapping was found to be negligible when the H<sub>2</sub>O flux to the surface was reduced to  $4.8 \times 10^{13}$  molecules cm<sup>-2</sup> s<sup>-1</sup>. This corresponds to an average of 0.4 H<sub>2</sub>O molecules depositing on an adsorbed CO<sub>2</sub> molecule and thereby trapping it in amorphous ice. Extrapolating the experimental data to mesospheric conditions shows that a mesospheric temperature of 100 K would be required (at a maximum mesospheric H<sub>2</sub>O concentration of 10 ppmv) in order to trap CO<sub>2</sub> in the ice particles. Given the rarity of this temperature being reached in the mesosphere, this process would be an unlikely occurrence.

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### 1. Introduction

Polar mesospheric clouds (PMCs) (also known as noctilucent clouds) are optically thin clouds that form in the summer polar mesopause region. PMCs primarily consist of small H<sub>2</sub>O ice particles with median radii of  $\sim$ 51 nm (von Cossart et al., 1999; Hervig et al., 2001), which can form when the temperature falls below the water frost point (150 K). Larger ice particles which sediment to the bottom of the ice layer, at altitudes between 82 and 85 km, scatter sufficient light to become optically visible (Rapp and Thomas, 2006).

A major source of ice nuclei for heterogeneous nucleation in the mesosphere are meteoric smoke particles (MSPs) (Rapp and Thomas, 2006; Plane, 2011). MSPs are a product of ablation and subsequent recondensation of meteoric material (Saunders and Plane, 2006), of which somewhere between 10 and 110 tonnes enters the atmosphere daily (Plane, 2012). Average temperatures in the high latitude upper mesosphere (78°N) during summertime are approximately 125 K, while gravity waves can drive temperatures lower than 100 K (Lübken et al., 2009). Murray and Jensen (2010) showed that homogeneous nucleation may also be possible

\* Corresponding author. *E-mail address:* j.m.c.plane@leeds.ac.uk (J.M.C. Plane). under extreme mesospheric temperature conditions below 110 K. These authors also suggested that amorphous ice may preferentially nucleate due to a lower energy barrier to nucleation compared to crystalline ice. Satellite measurements indicate PMCs are composed primarily of crystalline ice but could not rule out amorphous ice formation (Hervig and Gordley, 2010).

H<sub>2</sub>O is present in the high latitude mesopause during the summertime at concentrations of up to 10 ppmv (Rong et al., 2010). The source of  $H_2O$  vapour in the mesosphere is due to a combination of atmospheric upwelling and methane oxidation via reaction with O(<sup>1</sup>D) and OH radicals (Thomas, 1991). CO<sub>2</sub> is well mixed vertically with concentrations > 360 ppmv up to altitudes of approximately 80 km, above which UV photolysis of CO<sub>2</sub> (to CO and O) and gravitational separation with respect to lighter gas species become dominant removal processes. Due to increasing anthropogenic sources of CO2, a global decadal increase of  $23.5 \pm 6.3$  ppmv is currently observed (Emmert et al., 2012). Given the concentrations of trace species present in the mesosphere during the PMC cloud season and the extreme low temperatures, it is possible that CO<sub>2</sub> trapping in amorphous solid water (ASW) ice could be a viable process. However, it should be noted CO<sub>2</sub> cannot become trapped when ASW has irreversibly crystallised to crystalline solid water (CSW) ice, or when ice is initially deposited as CSW ice.

CO<sub>2</sub> trapping in ASW ice has been experimentally observed

under conditions relevant to the interstellar medium (ISM) (Galvez et al., 2007; Malyk et al., 2007; Mate et al., 2008). However, these studies looked at deposition temperatures below atmospheric relevance (  $\leq$  90 K) and low CO<sub>2</sub>:H<sub>2</sub>O ratios (~0.02:1 to 0.7:1) due to the composition of interstellar ice. Experimental studies have also investigated CO<sub>2</sub> trapping in the Martian regolith (Trainer et al., 2010), but these were at higher temperatures and pressures than those observed during the PMC cloud season. To our knowledge, studies of co-deposition of high CO2:H2O ratio gas mixtures at conditions relevant to the Earth's mesosphere have not been performed. In the present study, the co-deposition of high CO<sub>2</sub>:H<sub>2</sub>O ratio mixtures were studied using temperature programmed desorption (TPD) to determine the mesospheric conditions under which CO<sub>2</sub> trapping in ASW ice would be possible. The experimental procedure is described in Section 2, and the results and discussion in Section 3.

