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Validation of an efficiency calibration procedure for a coaxial n-type and a well-type HPGe detector used for the measurement ¹/₁₅ ¹/₂ of environmental radioactivity

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ABSTRACT

To obtain reliable measurements of the environmental radionuclide activity using HPGe (High Purity Germanium) detectors, the knowledge of the absolute peak efficiency is required. This work presents a practical procedure for efficiency calibration of a coaxial n-type and a well-type HPGe detector using experimental and Monte Carlo simulations methods. The method was performed in an energy range from 40 to 1460 keV and it can be used for both, solid and liquid environmental samples. The calibration was initially verified measuring several reference materials provided by the IAEA (International Atomic Energy Agency). Finally, through the participation in two Proficiency Tests organized by IAEA for the members of the ALMERA network (Analytical Laboratories for the Measurement of Environmental Radioactivity) the validity of the developed procedure was confirmed. The validation also showed that measurement of ²²⁶Ra should be conducted using coaxial n-type HPGe detector in order to minimize the true coincidence summing effect.

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1. Introduction

41 Many approaches have been proposed in the literature for effi-42 ciency calibration of HPGe detectors when environmental samples are 43 measured [1–3]. Set of standard point sources having single energy 44 emissions, standard solution of mixed radionuclides and Certified 45 Reference Materials commercially available, are commonly employed. 46 47 In most of these cases, the activity measured by this method has to be 48 corrected for summation effects induced by photons emitted in 49 coincidence and also for self-absorption when the measured sample 50 has a different matrix (density or composition) than the source used 51 in the calibration process. A direct calibration can be performed using 52 standard radioactive sources of the same geometrical dimensions, 53 density, and chemical composition, compared with the samples of 54 interest. However, standards are not often available for all environ-55 mental matrices or for all radionuclides of interest. 56

In addition, theoretical and computational methods have also been employed for efficiency calibration, coincidence-summing and self-absorption corrections. In several studies have been used general

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69 and specific software based on Monte Carlo (MC) codes like Geant, 70 71 72

MCNPX, EFFTRAN or DETEF [4–7]. A good agreement between experimental and calculated data can be reached using these codes; in addition the working time can be reduced considerably. In these cases, main limitations reside in the precise knowledge about the characteristics of the experimental geometry and sample compositions [8]. Generally, the efficiency values obtained experimentally and by MC simulation based on nominal values of the parameters supplied by the manufacturer show significant differences due to the inaccuracy in some critical parameter like the thickness of the dead layer or the active volume. The optimization of these parameters can result in a substantial decrease of the deviations between the experimental and calculated values [9,10]. However, even when precise geometrical data are available, it is necessary to refine the model by feeding it back with experimental results when accuracy is desired. This is because some parameters involved in the detector response cannot easily be assessed. They include the distribution of the electrical field in the crystal, its mounting and dimensions and properties of the dead layers [11,12].

One effective procedure to overcome these difficulties is to use an efficiency transfer function from reference geometry to other source configurations, using MC calculations and experimental measurements. This procedure consists of calculating the full energy peak efficiency (FEPE) by an energy dependent transfer 67

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factor, which is derived by comparing the direct calculated FEPEs with the source experimental values at a reference position. In the literature, different authors have reported differences below 5% after using the transfer function [13,14]; even when there was no adjustment of parameters of the detectors.

Independent of the chosen calibration method (experimental or computational), there is a group of limitations which will be required to overcome. However, the combination of these techniques can be a potential tool as a practical and cheap method for routine measurement purpose in many laboratories. The goal of this study is to develop a simple procedure for efficiency calibration of two different HPGe detectors, complementing experimental and MC simulation methods. The main advantages of this approach is that it can be applied to coaxial detectors as well as well detectors and it can be used for different environmental matrices. The method was verified measuring several IAEA reference materials and finally validated through the participation in two ALMERA Proficiency Tests organized by IAEA for the ALMERA network members.

2. Materials and methods

2.1. Detectors

Two high purity germanium (HPGe) detectors were used for experimental measurements. Detector 1 (D1) was a coaxial n-type detector (model NGC 3019, from DSG Detector Systems GmbH) with epoxy-carbon window and 31.5% of relative efficiency. It was coupled to an electronic chain, including a multichannel analyzer type TMCA. Detector 2 (D2) was a well-type HPGe detector (model EGPC100 P-15, from Canberra) with an absolute efficiency of 12.1% at 661 keV. The data acquisition system of this detector consists of a PSC822 preamplifier, Canberra amplifier model 7245 and electronic card MCA 5000 which includes a 7602 ADC with 8192 channels and InterFast multichannel analyzer. Both detectors are surrounded by a cylindrical low-background chamber made with the following elements from outside to the inner region: 240 mm of steel, 37 mm of lead, 1 mm of aluminum and 1 mm of copper. In both detectors the gamma spectra were recorded and analyzed using Winner[™] 6.0 software. The detectors resolution and energy calibration is periodically verified for stability using a set of point sources (²⁴¹Am, ¹³⁷Cs, ⁶⁰Co and ²²⁶Ra).

