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Multiparton interactions: From pp to pA☆

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Abstract

In the process of understanding nuclear collisions, reliable extrapolations from pp collisions, based on Glauber models, are highly desirable, though seldomly accurate. We review the inclusion of diffractive excitations and argue that they provide an important contribution to centrality observables in pA collisions. We present a method for distinguishing between diffractively and non-diffractively wounded nucleons, and a proof-of-principle for an extrapolation of multiparton interaction models built on this.

Keywords: QCD, Nucleus collisions, Fluctuations, Glauber models, Diffraction

1. Introduction

An important step towards fully understanding signals of QGP formation and collectivity in heavy ion collisions, is providing realistic extrapolations of the dynamics of pp collisions. Collisions of protons with nuclei is an important stepping stone, as the full nuclear geometry is already involved here, but the situation remains somewhat simpler than a full AA collision, as the number of sub-collisions is equal to the number of wounded nucleons in the target.

In ref. [1] we argued that the approximations normally used when extrapolating pp dynamics to pA collisions are too crude. We will present inclusion of fluctuations to the Glauber formalism, giving rise to a "wounded" cross section with contributions from diffractive excitations. We compare inclusion of fluctuations calculated in the DIPSY model with those from the Glauber–Gribov model, and use them to calculate distributions of wounded nucleons at LHC energies. Finally we present a simple model based on these principles, which allows for extrapolation of multiparton interaction models to pA, and comparisons to data.

2. Including fluctuations in pA collisions

2.1. Fluctuations in proton–proton

The DIPSY [2] model is a dynamic initial state model, built on the Mueller dipole model [3]. The model includes dynamics up to LL-BFKL, plus additional corrections from satuation and momentum conservation[1]. The initial state is built up through evolution from an initial proton consisting of three valence dipoles. The evolution is in impact-parameter space and rapidity, with the dipole splitting probability per unit rapidity:

$$\frac{d\mathcal{P}_g}{dY} = \frac{N_c\alpha_s}{2\pi^2} d^2\boldsymbol{x}_g \frac{(\boldsymbol{x}_1 - \boldsymbol{x}_2)^2}{(\boldsymbol{x}_1 - \boldsymbol{x}_g)^2(\boldsymbol{x}_g - \boldsymbol{x}_2)^2}. \quad (1)$$

An emission produces two new dipoles, $(\boldsymbol{x}_1, \boldsymbol{x}_g)$ and $(\boldsymbol{x}_g, \boldsymbol{x}_2)$. The interaction probability between two

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[1]The DIPSY initial state model is implemented in a full event generator, with final state radiation from the Ariadne shower [4]. Hadronization is carried out by PYTHIA8 [5], with added coherence effects from rope hadronization [6]. For more information, visit http://home.thep.lu.se/DIPSY.

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dipoles, one from the left-moving cascade and one from the right-moving cascade:

$$P = \frac{\alpha_s^2}{4}\left[\ln\left(\frac{(\boldsymbol{x}_1-\boldsymbol{x}_3)^2(\boldsymbol{x}_2-\boldsymbol{x}_4)^2}{(\boldsymbol{x}_1-\boldsymbol{x}_4)^2(\boldsymbol{x}_2-\boldsymbol{x}_3)^2}\right)\right]^2. \quad (2)$$

Using the optical theorem, cross section can now be calculated. With convenient normalization, the optical theorem in impact parameter space reads:

$$\Im(A_{el}) = \frac{1}{2}(|A_{el}|^2 + P_{abs}). \quad (3)$$

Where "abs" is short for "absorption", *i.e.* inelastic non-diffractive contributions. By inserting from equation (2), we obtain the real elastic amplitude in impact parameter space, including all fluctuations in projectile and target:

$$T(b) \equiv -iA_{el} = 1 - \exp\left(-\sum_{ij} f_{ij}\right). \quad (4)$$

Since fluctuations are related to diffraction through the Good-Walker formalism, calculation of several semi-inclusive proton-proton cross sections is possible with equation (4). Here the absorptive, single diffractive and double diffractive:

$$\frac{d\sigma_{abs}}{d^2b} = 2\langle T(b)\rangle - \langle T(b)\rangle^2, \quad (5)$$

$$\frac{d\sigma_{SD,(p|t)}}{d^2b} = \left\langle \langle T\rangle^2_{(t|p)}\right\rangle_{(p|t)} - \langle T\rangle^2_{p,t}, \quad (6)$$

