

Strain effect on the electronic properties of Ce-doped SnS₂ monolayer

Guizhen Qiu^a, Huimin Zhang^a, Yaming Liu^{b,*}, Congxin Xia^{c,**}

^a Xinxiang Medical University, Xinxiang, 453003, China

^b Henan Institute of Science and Technology, Xinxiang 453003, China

^c Henan Normal University, Xinxiang 453003, China

ARTICLE INFO

Keywords:

First-principles
Ce-doped SnS₂ monolayer
Effect of strain
Electronic property

ABSTRACT

First-principles calculations based on density functional theory are carried out to investigate the formation energy, transition energy level, electronic structure and effect of strain on Ce-doped SnS₂ monolayer nanosheet. Numerical results show that the doped nanocompound is energetically stable, while the introduced dopant state is composed completely of Ce 4*f* electrons, and the adopted perfect SnS₂ monolayer leads to the disagreement with experimental Ce³⁺ ionic state. Applying biaxial strain from −10% to +10%, the doped system holds the indirect semiconducting characteristics in the strain range −4%~+6%. The variation tendency of bandgaps is almost to be a horizontal line about 1.9eV under compressive conditions, and the bandgaps increase slowly at first and then decrease rapidly beyond +4%. This work is useful for further improving Ce-doped SnS₂ nanostructures and paves the way in the potential application in photocatalysis and lithium-ion battery.

1. Introduction

Due to the peculiar geometrical structure, abundant reserves, excellent chemical stability in acid conditions, nontoxicity and environmental friendliness, the layered two-dimensional (2D) semiconductor SnS₂, has been receiving intensive attention for the applications in next-generation nanoelectronics, optoelectronics, photocatalysis and lithium-ion battery [1–9]. A series of experiments show that doping with Cu, Zn, In, Ag, etc, the composites exhibit enhanced visible light-driven photocatalytic activity, and can be candidates for environmental remediation and energy conversion [10–15]. Doping with V, W, Fe, etc, the bandgaps are engineered, and the systems can be used as solar energy absorbers in photovoltaic devices [16–20]. Recent experiments show that Ce-doped SnS₂ nanocomposites exhibit both higher photocatalytic activity and significantly improved electrode cycling performance and rate retention ability in lithium-ion batteries [21,22]. Except for these metal dopants, controlling the Se composition in SnSe_xS_{2-x} ternary alloys, the carrier mobility of the 2D nanosheets can be effectively tuned and showing excellent optoelectronic properties [23,24]. To address the characterization of the nature of these dopants, theoretical works are highly in demand, and tremendous reports are carried out to explore the mechanism and underlying reasons [25–35].

To the best of our knowledge, however, theoretical investigations related to the Ce-doped SnS₂ systems have not been reported yet. The

doped geometrical structures, variation of electrical properties, effect of strain as well as the mixing behavior, are still unclear from a theoretical point of view. In this work, we focus on the Ce-doped SnS₂ monolayer, and the mentioned issues will be systematically calculated under density functional theory (DFT) framework.

2. Computational methods

All the first-principles calculations in this work are carried out using VASP [36,37]. The Perdew-Burke-Ernzerhof (PBE) parameterization of the generalized-gradient approximation (GGA) is employed to model the exchange and correlation interactions [38,39]. Projected augmented-wave (PAW) method is adopted to describe the electron-ion potential [40]. A kinetic energy cutoff of 520eV is selected for the plane wave expansion, and the Γ -centered Monkhost-Pack *k*-points grid of $5 \times 5 \times 1$ and $11 \times 11 \times 1$ are chosen in the atomic and electronic calculations [41]. The geometrical structure of the crystal is determined by optimizing the atomic positions using CG method with a force convergence criterion of 10^{-3} Å to precisely harvest the energy bandgaps, the valence electron configurations considered here are S ($3s^2 3p^4$), Sn ($4d^1 05s^2 5p^2$) and Ce ($5s^2 5p^6 6s^2 4f^1 5d^1$). The effect of on-site coulomb repulsion of Sn 4*d* and Ce 4*f* states is treated by GGA + *U* approach [42], and the *U* values are choosed respectively as 9.0eV and 6.3eV for Sn 4*d* and Ce 4*f*. These values are successfully used in several

* Corresponding author. Henan Institute of Science and Technology, Xinxiang 453003, China.