2.2. Monte Carlo simulation

The first step of the calibration procedure was to obtain the efficiency calibration curves for both detector (D1 and D2) by Monte Carlo simulation methods. It was employed the code MCNPX 2.6. The efficiency response for both detectors have been previously reported in the studies [6,15]. Nominal values of the parameters supplied by the manufacturer are described in those studies.

For detector D1 the Monte Carlo efficiency transfer method was applied using directly the manufacturer supplied data in all MC calculations. The efficiencies were calculated using the expression $\epsilon_x = \epsilon_{ref} \left(\epsilon_x^{MC} / \epsilon_{ref}^{MC} \right)$ where ϵ_x is the efficiency for a particular geometry and energy, ϵ_{ref} is the experimental efficiency for a reference case, and ϵ_{ref}^{MC} and ϵ_x^{MC} are calculated efficiencies (via Monte Carlo) for the reference case and the geometry of interest, respectively. The characteristics of the reference source and source-detector configuration are described in [6]. We used the detector model described in [6] to compute the efficiencies of the samples that were measured here (ϵ_x^{MC}) and we used the same values of ϵ_{ref} and ϵ_{ref}^{MC} measured and calculated in this work, respectively.

For detector D2 the Monte Carlo efficiency calculations were made directly. In this case, a tuning of some critical parameters of the detector was made. This tuning showed that the thickness of external dead layer (EDL) and distance between the Ge crystal and the Al end cap (DGA) are critical parameters and they were optimized. For the rest of parameters we used the nominal values.

We used the pulse-height tally (F8) per photons emitted from the source to compute the absolute efficiency and we generally obtained relative errors lower than 1% with a number of histories about 10^5 – 10^6 and 12–14 min of computational times for each energy. This computational time per energy allows us to build an efficiency calibration curve in approximately 4 h. Therefore, the proposed calibration method is good for practical application in everyday measurements. All MC calculations covered the energy range 40–1460 keV.

2.3. True coincidence summing corrections

The second step consisted of the determination of the coincidence summing correction factors (TCSs). For this purpose a simple experimental technique was applied. The corrections were estimated by measuring of a sample, containing the radionuclides of interest with summing effect and one single-emitter radionuclide as reference, in two detector geometries (near and far) [16,17]. The coincidence summing correction factor is defined by the relations $TCS = {\binom{R_r}{R_s}}_h / {\binom{R_r}{R_s}}_g$ where R_r is the count rates of the single-emitter radionuclide of reference and R_s is the count rates of the radionuclide to correct. The ratios ${\binom{R_r}{R_s}}_h$ and ${\binom{R_r}{R_s}}_g$ are calculated at height *h* from the detector where coincidence summing is negligible and over the detector (indicated by *g*), respectively. This method is simple to use and it is independent of the sample activity which contribute to minimize several uncertainties.

Water solution of unknown activity with the following radionuclides of interest: ¹³⁴Cs, ¹⁵²Eu, ⁶⁰Co and ¹³⁷Cs, was used. For the reference we used ¹³⁷Cs in the middle energy of the spectra (661.7 keV). For detector D1 the sample was measured at 10 cm from the end cap and over the detector while for detector D2 the 100 sample was measurement out of the well (at 6 cm from the top of 101 the well) and into the well of the detector. In order to calculate the 102 corrections factors, the TCS expression was applied to the following 103 energy lines: ¹³⁴Cs (604.7 and 795.7 keV), ¹⁵²Eu (344.3 keV) and 104 ⁶⁰Co (1173.2 and 1332.5 keV). It should be noted that no correction 105 was applied to ²²⁶Ra and ²³²Th because the available water source 106 did not contain these radionuclides. All measurements were carried 107 out for about 6-10 h ensuring a meaningful statistics in each source 108 to detector configuration and in all cases the dead time corrections 109 were intrinsically taken into account by the software WinnerTM. 110

2.4. Experimental verification

The experimental verification included the internal validation of MC efficiency curves and the participation in two ALMERA Proficiency Tests: IAEA-TEL-2013-04 and IAEA-TEL-2014-04 [18].

The validations of the efficiency calibration curves obtained by 117 Monte Carlo simulation was performed with high purity KCl and 118 several Certified Reference Materials (CRMs) provided by IAEA 119 [19]: Uranium Ore IAEA-RGU-1, soil IAEA-375, IAEA-326 and IAEA-120 444, marine sediments IAEA-300 and IAEA-306, and grass IAEA-121 373. For detector D1 the samples were encapsulated in a cylind-122 rical container of polystyrene of 1.8 mm thickness with an internal 123 diameter of 72 mm and filling height of 20 mm. The measurement 124 geometry of detector D2 was a cylindrical polystyrene vial of 1 mm 125 thickness with an internal diameter of 10 mm and filling height of 126 45 mm. Both capsules do not allow radon diffusion. In the sample 127 IAEA-RGU-1, the radionuclides ²²⁶Ra and ²¹⁰Pb are in secular 128 equilibrium. The specific activity of ⁴⁰K was determined con-129 sidering the elemental weights for KCl and ⁴⁰K natural abundance. 130 131

The 2013 ALMERA proficiency test (PT) consisted of three water 131 samples and one flour sample. The participating laboratories were 132

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