#### 2. Experimental procedure

The experiments were undertaken in an ultra-high vacuum (UHV) chamber with an internal volume of  $\sim$  25 L (see Fig. 1). This system has been previously described by Vondrak et al. (2006). Briefly, the chamber is pumped *via* a  $550 \text{ Ls}^{-1}$  turbo-molecular pump (Varian Turbo-V 551 Navigator) backed by a rotary pump (Varian, DS302), achieving a base pressure of typically < 7  $\times$  10<sup>-10</sup> mbar. The chamber is equipped with a quadrupole mass spectrometer (QMS) (Hiden, HAL 3F 301 RC PIC), a needle valve (NUPRO, SS4BK) for gas dosing and an inert ion (Ar<sup>+</sup>) sputter source (PSP technology Ltd., ISIS 3000) which was modified to function as a leak valve for this study. Within the centre of the UHV chamber is a cylindrical Cu (111) crystal sample of 12 mm diameter and 2.5 mm thickness, polished to 1 µm and oriented to  $\pm 0.5^{\circ}$  of the (111) plane. The crystal is mounted *via* tungsten (W) heating wires onto an oxygen free high conductivity (OFHC) cold finger (shown in Fig. 1), which in turn is mounted onto an  $xyz\theta$  manipulator. The sample is liquid nitrogen cooled to a base temperature of approximately 98 K and can be heated resistively by the W wires embedded in the perimeter of the Cu (111) crystal. The surface temperature is monitored using a K-type thermocouple positioned in a small hole on the side of the Cu (111) crystal. Prior to each experiment the Cu (111) crystal was annealed to 800 K for 30 min.

H<sub>2</sub>O (deionised) was subject to three freeze-pump-thaw cycles and CO<sub>2</sub> (from dry ice) was first purified by pumping. CO<sub>2</sub>:H<sub>2</sub>O gas mixtures at ratios varying from 1:1 to 13:1 were prepared using manometric techniques on a glass gas-handling line, and then dosed onto the Cu (111) crystal via the needle valve in the form of a collimated effusive beam. This beam has a diameter only slightly larger than the Cu (111) crystal and therefore minimises adsorption (and therefore desorption) of either species onto other nearby surfaces such as the OFHC cold finger (see Fig. 1). Any species adsorbed onto the W wires would desorb immediately upon heating. The dosing rates were calibrated for CO<sub>2</sub> and H<sub>2</sub>O separately through beam flux calibrations based on the procedure of Oakes (1994). Briefly, the beam flux ( $F_{beam}$ ) is calculated from the exponential decay of a gas species from the point where the input gas flow has been terminated. The exponential decay is described by:

$$P = P_0 \exp^{-C_m t} \tag{1}$$

where *P* is the pressure,  $P_0$  is the initial pressure at the termination of the dose,  $C_m$  is the pumping coefficient for the gas and *t* is the time from the start of the decay. Eq. (1) was linearised and  $C_m$  obtained from the gradient. From this  $F_{beam}$  is given by:



Fig. 1. Schematic diagram of the UHV chamber upper level (top) and Cu (111) crystal mount (bottom).

$$F_{beam} = \frac{V_{system}C_m P_0}{k_B T A_{beam}}$$
(2)

where  $V_{system}$  is the volume of the chamber,  $k_B$  is the Boltzmann constant, *T* the temperature of the gaseous species and  $A_{beam}$  the cross sectional area of the molecular beam. The CO<sub>2</sub> and H<sub>2</sub>O fluxes varied from  $(3.4 - 6.3) \times 10^{14}$  molecules cm<sup>-2</sup> s<sup>-1</sup> and  $(0.5-3.4) \times 10^{14}$  molecules cm<sup>-2</sup> s<sup>-1</sup>, respectively. The total flux of CO<sub>2</sub> and H<sub>2</sub>O remained the same (within experimental error) for all the experiments  $(6.8 \times 10^{14} \text{ molecules cm}^{-2} \text{ s}^{-1})$ .

In a typical experiment, the gas mixture was dosed for 7200 s onto the Cu (111) crystal at a base temperature of 98 K. This enabled the H<sub>2</sub>O film to form as ASW ice. The absorbed species were then analysed *via* TPD, where the Cu crystal was heated between 98 and 200 K at a heating ramp of 0.25 K s<sup>-1</sup>. This ensured that all dosed species had desorbed. The CO<sub>2</sub> TPD profiles described in the following section were corrected against a pure CO<sub>2</sub> TPD trace to remove the surface CO<sub>2</sub> desorption. This surface CO<sub>2</sub> is weakly physisorbed to the ASW ice surface and rapidly desorbs upon sample heating, peaking at temperatures of approximately 125 K, which is more than 20 K lower than the observed initial

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