** Corresponding author. Henan Normal University, Xinxiang 453003, China.

E-mail addresses: liuym@hist.edu.cn (Y. Liu), xiacongxin@htu.cn (C. Xia).

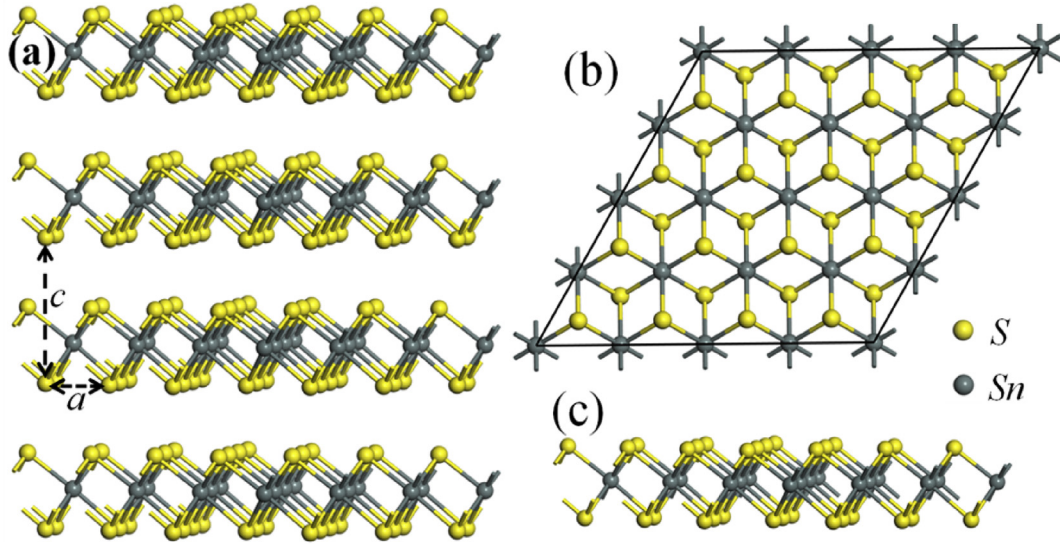


Fig. 1. Schematics of SnS_2 in different structures. (a) 2H-type ($P-3m1$), lattice parameters a and c are marked in dashed arrows (b) top and (c) side view of 4×4 supercell of monolayer SnS_2 .

present theoretical works [28,43]. To describe the non-bonding weak van der Waals interactions between the layers of bulk SnS_2 , the Grimme's DFT-D2 method is considered [44].

3. Results and discussions

3.1. Structures of bulk and monolayer SnS_2

Bulk SnS_2 is an n -type layered semiconductor, each layer has an S–Sn–S sandwich-like structure, and the layers are weakly coupled by van der Waals interactions. Periodically stacking along c -direction, the triple-layers result in many different polytypes. The simplest and most common polytype belongs to space group $P-3m1$ (No.164), and denoted as 2H– SnS_2 , (H-for hexagonal) see Fig. 1. Here, we choose the 2H– SnS_2 as the parent bulk crystal phase. Independent of the different bulk phase, the exfoliated monolayer SnS_2 nanosheet adopts only the space group $P-3m1$. To simulate the monolayer, a supercell consisting of 4×4 formula units with 16 Sn atoms and 32 S atoms (*i.e.* $\text{Sn}_{16}\text{S}_{32}$) is constructed.

The optimized geometrical and electronic structures of bulk and monolayer SnS_2 are listed in Table 1, and compared with previous experimental and theoretical data.