$$\frac{d\sigma_{DD}}{d^2b} = \left\langle T^2\right\rangle_{p,t} - \left\langle \langle T\rangle_t^2\right\rangle_p - \left\langle \langle T\rangle_p^2\right\rangle_t + \langle T\rangle^2_{p,t}. \quad (7)$$

where subscripts p and t indicates averages over projectile and target respectively. The DIPSY formalism applies for protons and nuclei alike, and has recently been applied directly to pA collisions [7]. In this work DIPSY is only used to calculate fluctuations in the proton transverse structure, and a simpler model is used for extrapolation to pA (see section 2.3).

2.2. Parametrization of cross section fluctuations

When extrapolating from pp to pA collisions, it is important to keep in mind that many measurements will rely on a centrality measure, *e.g.* particle production in the forward (Pb going) direction. The relevant semi-inclusive cross section for an interaction between projectile and target will therefore have contributions from absorptive, single diffractive and double diffractive processes. We dub this the "wounded" cross section, inspired by the wounded nucleon model by Białas *et al* [8]:

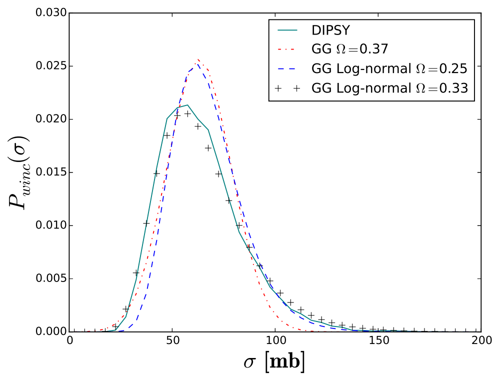

Figure 1: Fluctuations in σ_w for DIPSY and three versions of GG fluctuations. GG fluctuations using a log-normal parametrization of $P_{tot}(\sigma)$ seems to be able to describe the DIPSY fluctuations best.

$$\frac{d\sigma_w}{d^2b} = \frac{d\sigma_{abs}}{d^2b} + \frac{d\sigma_{SD,t}}{d^2b} + \frac{d\sigma_{DD}}{d^2b} = 2\langle T\rangle_{p,t} - \left\langle \langle T\rangle_t^2\right\rangle_p. \quad (8)$$

We can now compare the fluctuations in σ_w produced by DIPSY with the often used parametrization, Glauber–Gribov Colour Fluctuations (GG) [9]. Here the fluctuations are parametrized with a distribution $P_{tot}(\sigma)$ such that:

$$\sigma_{tot} = \int d\sigma \sigma P_{tot}(\sigma) \quad (9)$$

$$= \int d\sigma \rho \frac{\sigma^2}{\sigma+\sigma_0}\exp\left[-\frac{(\sigma/\sigma_0-1)^2}{\Omega^2}\right], \quad (10)$$

where the usual choice for $P_{tot}(\sigma)$ has been inserted in equation (10). The parameters of the model can be fitted to semi-inclusive cross sections by assuming a functional form for T. Here we use a semi-transparent disk with:

$$T(b,\sigma) = T_0\Theta\left(\sqrt{\frac{\sigma}{2\pi T_0}} - b\right). \quad (11)$$

The semi-inclusive cross sections are:

$$\sigma_{tot} = \int d^2b \int d\sigma P_{tot}(\sigma) 2T(b,\sigma) \quad (12)$$

$$\sigma_{el} = \int d^2b \left|\int d\sigma P_{tot}(\sigma) T(b,\sigma)\right|^2 \quad (13)$$

$$\sigma_w = \int d^2b \int d\sigma P_{tot}(\sigma)\left[2T(b,\sigma) - T(b,\sigma)\right] \quad (14)$$

In figure 1 we show fluctuations in σ_w at $\sqrt{s_{NN}} = 5.02$ TeV with DIPSY as well as GG, compared to GG with an modified P_{tot} distribution – a log-normal distribution – which we find describes the DIPSY fluctuations better:

$$P_{tot}(\sigma, b) = \frac{1}{\Omega\sqrt{2\pi}}\exp\left(-\frac{\log^2(\sigma/\sigma_0)}{2\Omega^2}\right). \quad (15)$$