Due to the well-known shortcoming of DFT, the PBE functional overestimates the in-plane lattice constant a of the bulk phase by 1.4%, whereas the error of out-of-plane lattice parameter c is approximately 14.8%. Applying the Grimme's DFT + D2 method to the bulk phase, a proper description with error of c are less than 4% is obtained. As for the monolayer SnS_2 , PBE functional results in better lattice constant than that of PBE + U approach. However, compared with the reported experimental data, the effect of on-site Hubbard U can lead to excellent

bandgaps. To harvest moderate electronic properties, the PBE + U method will be applied to all the following calculations.

Fig. 1(c) shows the electronic band structures of monolayer and bulk 2H– SnS_2 . We find that the valence band maximum (VBM) of both panels locate along Γ –M high symmetry line and near the Γ -point. While the conduction band minimum (CBM) lies respectively at M and L point for monolayer and 2H– SnS_2 , indicating the indirect semiconducting characteristic. Due to the reduction of dimensionality from 3D-bulk to 2D-monolayer structure, the M point of monolayer is equivalent to L-point of bulk phase in the Brillouin zone [20]. The calculated bandgaps agree well with the present experimental reports. Rechecking the band structures of 2H– SnS_2 , the energy curves along the Γ –K–M– Γ are similar to that along A–H–L–A, which implies that the interaction along c -axis is weak, and coincides with the characteristic of layered van der Waals compounds.

3.2. Structure and stability of Ce-doped monolayer SnS_2

Recent experiments found that Ce-doped SnS_2 nanostructures exhibit higher photocatalytic activity and better performance in lithium-ion battery. To have a microscopic scenario and better understanding of the substitutional behavior, the formation energy (E_{form}), transition energy level, geometrical structure and electronic properties of Ce-doped monolayer SnS_2 is systematically discussed. Though the following equation

$$E_{\text{form}}(\text{Ce}^q) = E[\text{Ce}^q] - E[\text{SnS}_2] + \mu_{\text{Ce}} + q[E_{\text{F}} + E_{\text{V}} + \Delta V],$$

the formation energy is calculated, where $E[\text{Ce}^q]$ and $E[\text{SnS}_2]$ are the calculated total energies of Ce-doped and pristine monolayer SnS_2 with the same size. E_{F} is the Fermi level referenced to the energy position (eV) of the valence band maximum (VBM) of pristine monolayer SnS_2 . The correction term ΔV is used to align the reference potential of the doped compounds to that of the pristine supercell. μ_{X} is chemical potential of corresponding X ($X = \text{Sn}, \text{S}, \text{Ce}$) and the magnitudes of μ_{X} depend on growth conditions. In X-rich environments, μ_{X} is the total energy per atom in its reference phase. Here, we adopt the most energetic stable cubic phase for Ce (No.225) and Sn (No.227), while S_2 molecular is used for sulfur atom. To maintain a stable SnS_2 nanosheet, μ_{X} of S and Sn must satisfy $\mu_{\text{Sn}} + 2\mu_{\text{S}} = \mu_{\text{SnS}_2}$. μ_{SnS_2} is the total energy of one formula in nanosheet reference phase. Thus, the formation energies of the Ce-doped compounds are 1.249 eV and -3.201 eV for Sn-rich and S-rich growth conditions. The negative E_{form} under S-rich condition

Table 1

Lattice constants, S–Sn bond lengths $d_{\text{S-Sn}}$ and band gaps E_{g} of SnS_2 in different crystal structures calculated with different DFT frameworks, and compared with previous theoretical and experimental reports.

	Bulk (2H-phase)		monolayer	
	D2 + U (+ U)	Reference	PBE (+ U)	Reference
$a(\text{\AA})$	3.573 (3.701)	3.649 [1]	3.689 (3.482)	3.700 [3]
$c(\text{\AA})$	6.159 (6.774)	5.899 [1]	–	–
$d_{\text{S-Sn}}(\text{\AA})$	2.485 (2.596)	2.467 [28]	2.591 (2.437)	2.600 [3]
E_{g} (eV)	2.141 (1.660)	1.926 [28]	1.650 (2.362)	2.230 [3]

Download English Version:

<https://daneshyari.com/en/article/8160090>

Download Persian Version:

<https://daneshyari.com/article/8160090>

[Daneshyari.com](https://daneshyari.com)